

# Probing Electronic Structure in Novel Engineered Quantum States of Matter

*In situ ARPES Studies of Epitaxial Novel Oxides Thin Films*

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# Acknowledgement

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❖ Mingying Li, Haifeng Yang and Zhengtai Liu, and Hongping Mei  
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*Fudan University*

– The set-up of OMBE/ARPES systems at SIMIT and partial work on SrRuO<sub>3</sub> epitaxial thin films

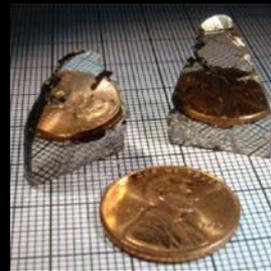
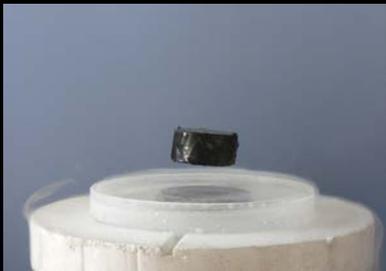
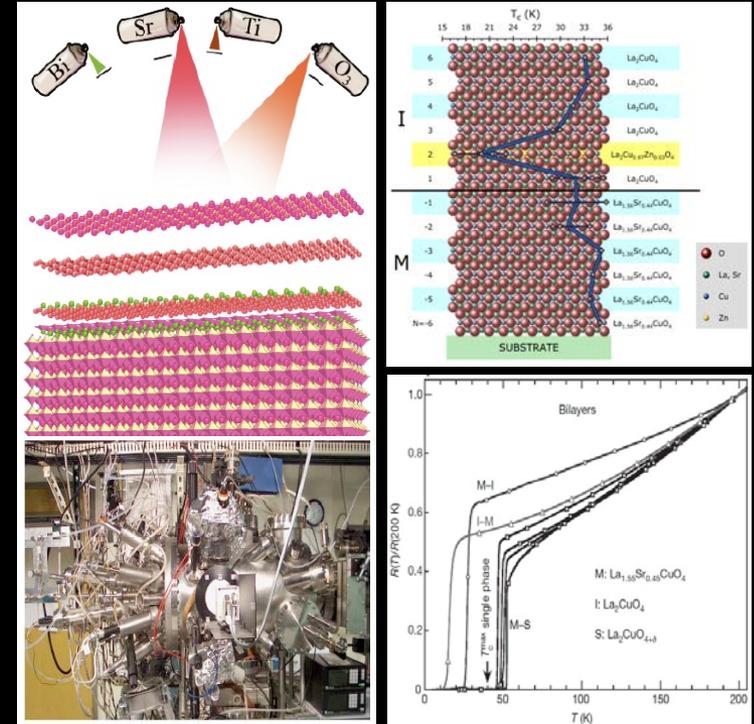
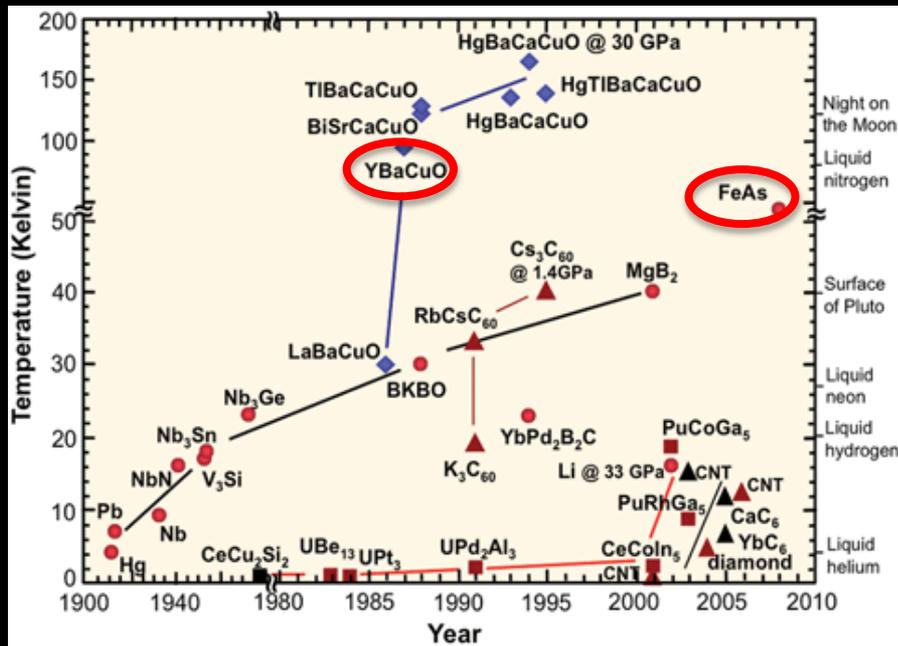
– The design and set-up of OMBE/ARPES systems at Cornell and most work on SrRuO<sub>3</sub> epitaxial thin films

– The work on the dead-layer of LSMO/STO interface

# The Trend: Discovering and Engineering Novel Quantum Matters

“At the extreme forefront of research in superconductivity is the empirical search for new materials”

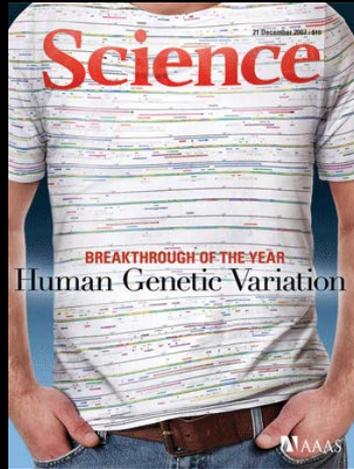
Bednorz and Mueller ('86)



Atomic-scale layer-by-layer growth

G. Logvenov et al., Science 326, 699 (2009)  
A. Gozar et al., Nature 455, 782 (2008)

# Rational Materials Design: Desired Correlated-electron properties

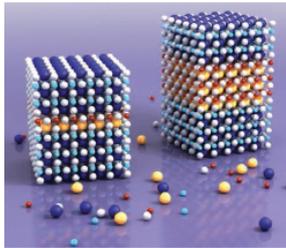


**5 BEYOND SILICON?** Sixty years ago, semiconductors were a scientific curiosity. Then researchers tried putting one type of semiconductor up against another, and suddenly we had diodes, transistors, microprocessors, and the whole electronic age. Startling results this year may herald a similar burst of discoveries at the interfaces of a different class of materials: transition metal oxides.

Transition metal oxides first made headlines in 1986 with the Nobel Prize-winning discovery of high-temperature superconductors. Since then, solid-state physicists keep finding unexpected properties in these materials—including colossal magnetoresistance, in which small changes in applied magnetic fields cause huge changes in electrical resistance. But the fun should really start when one oxide rubs shoulders with another.

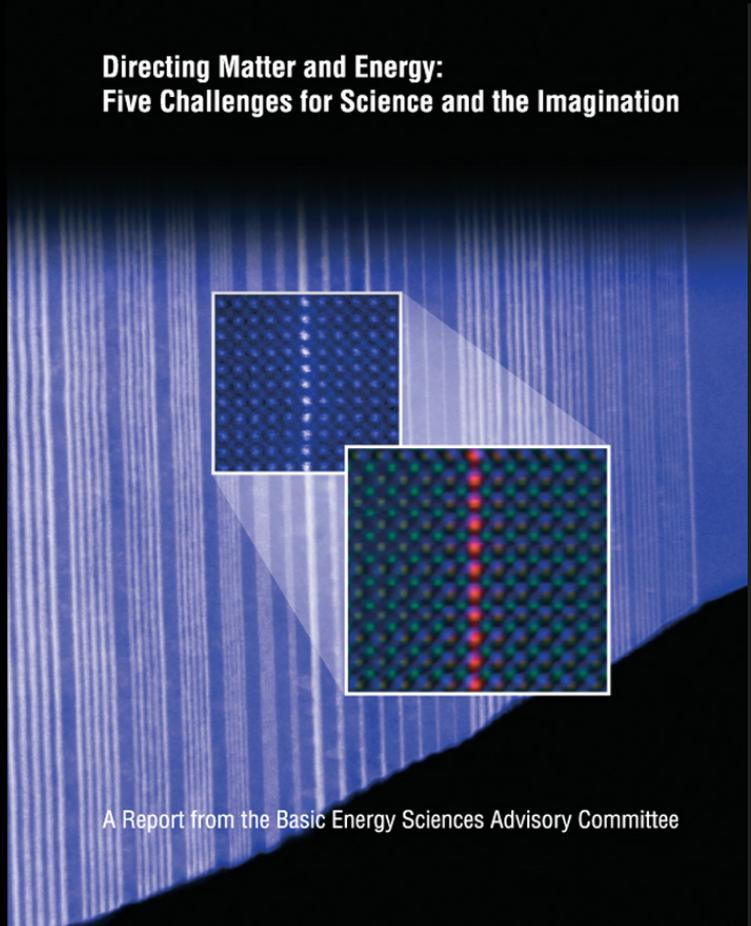
If different oxide crystals are grown in layers with sharp interfaces, the effect of one crystal structure on another can shift the positions of atoms at the interface, alter the population of electrons, and even change how

**Tunable sandwich.** In lanthanum aluminate sandwiched between layers of strontium titanate, a thick middle layer (*right*) produces conduction at the lower interface; a thin one does not.



**Rational Materials Design:** How can we control the properties of matter emerging from complex correlations **by design**?

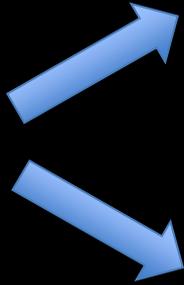
Directing Matter and Energy:  
Five Challenges for Science and the Imagination



A Report from the Basic Energy Sciences Advisory Committee

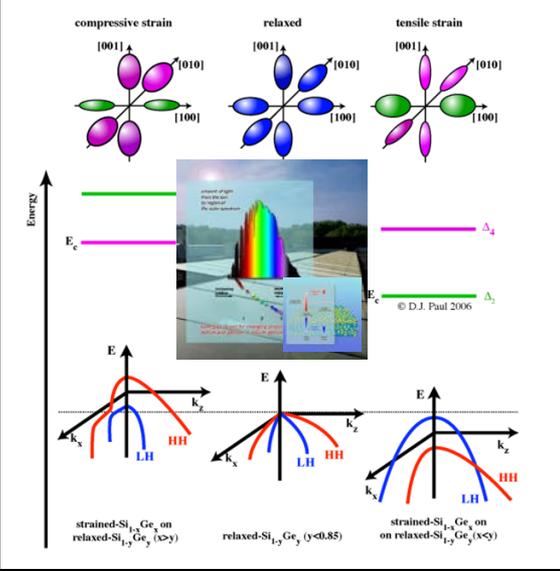
# Great Success of Band-gap engineering in conventional semiconductors

