

Spin-orbit coupling in Fe-based superconductors

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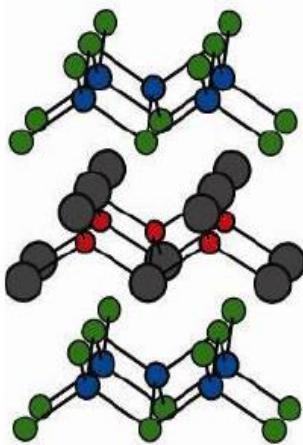
Outline

- **Introduction:** structure, Fermi surface, etc.
- **Spin-resonance as evidence for the s_{\pm} gap symmetry**
- **Experimental evidences for spin-space anisotropy**
 - Inelastic neutron scattering
 - NMR
- **Spin-orbit coupling**
 - Band structure renormalization
 - Fermi surface changes
 - Spin-resonance splitting
- **Conclusions**

1111 vs. 122 vs. 111 vs. 11 materials

Big family of pnictides and chalcogenides

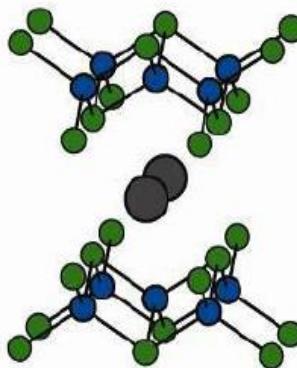
LaFeAsO



$T_c=26\text{K}$

(55K for Sm, Chen et al., Nature 2008; Ren Chin. Phys. Lett. 2008)