**C**alculation of physical properties

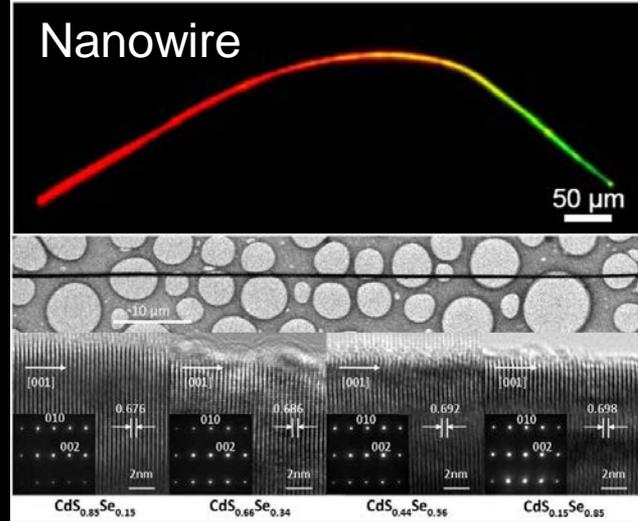


**C**onstruction of desired materials

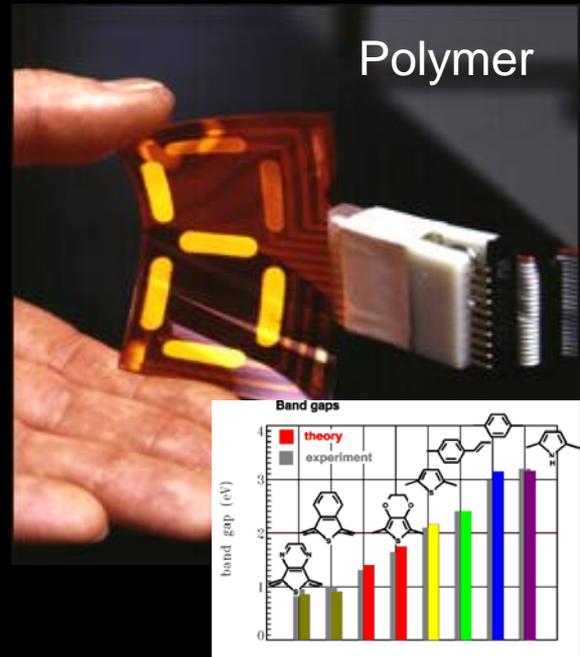
**C**ontrol of properties through **epitaxial strain & chemical doping etc.**



Strain band gap engineering

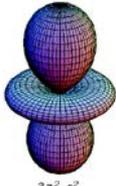
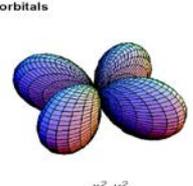
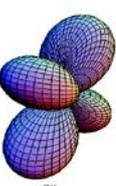
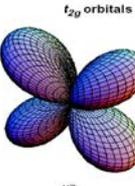
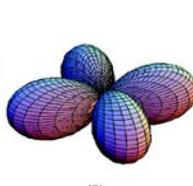


Quantum confinement effects



Polymer

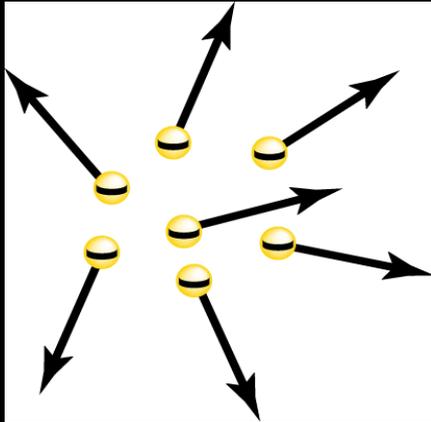
# Correlated Transition Metal Oxides

H											He							
Li	Be	    					B	C	N	O	F	Ne						
Na	Mg											Al	Si	P	S	Cl	Ar	
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr	
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe	
Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn	
Fr	Ra	Ac	Rf	Ha	106	107	108	109	110	111	112							
			Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu		
			Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr		

$$\mathcal{H} = \underbrace{-\sum_j^{N_e} \frac{\hbar^2}{2m} \nabla_j^2}_{e^- \text{ KE}} - \sum_{\alpha}^{N_i} \frac{\hbar^2}{2M_{\alpha}} \nabla_{\alpha}^2 - \sum_j^{N_e} \sum_{\alpha}^{N_i} \frac{Z_{\alpha} e^2}{|\vec{r}_j - \vec{R}_{\alpha}|} + \underbrace{\sum_{j \ll k}^{N_e} \frac{e^2}{|\vec{r}_j - \vec{r}_k|}}_{e^- / e^- \text{ Int.}} + \sum_{\alpha \ll \beta}^{N_j} \frac{Z_{\alpha} Z_{\beta} e^2}{|\vec{R}_{\alpha} - \vec{R}_{\beta}|} \quad \text{Nuclei / Nuclei Int.}$$

# *Innovation in the design of correlated quantum matters is needed!*

## Non-Interacting



### We understand well:

- Systems of non-interacting or weakly interacting particles;
- Interaction energy is smaller than the kinetic energy, perturbation theory works well.
- DFT do nice predictions.

## Correlated

**C**alculation of physical properties

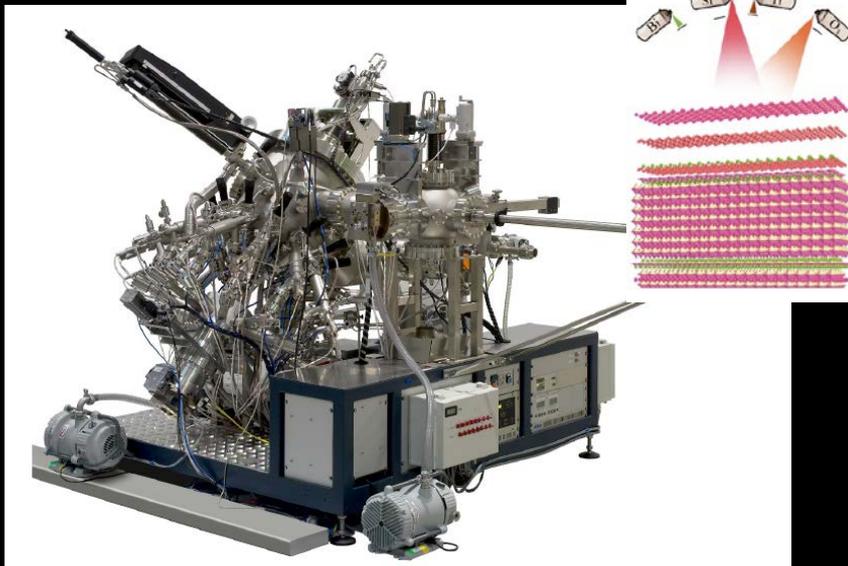
**C**onstruction of desired materials

**C**haracterization of **microscopic interactions** by **high-resolution experiments**

**C**ontrol of properties through **epitaxial strain & chemical doping**

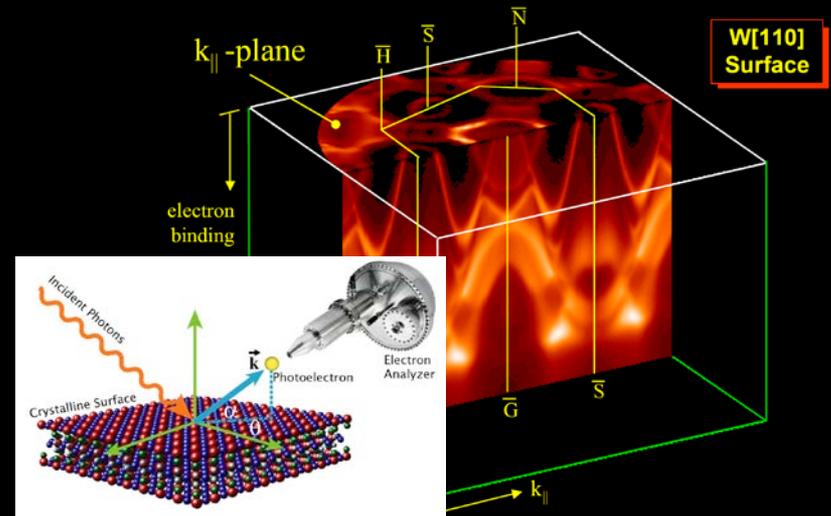
# Outline

1. New tools for investigating novel quantum electronic states
2. Correlation in epitaxial Ruthenate Thin Films
3. Dead-layer at the  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3/\text{SrTiO}_3$  interface



Oxide Molecular Beam Epitaxy

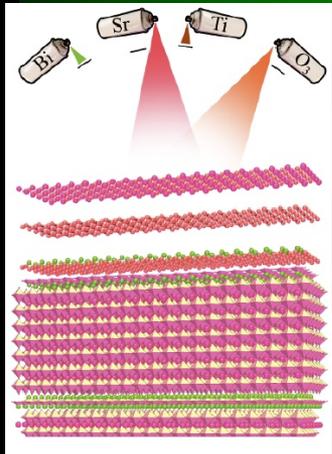
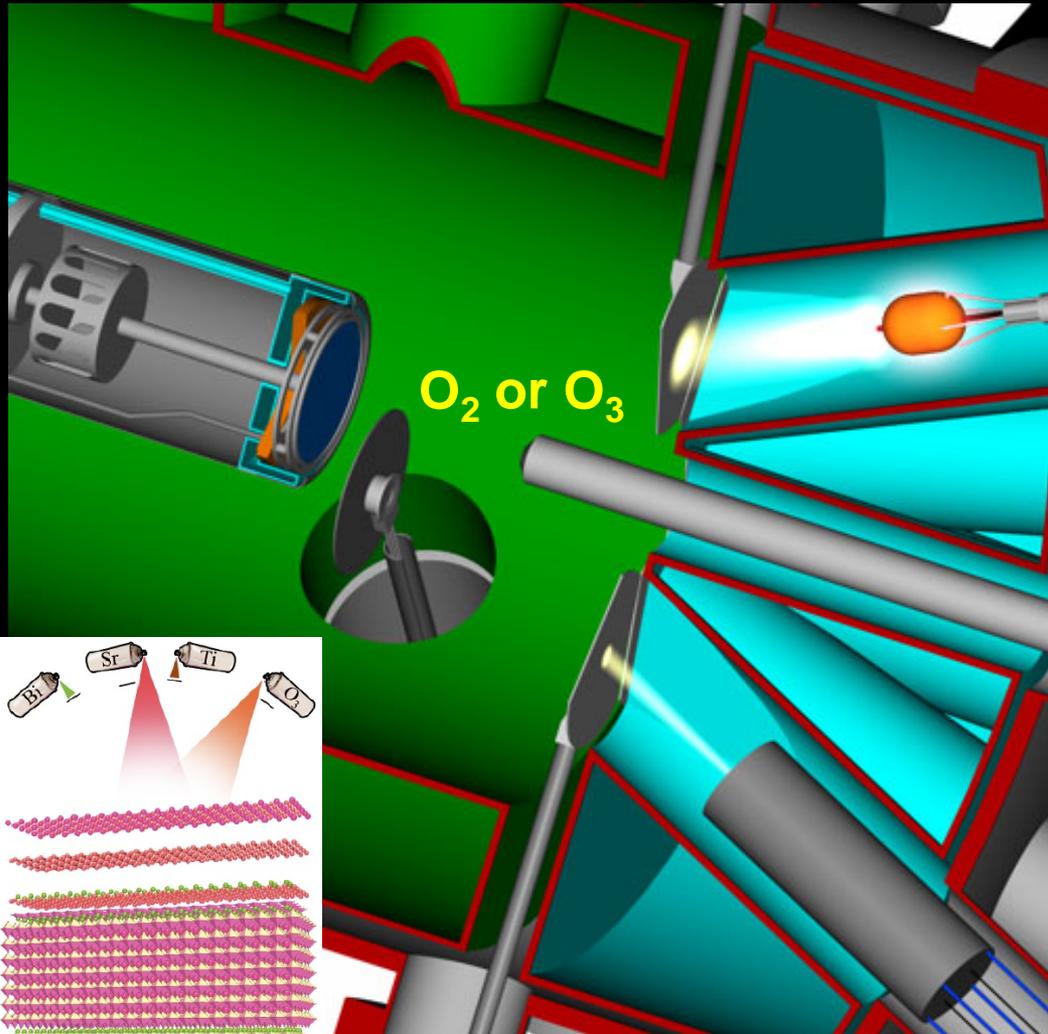
## Band Mapping and Fermi Contours



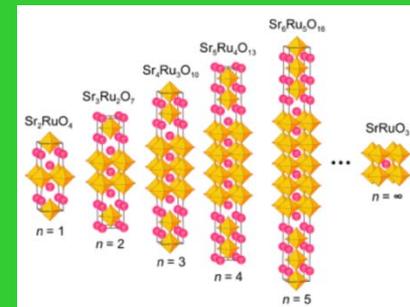
Angle-Resolved Photoemission Spectroscopy

# **Oxide** Molecular Beam Epitaxy : “Atomic Spray Painting”

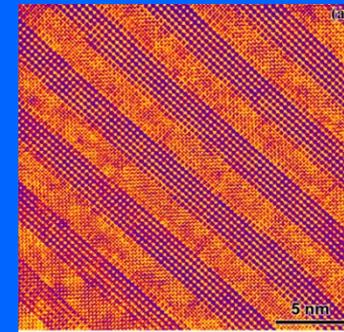
Shuttered growth utilizes alternating molecular beams so that each element is evaporated separately.



**RHEED oscillation (<1%)**



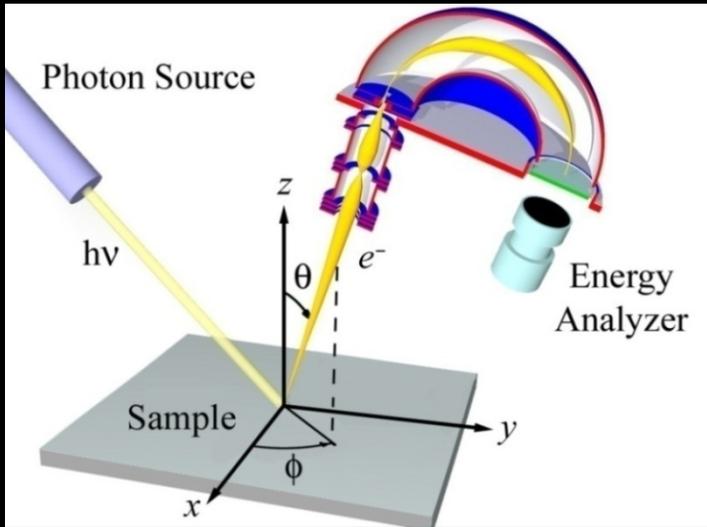
**New Phase**



**Superlattice, Interface**

Courtesy of Dr. Carolina Adamo

# Characterization by Angle-Resolved Photoemission Spectroscopy



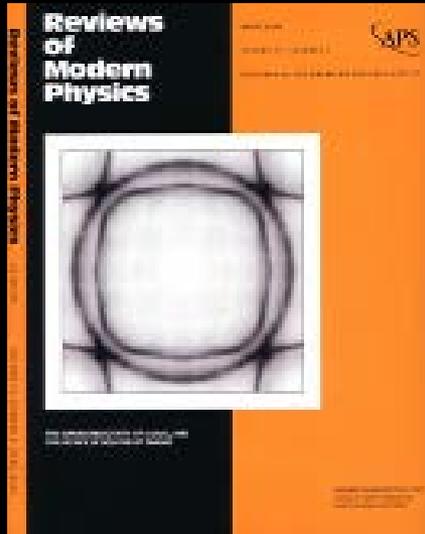
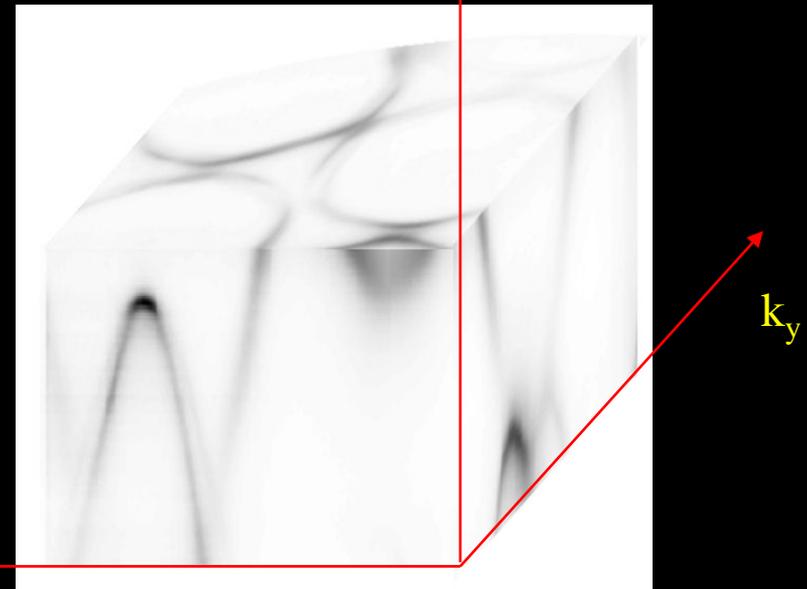
From momentum/energy conservation rules:

$$k_{pe} = -k_{N-1}$$

$$h\nu - E_{pe} = E_{N-1} - E_N$$

<http://en.wikipedia.org/wiki/ARPES>

Direct Mapping of "Band"  $E$   $\text{Sr}_2\text{RhO}_4$



Courtesy of Prof. Donglai Feng

[http://arpes.stanford.edu/research\\_strontium\\_ruthnates.html](http://arpes.stanford.edu/research_strontium_ruthnates.html)

## Limitations of ARPES

Highly surface sensitive  
(atomically clean surfaces)

Mostly restricted to cleavable  
samples (need well-defined  
layer-structured crystals)

Requires sizable single  
crystals  
~ mm-sized

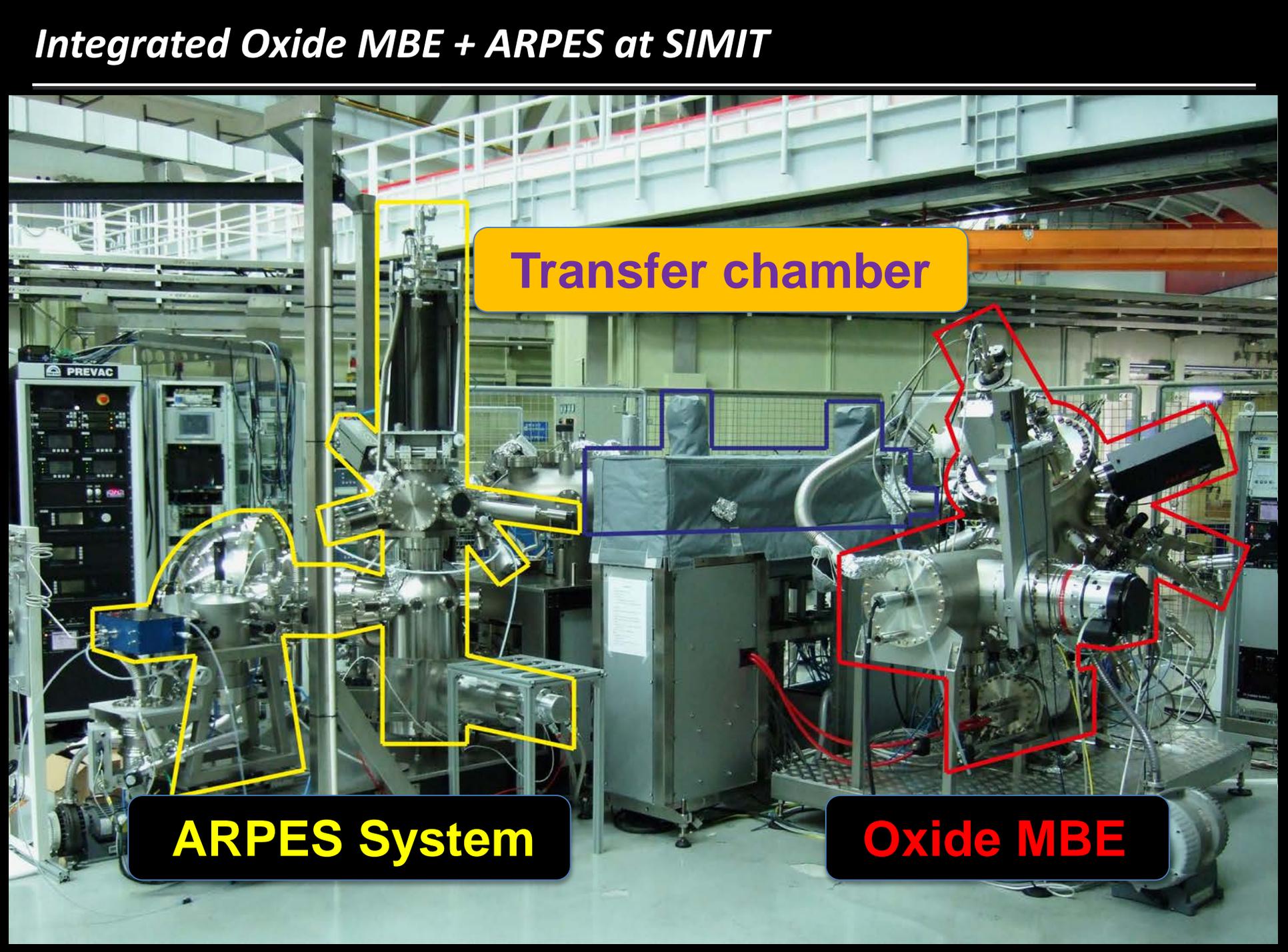
## Advantages of OMBE

Films grown under ultra-  
high vacuum conditions,  
not energetic deposition

Epitaxial growth allows  
for study of  
“uncleavable” materials

Allow for growth of  
“artificial” materials :  
heterostructures,  
interfaces, epitaxially  
strained materials