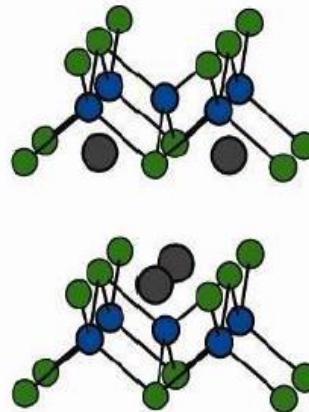
BaFe_2As_2



$T_c=38\text{K}$

Rotter et al., PRL 2008, Ni et al. PRB 2008

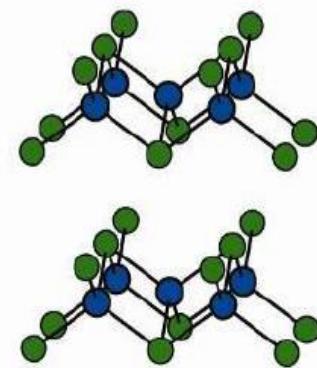
LiFeAs



$T_c=18\text{K}$

Wang et al., arXiv:0806.4688

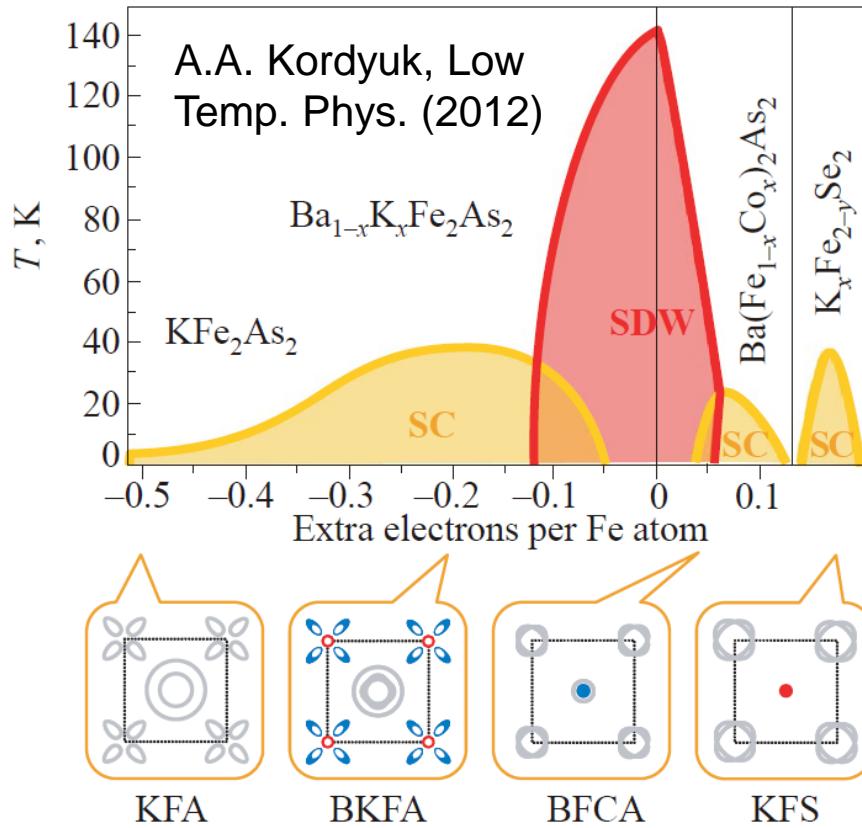
FeSe