# *Integrated Oxide MBE + ARPES at SIMIT*

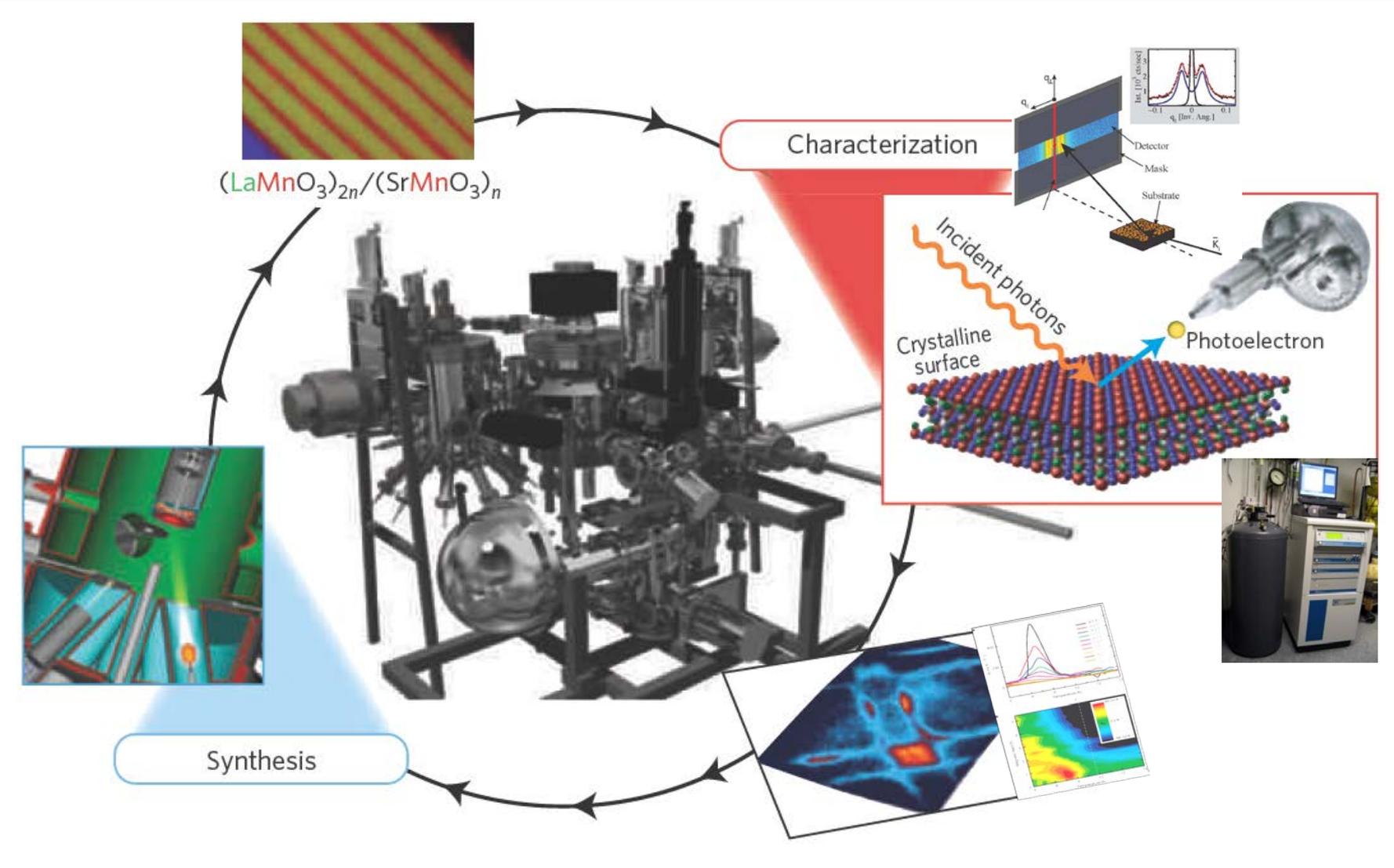


**Transfer chamber**

**ARPES System**

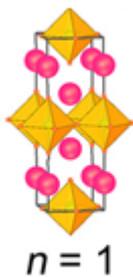
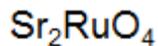
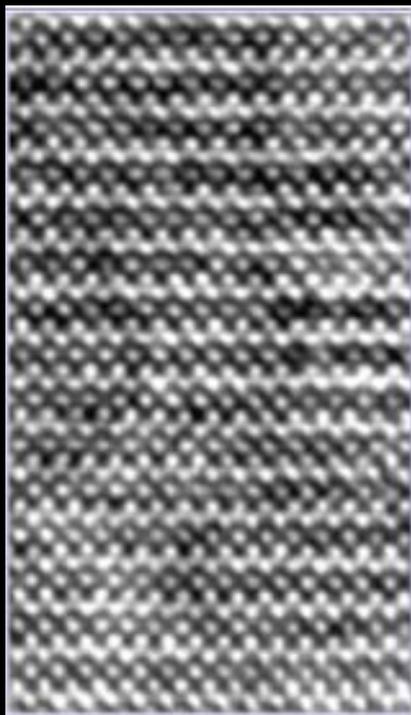
**Oxide MBE**

# The mode of rational materials design for complex quantum matters



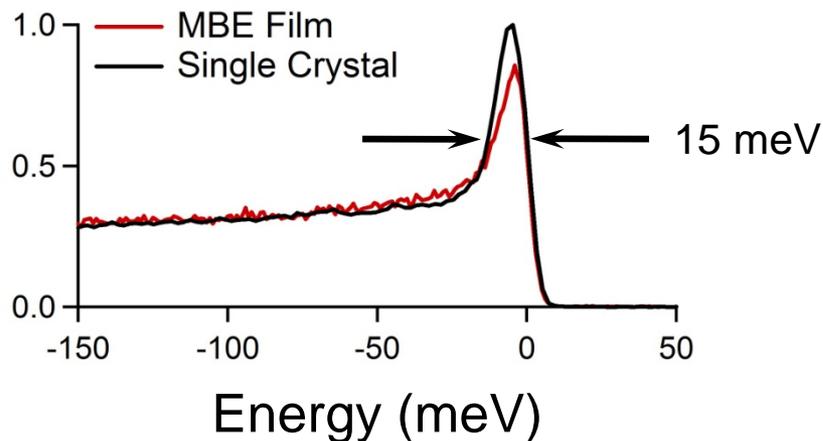
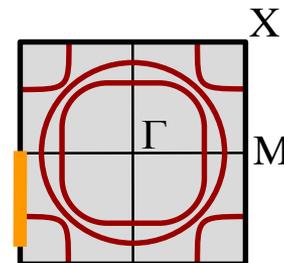
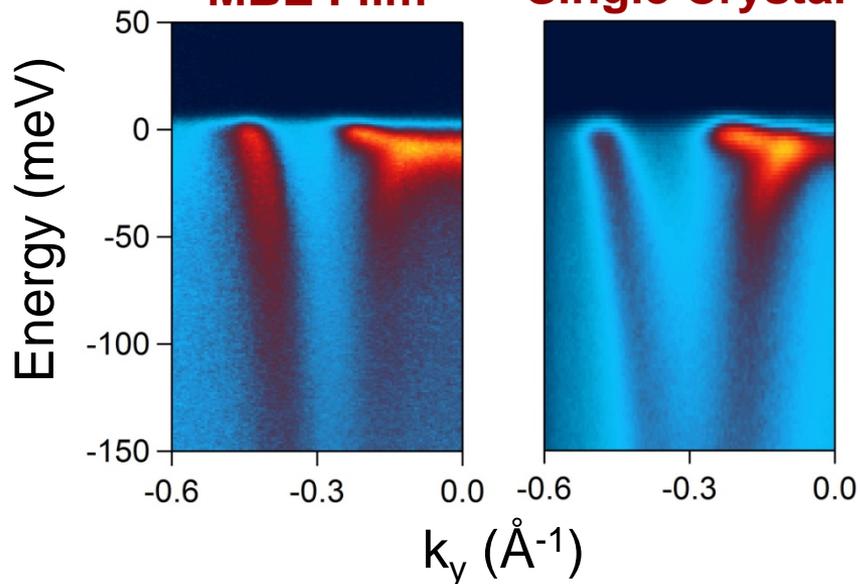
# Preliminary reliability testing of the integrated OMBE/ARPES

High-resolution transmission electron microscopy (HRTEM) image of  $\text{Sr}_2\text{RuO}_4$  films



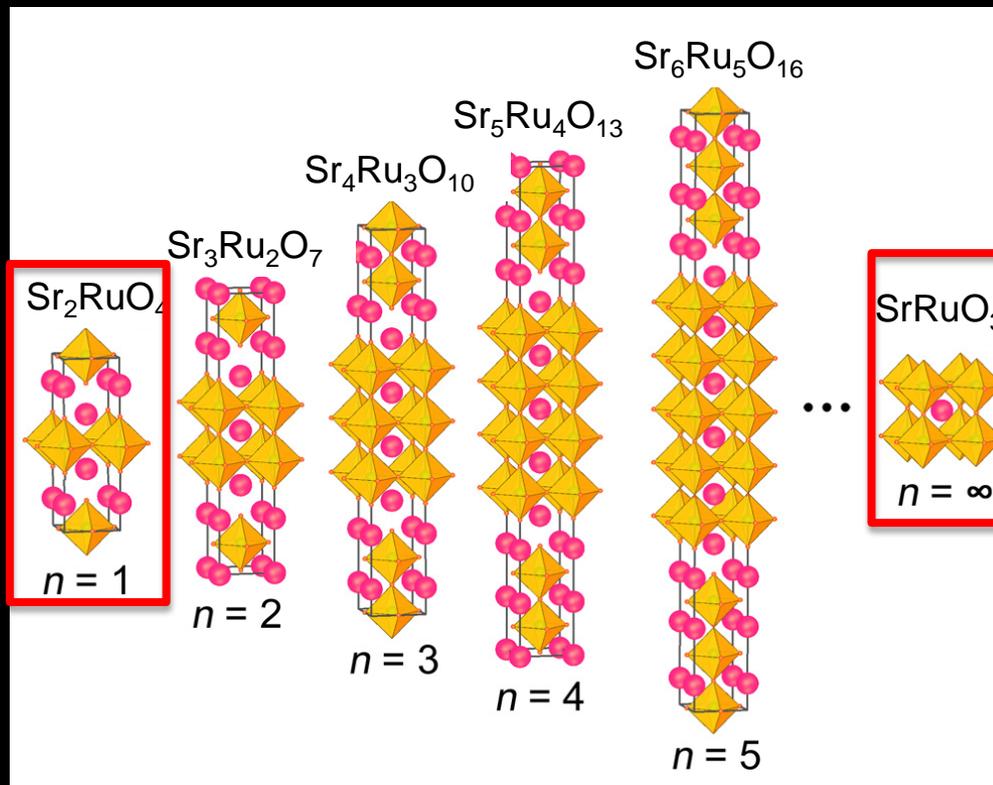
MBE Film

Single Crystal

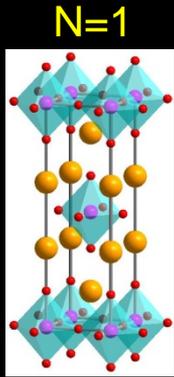


# Outline

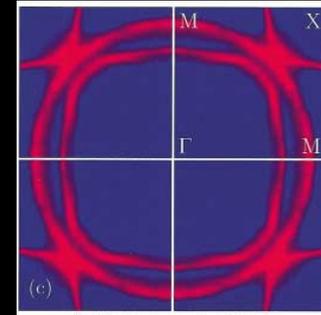
1. New tools for investigating novel quantum electronic states
2. Correlation in epitaxial ruthenate thin films
3. Dead-layer at the  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3/\text{SrTiO}_3$  interface



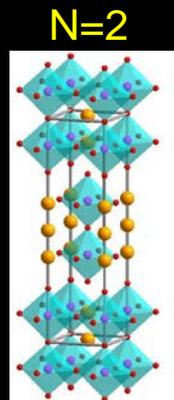
# Sr-based Ruthenates $SrO(SrRuO_3)_n$ : Rich physics



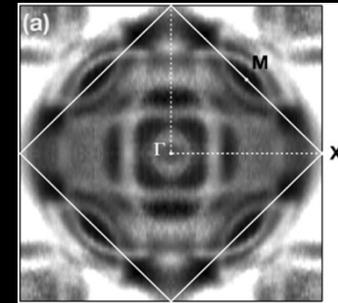
Spin triplet super-  
Conductivity,  
possible  
topological SC



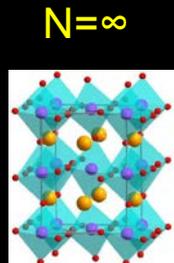
*Damascelli et al., PRL (2000)*



Electronic  
nematic phase



*Tamai et al., PRL (2008)*

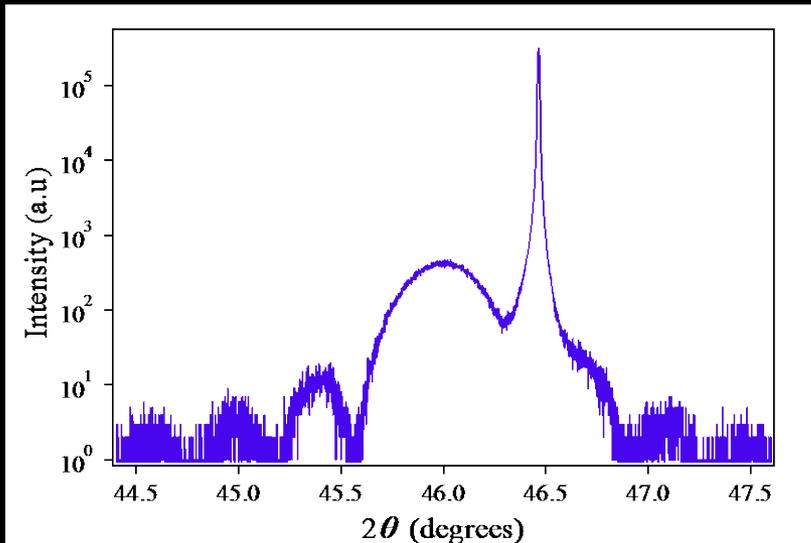


Significant  
correlations,  
Ferromagnetism  
(origin?)



# Characterization of SrRuO<sub>3</sub> Epitaxial films

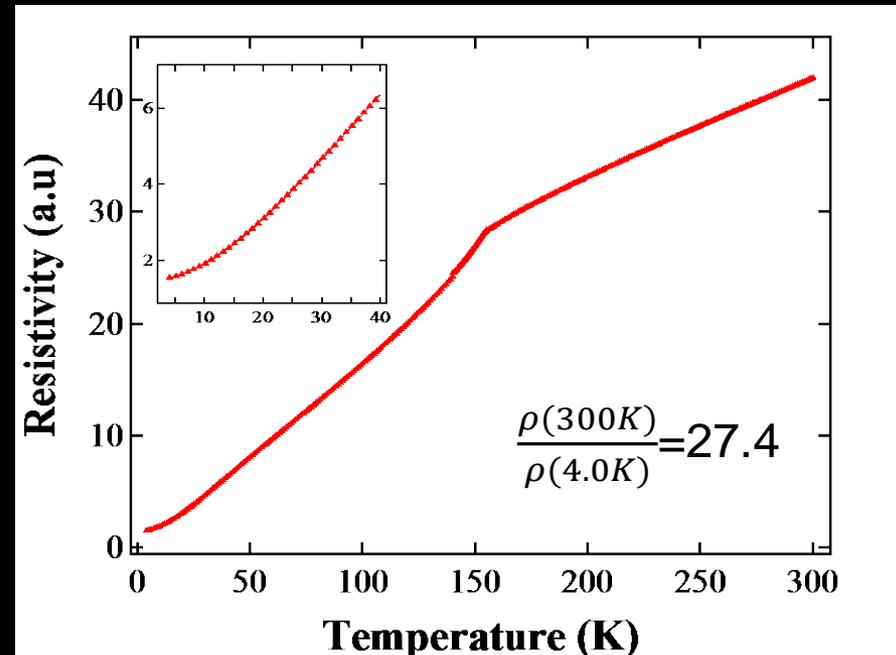
## X-ray diffraction



- The fringes indicate good crystalline quality of the films
- Right 2θ value of the (002) peak for strained SrRuO<sub>3</sub> films

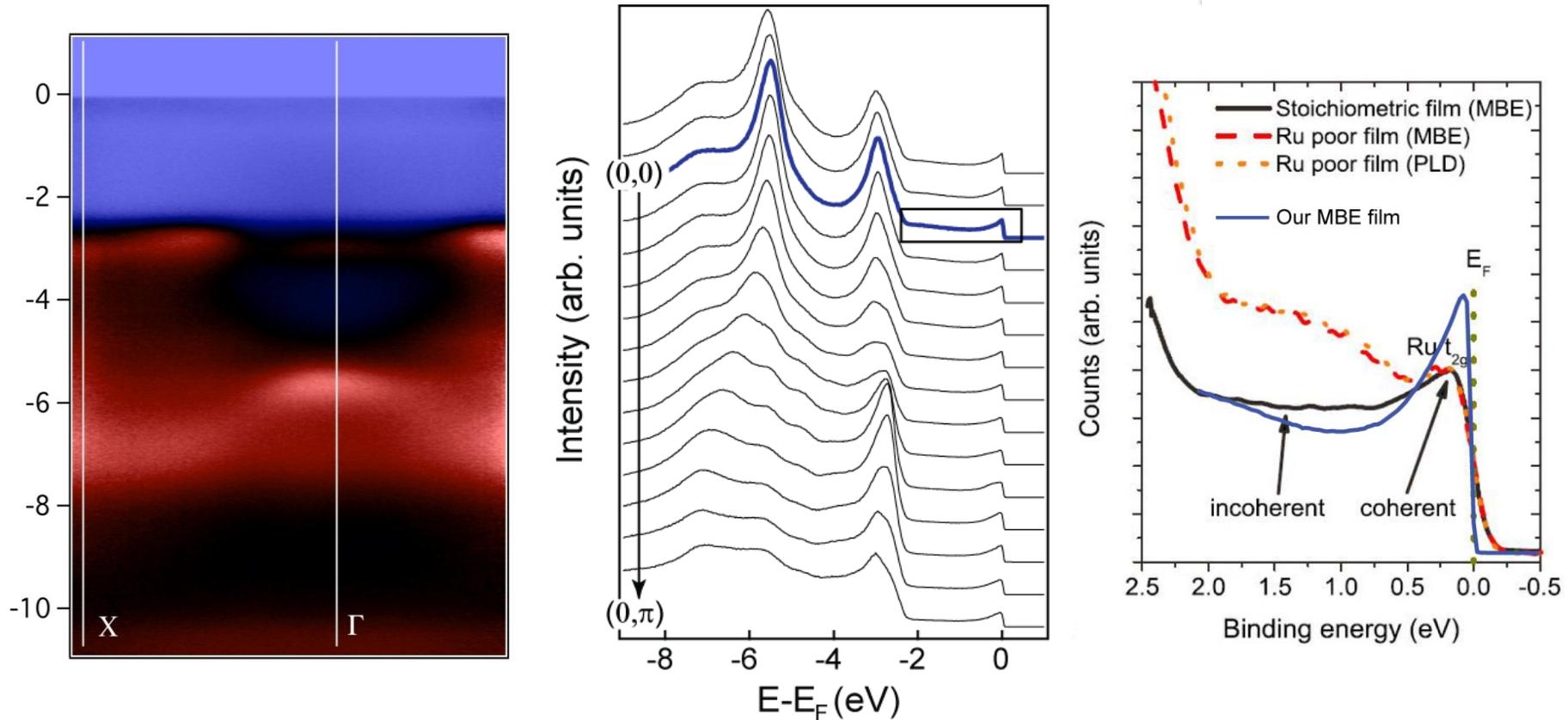
## Transport measurements

Residual resistivity ratio RRR ~27.4, the record high quality of SrRuO<sub>3</sub> thin film



**So far the best SrRuO<sub>3</sub> thin films (~20nm) reported**

# What does the electronic structure look like?



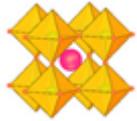
Siemons et al., PRB 76, 075126(2007)

**Photoemission spectra are sensitive to the stoichiometry**

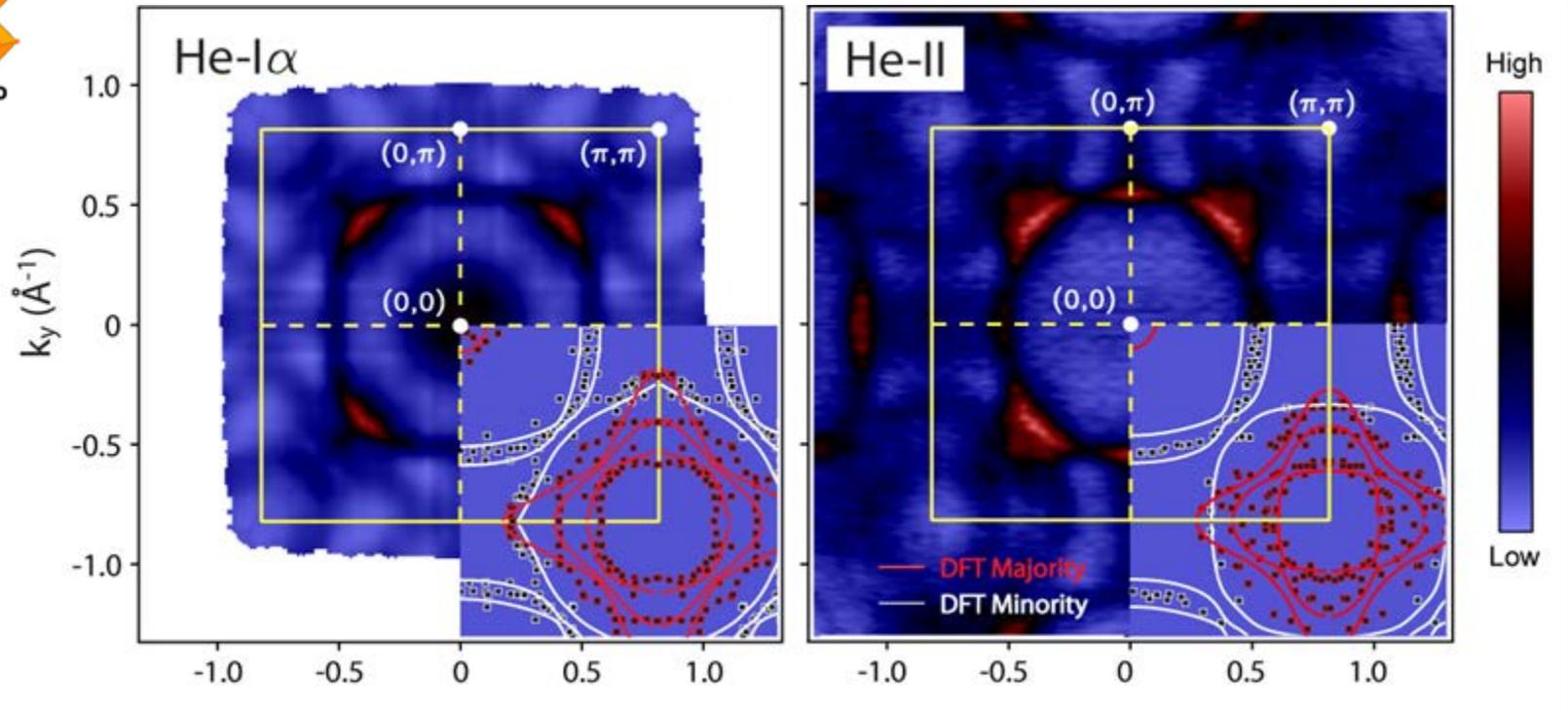
# The first probe of the $k$ dependent electronic structure