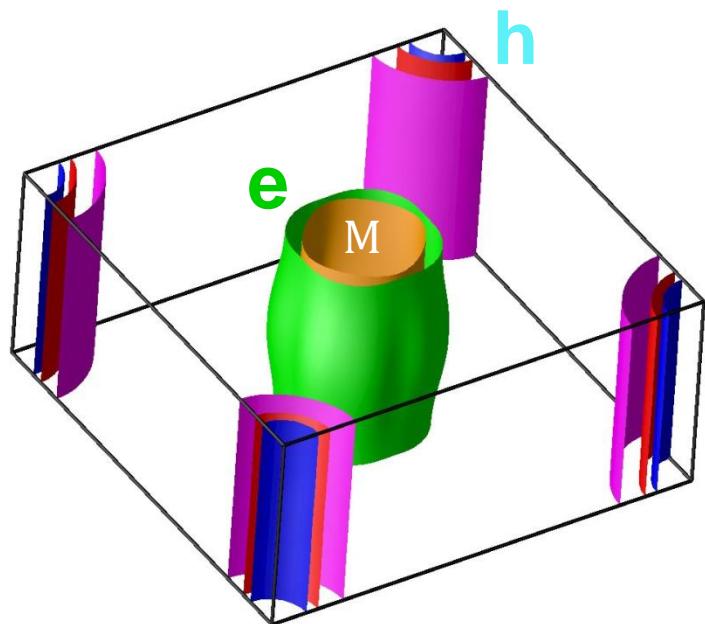
$T_c=8\text{K}$

Hsu et al., arXiv:0807.2369

Itinerant picture: LDA electronic structure (LaFeAsO)



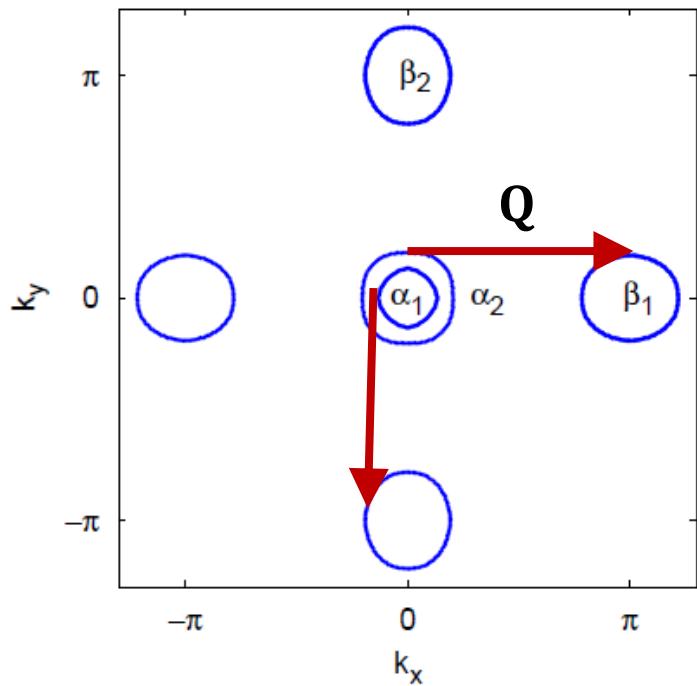
Weak CEF splitting: all 5 orbitals ($d_{x^2-y^2}$, $d_{3z^2-r^2}$, d_{xy} , $d_{xz}+d_{yz}$) are near the Fermi level