## Fermi Surface Topology Maps (4-fold symmetrized) $T=10\text{K}$ , FM phase

SrRuO<sub>3</sub>



$$\eta = \infty$$

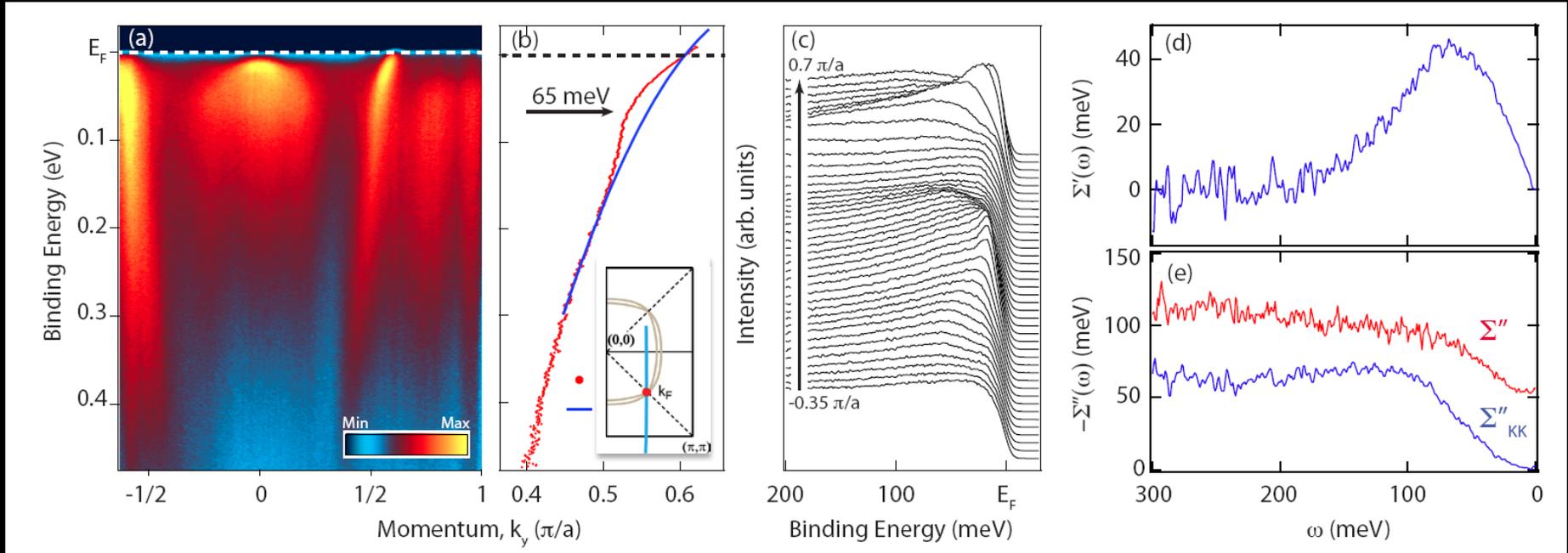


HeI $\alpha$  (21.2eV)

HeII (40.8eV)

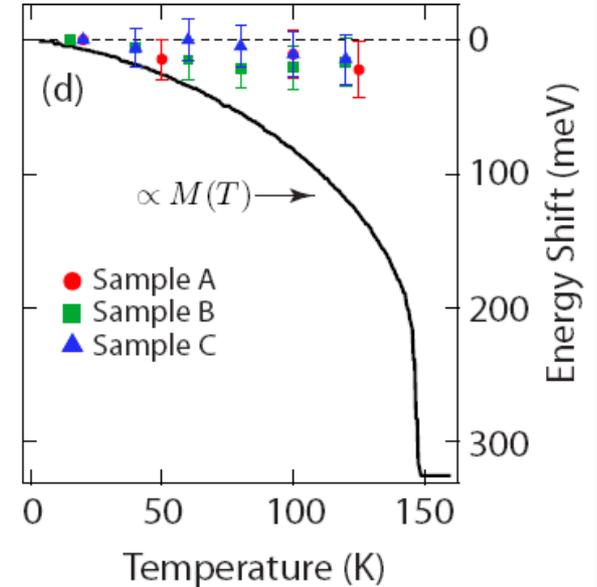
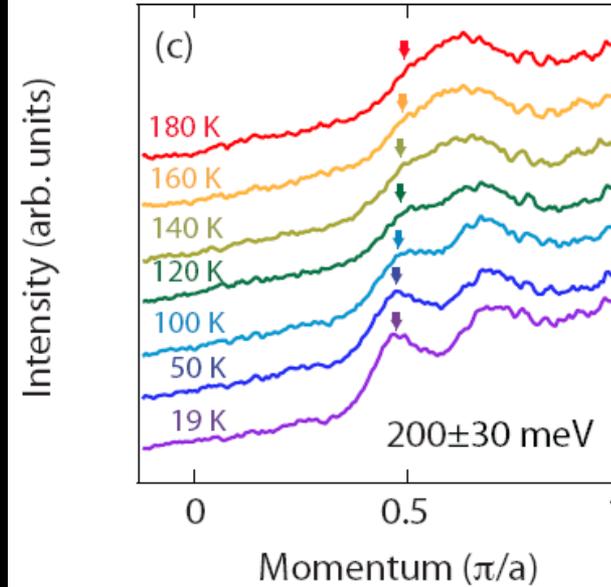
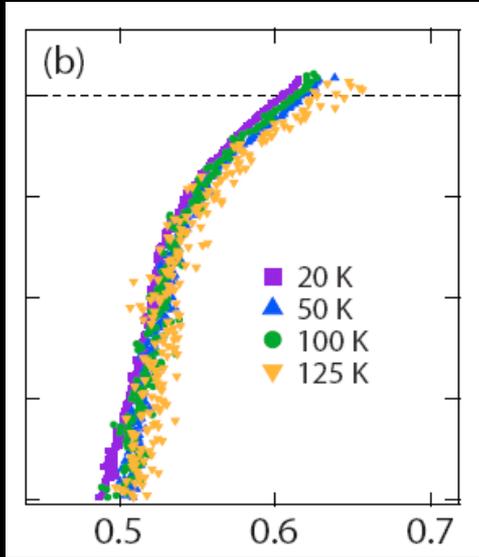
$k_F$ 's are consistent with both DFT prediction and previous bulk sensitive quantum oscillation measurements on SrRuO<sub>3</sub>.

# What is the origin of the strong correlation?



- Evident ‘kink’ at  $\sim 65$  meV: **signature of a strong coupling to some bosonic modes:**
  - Low binding energy: heavily dressed QP ( $M^* = 4.5 M_{LDA}$ ), in good agreement with dH-vA oscillations;
  - High binding energy: similar slope as DFT, **weak e-e correlation.**
- Naturally explain the long-standing puzzle about the strong correlation effects.

# What is the origin of the ferromagnetism in $\text{SrRuO}_3$ ?



- The Stoner model: the exchange splitting of an itinerant FM should decrease from its saturation value at low temperatures to zero at  $T_c$ .
- A strong deviation of the band dispersion from the Stoner model.
- The magnetism has a strong local character and cannot be explained in a simple picture of itinerant Stoner FM.

## ***Strong correlation in SrRuO<sub>3</sub> thin films***

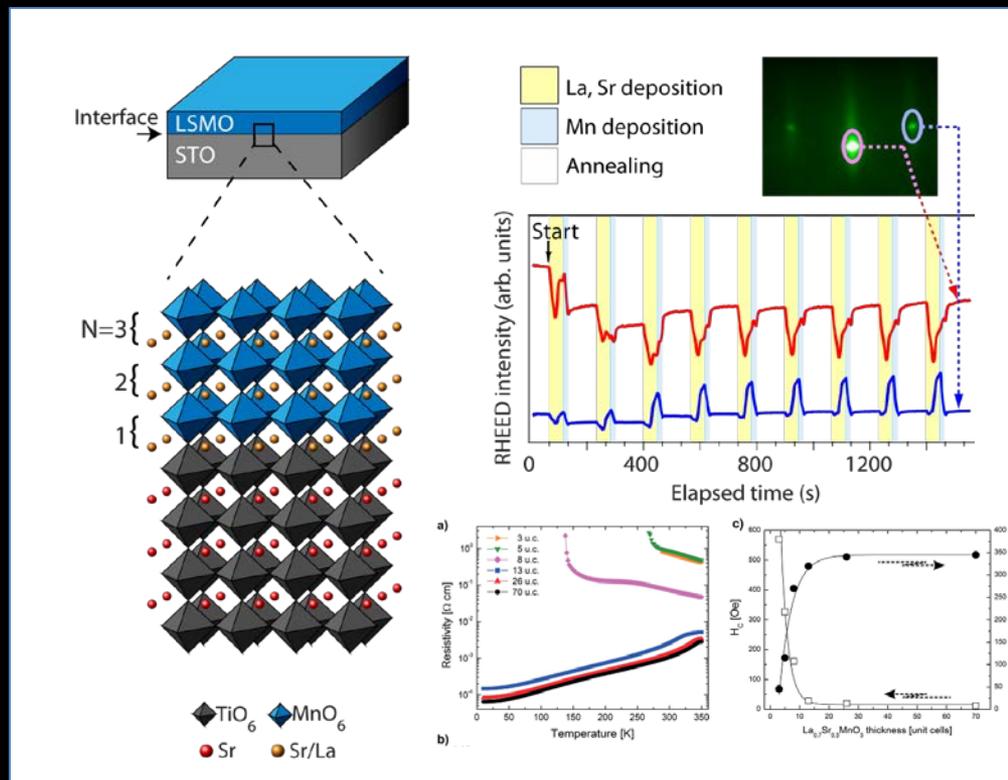
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- ***For the first time, high resolution ARPES is conducted on epitaxial SrRuO<sub>3</sub> thin films to extract the electronic structure:***
  - High accuracy Fermi surface maps, in agreement with DFT calculation;
  - Strong kink structure in QPs' dispersion
- ***Many-body interactions in SrRuO<sub>3</sub> thin films are anatomized:***
  - Strong coupling to some bosonic modes, e-ph? Magnons?
  - Relatively weak e-e correlations.
  - The magnetism has a strong local character.

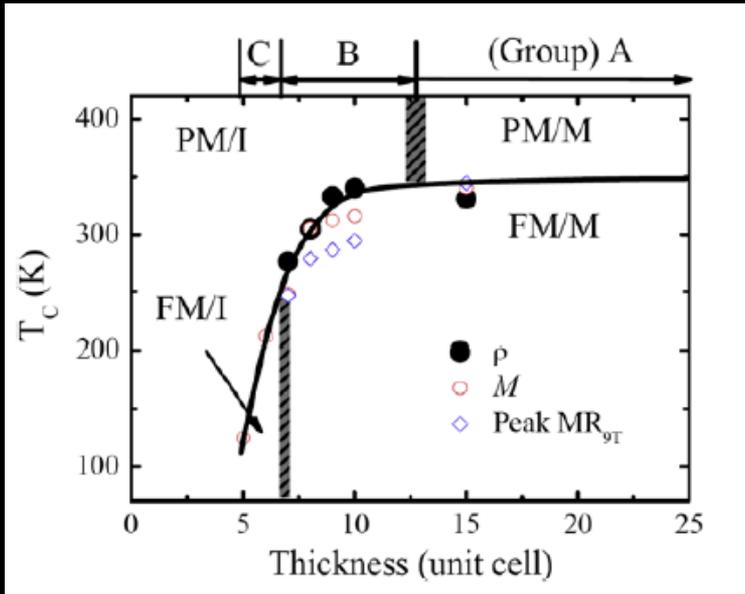
***Naturally explain the long-standing puzzle about the strong correlation effects.***

# Outline

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3. Dead-layer at the  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3/\text{SrTiO}_3$  interface



# Dead-layer phenomenon at $\text{La}_{0.67}\text{Sr}_{0.33}\text{MnO}_3/\text{SrTiO}_3$ interfaces

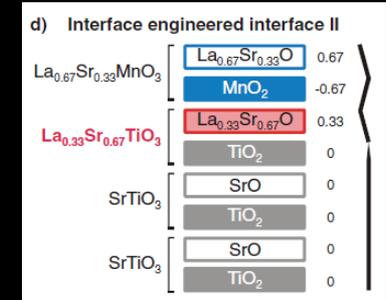
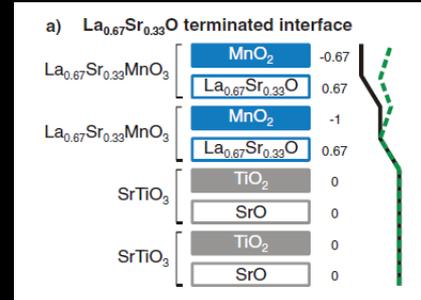


7 unit cells on LSMO/STO

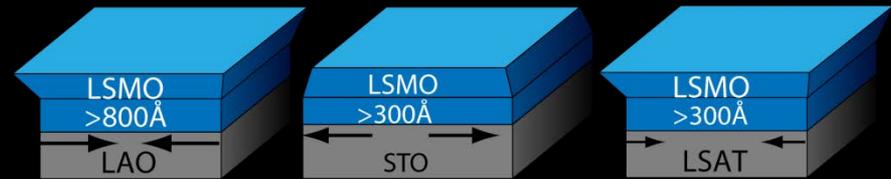
Dead layer—strongly affected by defects

## Possible mechanisms:

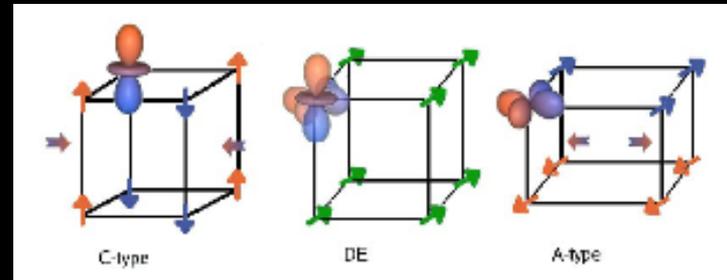
- Charge redistribution



- Strain effect



- Magnetic and orbital reconstruction



## Oxide MBE:

- Sharp interface, low energy
- low defect densities

Nakagawa, *Nature materials* (2006)

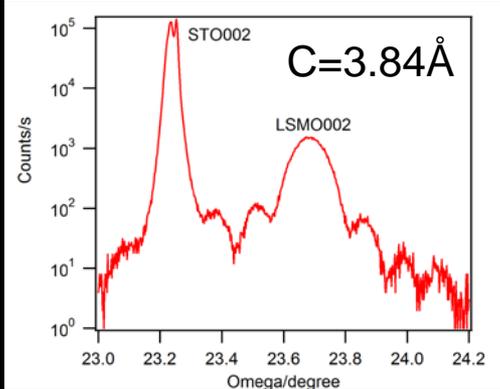
Tebano et al., *PRL* 100, 137401 (2008)

Kim et al., *Solid State Communications* 150(13–14): 598-601

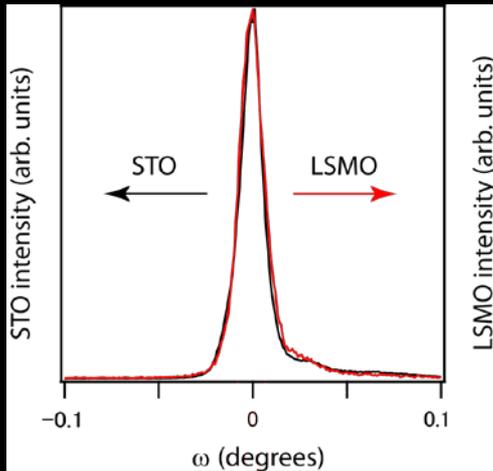
H. Boschker et al., *AdFuncMat* 22,2235 (2012)