1. Both electron and hole pockets do exist (quantum oscillations, ARPES)
2. There is a Fermi surface even at zero doping (ARPES)

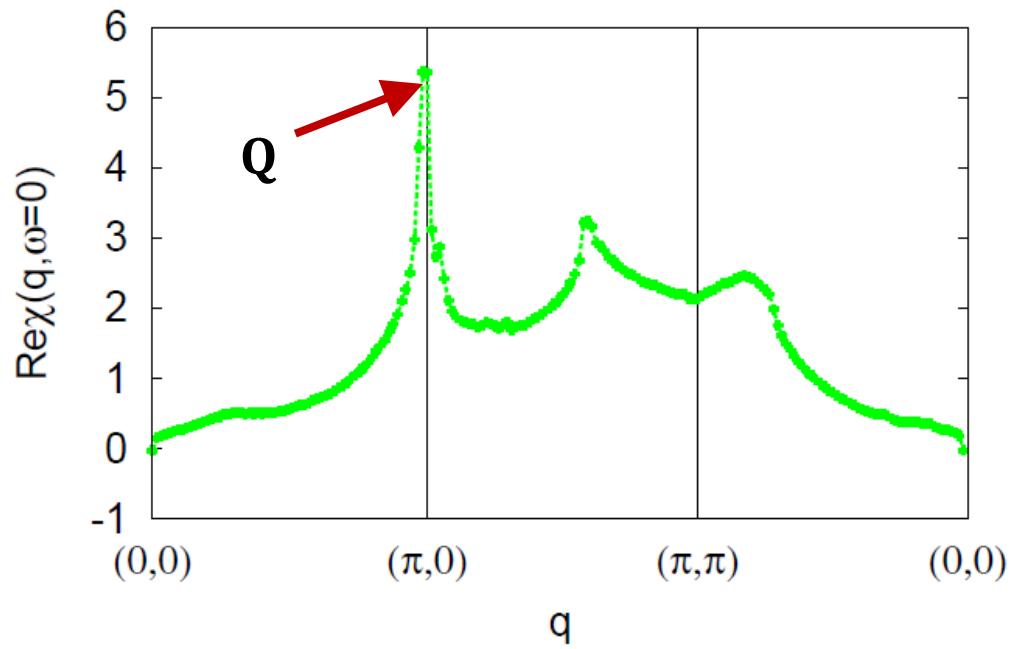
S. Lebegue, PRB 75, 035110 (2007); D.J. Singh and M.-H. Du, PRL 100, 237003 (2008); I.I. Mazin et al., PRL 101, 057003 (2008)

Spin susceptibility $\chi(q, \omega = 0)$

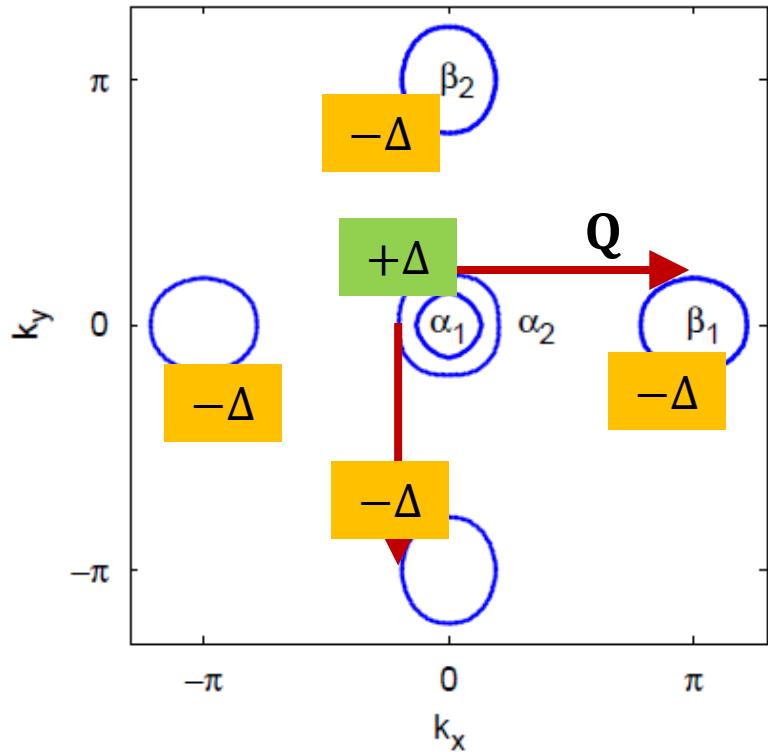


Local parameters (eV):
 $U=1.2$, $J=0.3$, $U'=0.6$, $J'=J$

Spin fluctuations – repulsive interaction at $\mathbf{Q} = (0, \pi)$
 (in the unfolded BZ)