# Characterization of $\text{La}_{0.67}\text{Sr}_{0.33}\text{MnO}_3$ thin films

## Fine scan around (002)

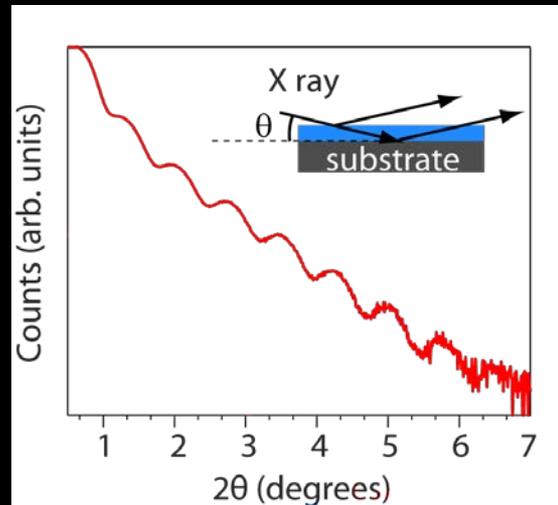


## Rocking around (002)



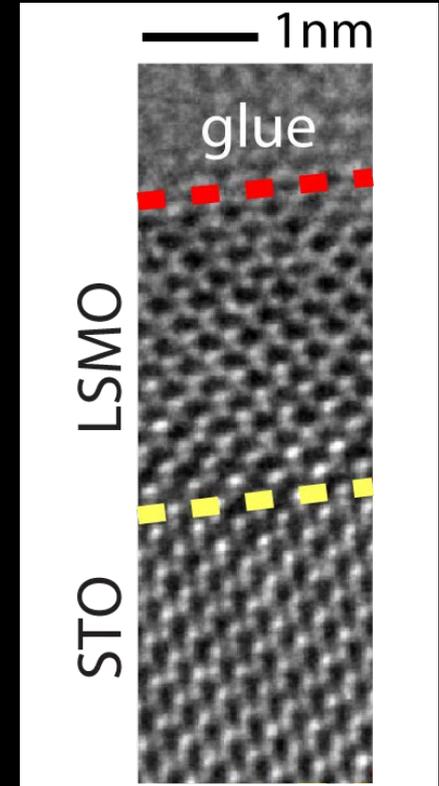
FWHM=0.015, limited by the substrate

## X-ray reflection



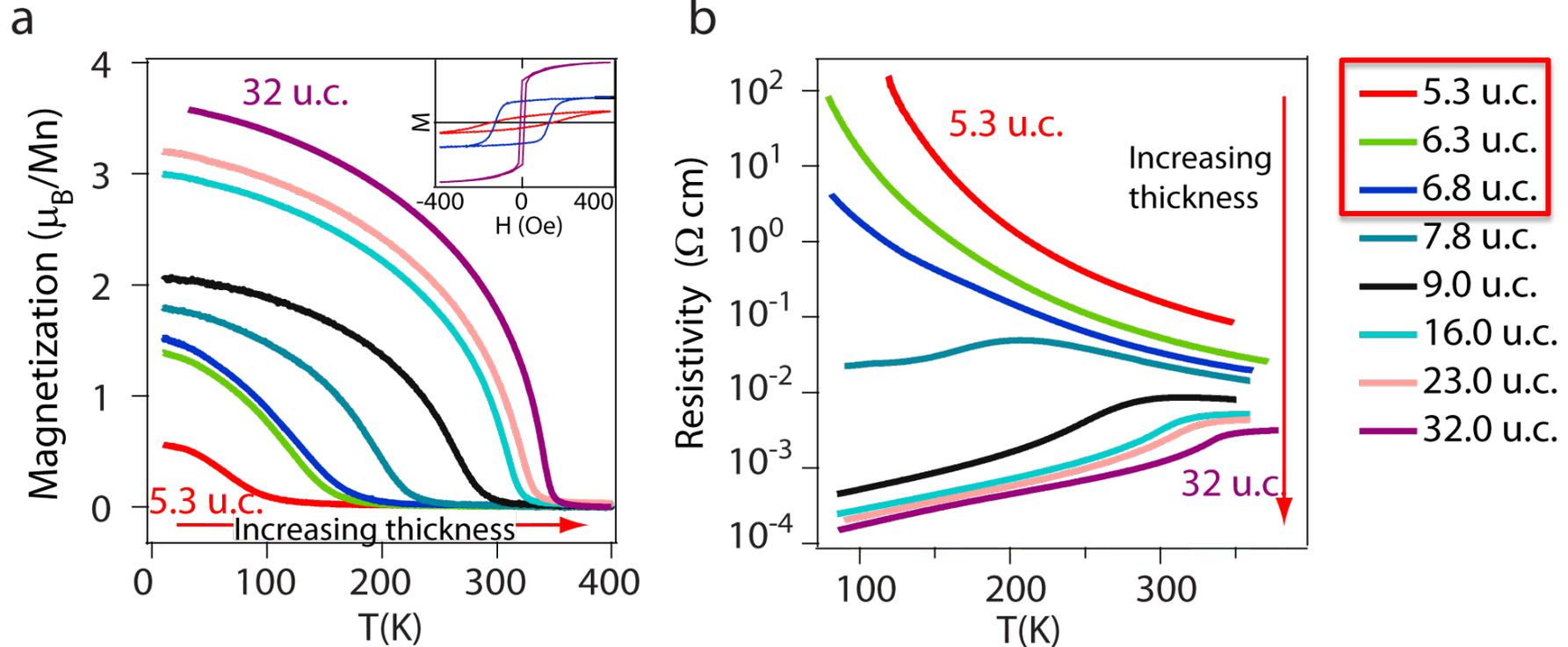
Absolute thickness error within 2%

## Cross-sectional TEM image



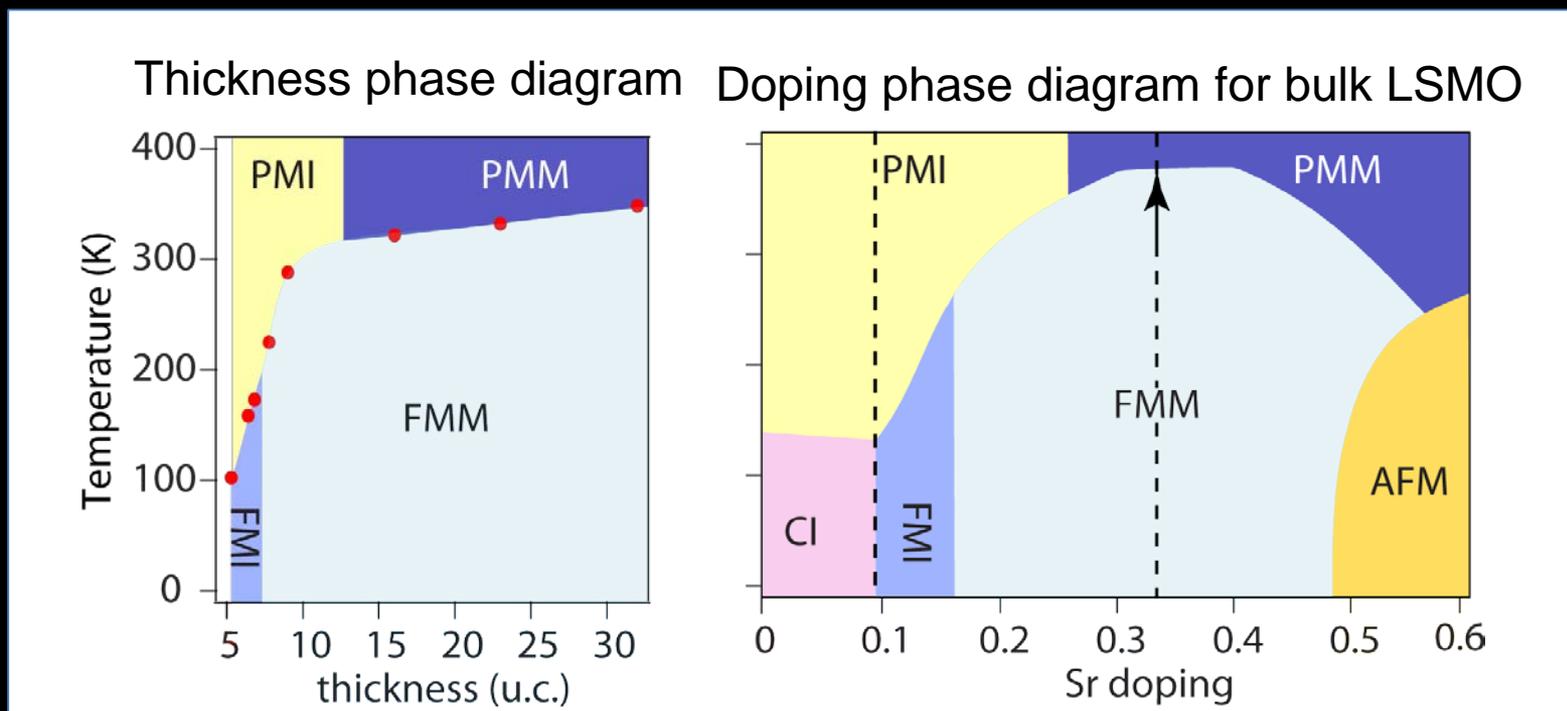
Confirm the 7 $\mu\text{c}$  thickness

# Transport properties v.s. thickness

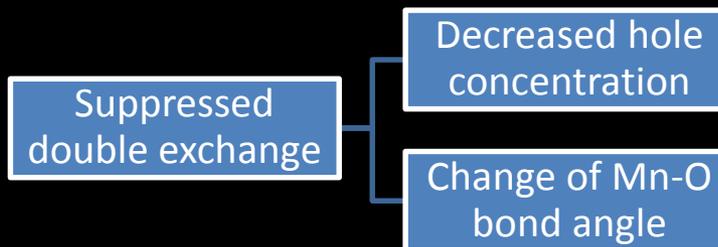


- 32 u.c. ~ bulk;
- Both of Curie temperature and conductivity drop upon decreasing thickness;
- Below 7u.c., the films are insulating below  $T_c$ , as a **typical dead-layer behavior**.

# Remarkable similarity in the phase diagram to the doped bulk

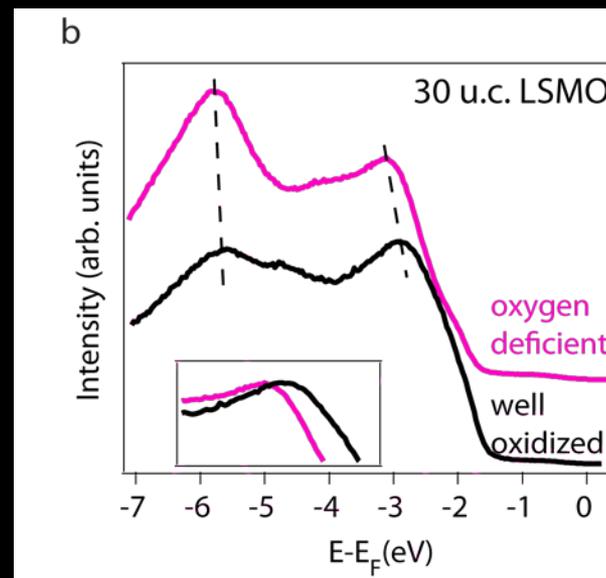
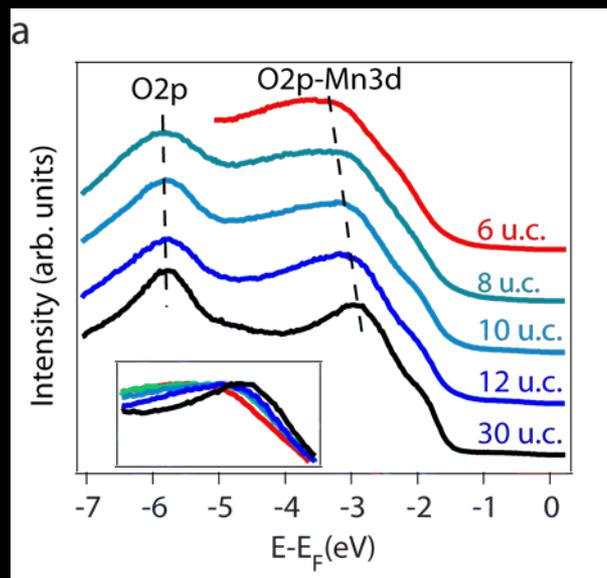


In bulk: double-exchange is suppressed → metal-insulator transition



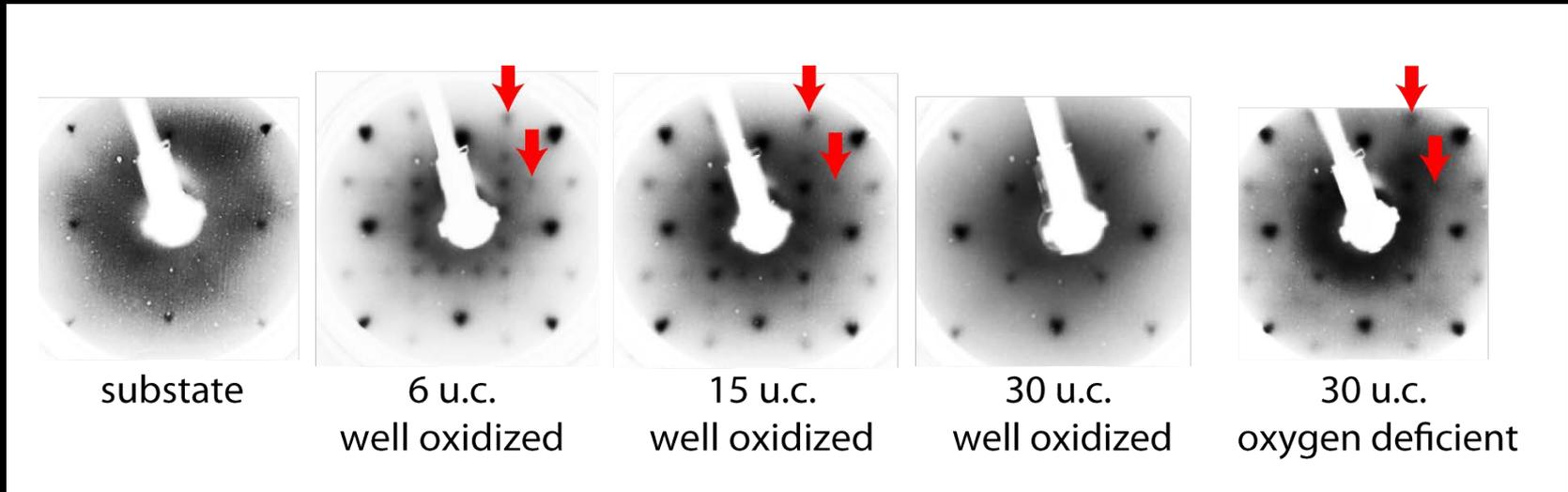
The same mechanism for ultrathin films?

# Oxygen vacancies $\rightarrow$ hole concentration and bond angle change



- The Mn-O hybridized band, shifts towards higher binding energy with decreasing thickness.
- Similar shift is observed in the 30 u.c. film intentionally grown under oxygen deficient condition
- **This valence band shift closely related to the oxygen deficiency**

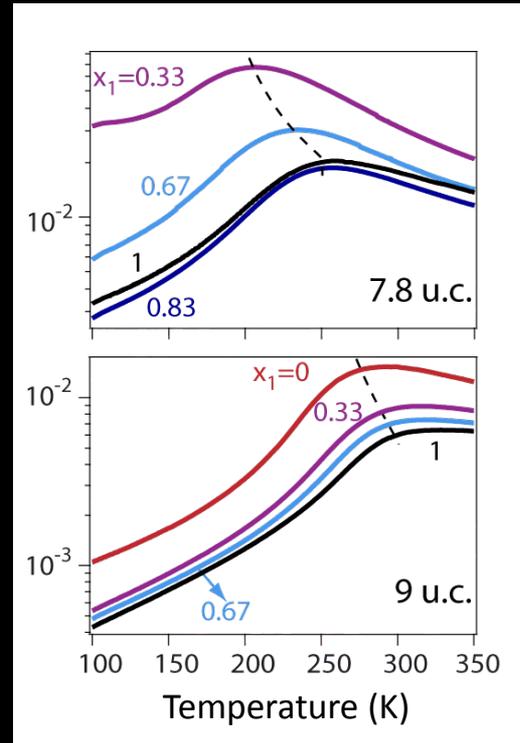
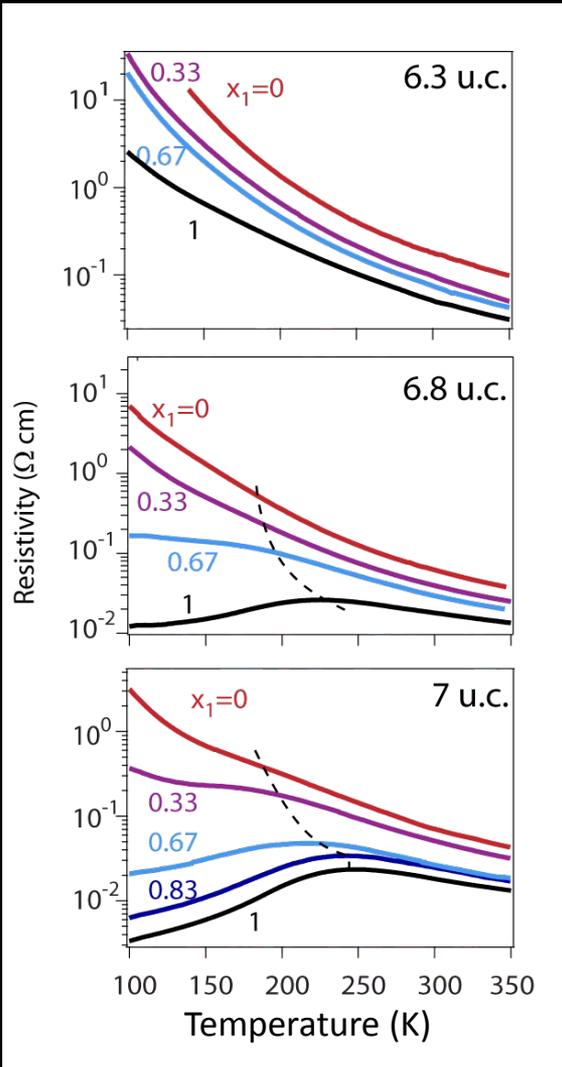
## Other evidence from LEED



- Both the ultra-thin films and the intentionally oxygen deficient thick film show an additional surface reconstruction.
- Further annealing in ozone did not cause any variation in the valence band or LEED patterns of the ultra-thin films.

**Intrinsic oxygen vacancy formation appear in ultra-thin films ----- change the carrier concentration, and suppress the double exchange**

# Interface engineering based on this model



The deteriorated metallicity and ferromagnetism with decreased thickness can be compensated by higher interfacial Sr doping.

The dead-layer is successfully reduced to 6 u.c.

## *Summary and Outlook*

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- ARPES is a powerful tool to study the strongly interacting solid state materials; BUT, it is greatly limited to the cleavable samples.
- MBE+ARPES technique can greatly extend our research scope to originally unachievable materials phase;
- It also provides a great tool to customize and characterize the oxide thin films with desired functions.

*Open  
ARPES?  
OMBE?*



**Thanks for your attention!**