Pnictides: nesting wave vector and the gap symmetry



$$\Delta_k(T) = - \sum_{k'} \frac{V(k - k')}{2E_{k'}} \Delta_{k'}(T) \tanh\left(\frac{E_{k'}}{2k_B T}\right)$$

$V(q = Q) > 0$ (repulsive interaction) –
 Δ have to change sign!

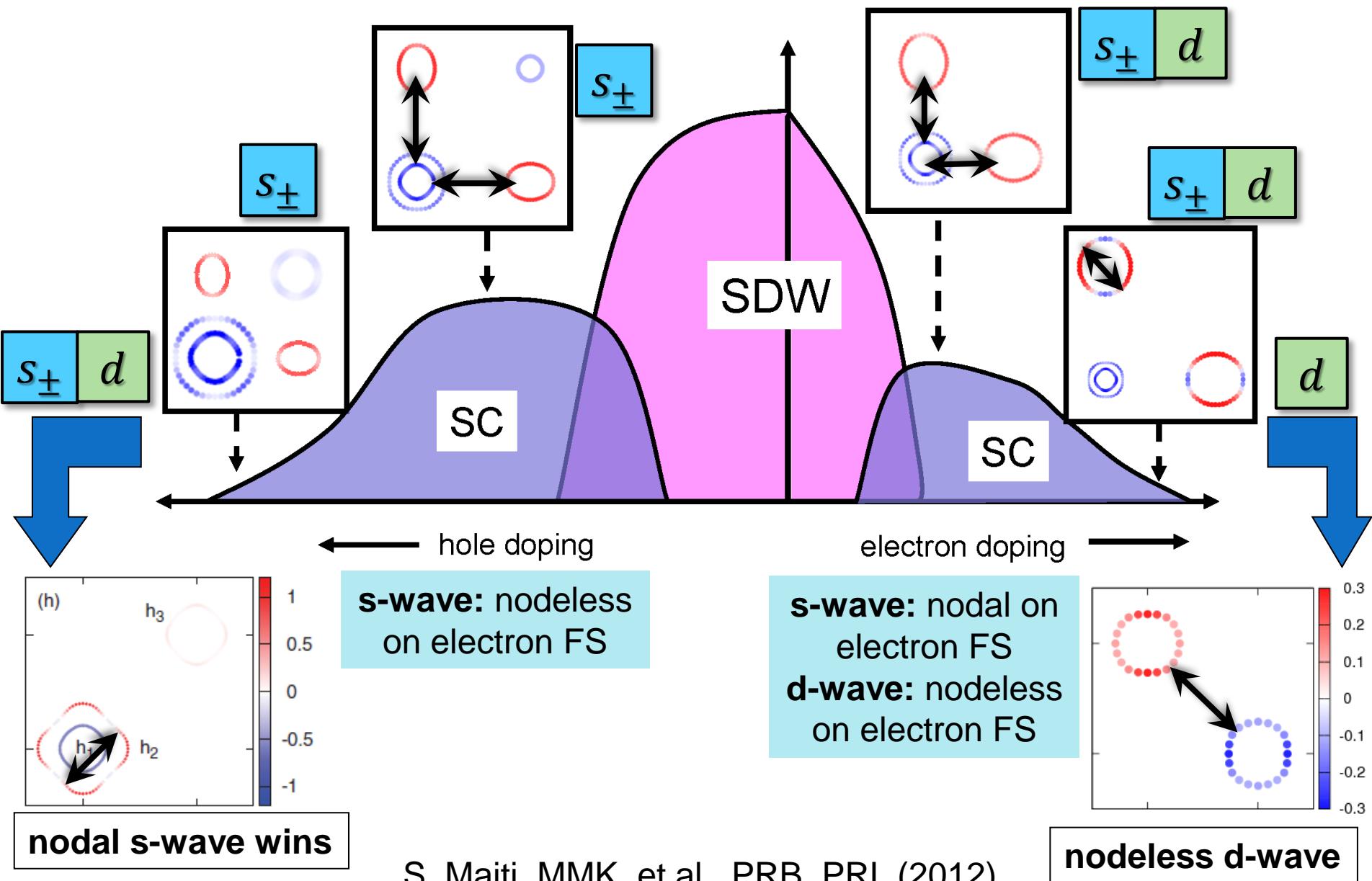
$$\Delta_1 = \Delta_2 e^{-i\pi}$$

Interband scattering with the SDW wave vector $\mathbf{Q}=(\pi,0)$ enhance intraband Cooper-pairing: extended s-wave (s_{\pm})

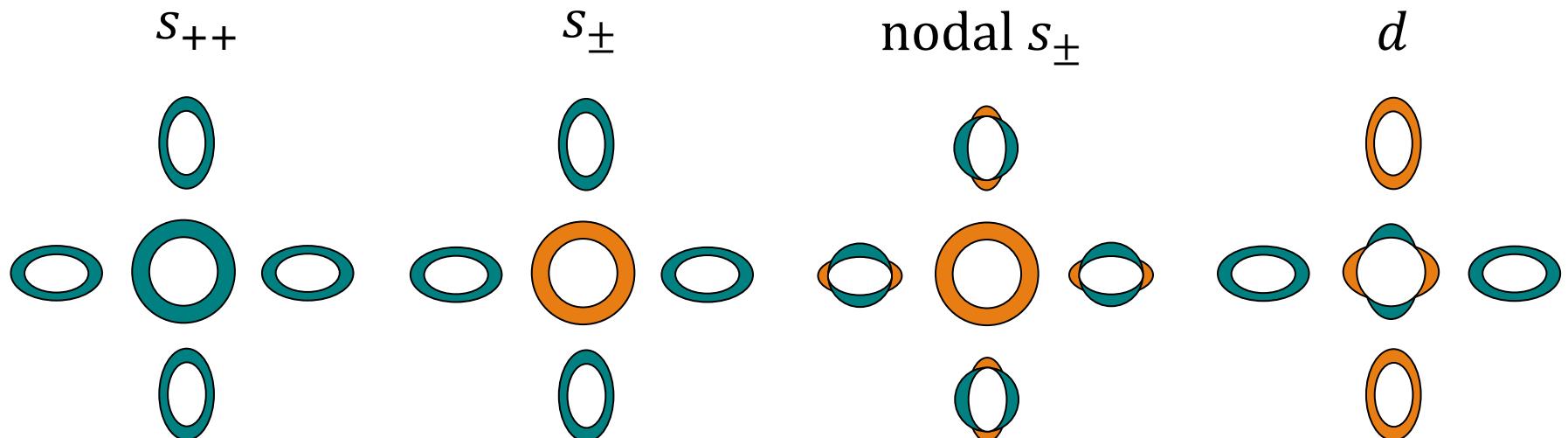
I.I. Mazin et al., PRL 101, 057003 (2008)

This symmetry most naturally comes out of the spin-fluctuation pairing theories

Spin-fluctuations pairing expectations in FeBS



Symmetry proposals for FeBS



Orbital fluctuations,
Phonons

Spin fluctuations

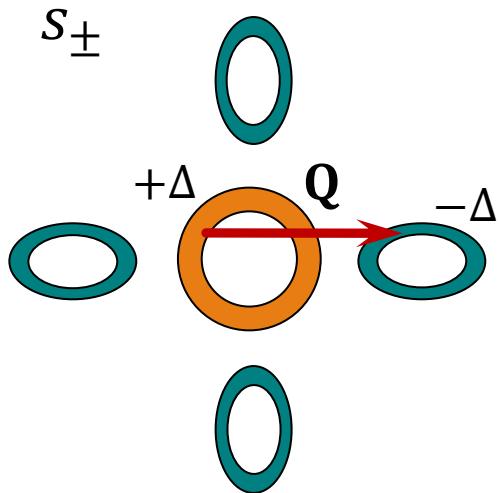
“Realistic” Spin
fluctuations

Mott-Hubbard-type
theories

Family	Full gap	Highly anisotropic	Strong nodal
1111	PrFeAsO _{1-y} [52K] [290] SmFeAs(O, F)[55K] [211]	LaFeAs(O, F)[26K] [212] NdFeAs(O, F) [212]	LaFePO[6K] [201, 202, 291]
122	(Ba, K)Fe ₂ As ₂ [40K] [144, 234, 240, 292] Ba(Fe, Co) ₂ As ₂ [OP, 23K] [236, 206]	Ba(Fe, Co) ₂ As ₂ [OD] [236, 239]* Ba(Fe, Ni) ₂ As ₂ [293]* Ba(Fe, Co) ₂ As ₂ [UD] [239]*	KFe ₂ As ₂ [4K] [209, 304] BaFe ₂ (As, P) ₂ [OP, 31K] [203, 147] (Ba, K)Fe ₂ As ₂ [UD] [240] LiFeP[6K] [295]
111	LiFeAs[18K] [294, 256]		
11		Fe(Se, Te)[27K] [229, 244]	

Note: Symbol * indicates possible evidence for ‘c-axis nodes’.

Resonance peak in INS: hallmark of s_{\pm} symmetry

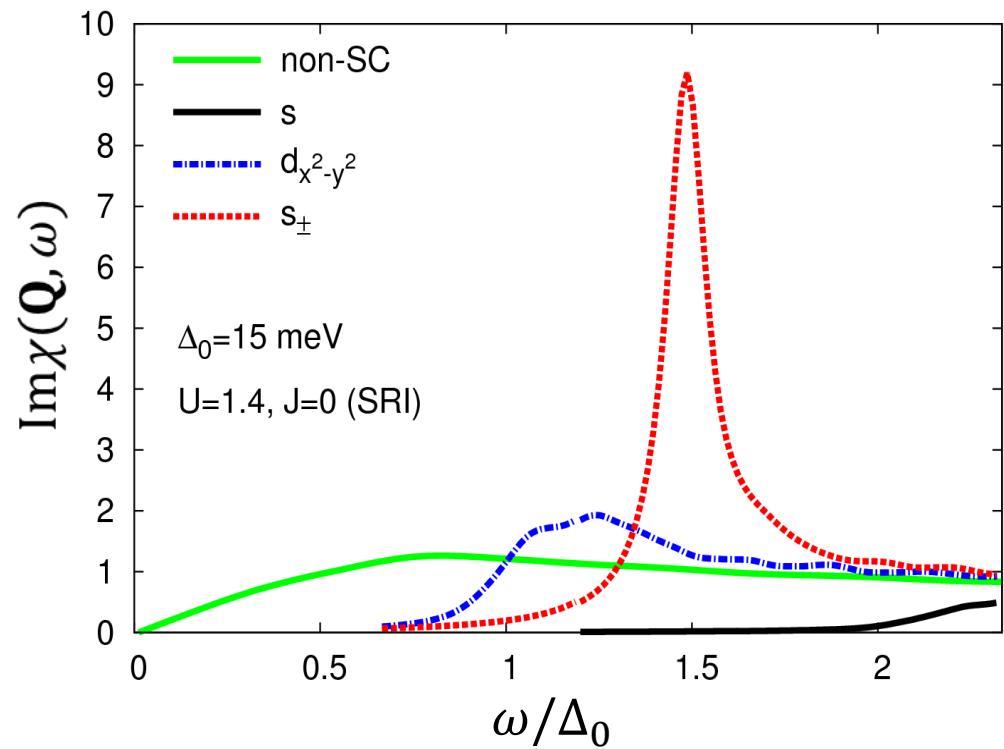


s_{\pm} superconductivity →
resonance peak in INS,
 $\omega_R \leq 2\Delta$

Observed experimentally:
A.D. Christianson et al., Nature 456, 930 (2008); M.D. Lumsden et al., PRL 102, 107005 (2009)

For s_{\pm} symmetry $\Delta_{\mathbf{k}}^{\alpha} = -\Delta_{\mathbf{k}+\mathbf{Q}}^{\beta}$

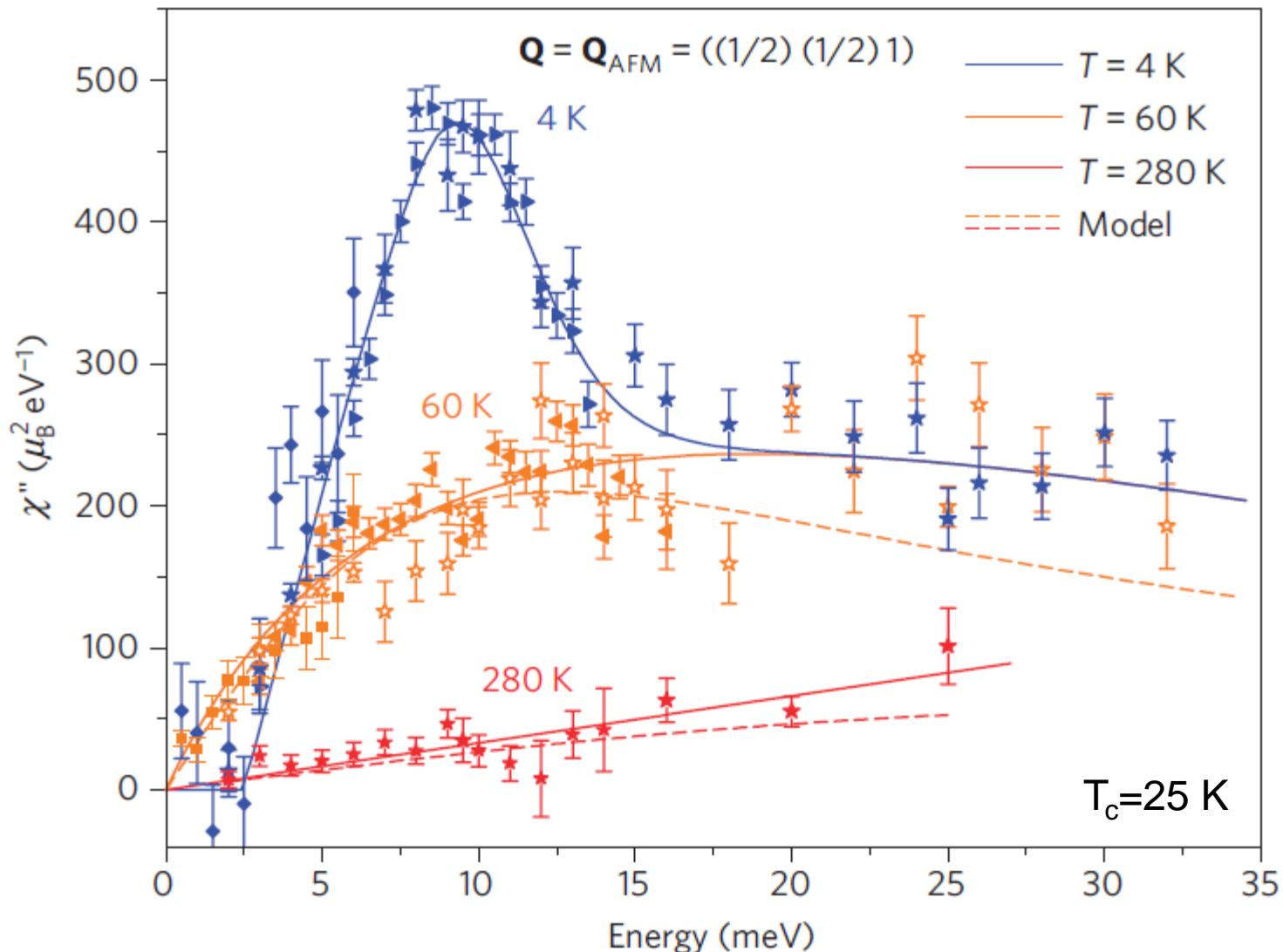
Coherence factors in $\text{Im}\chi$: $\left(1 - \frac{\Delta_{\mathbf{k}}\Delta_{\mathbf{k}+\mathbf{Q}}}{|\Delta_{\mathbf{k}}||\Delta_{\mathbf{k}+\mathbf{Q}}|}\right) \neq 0$



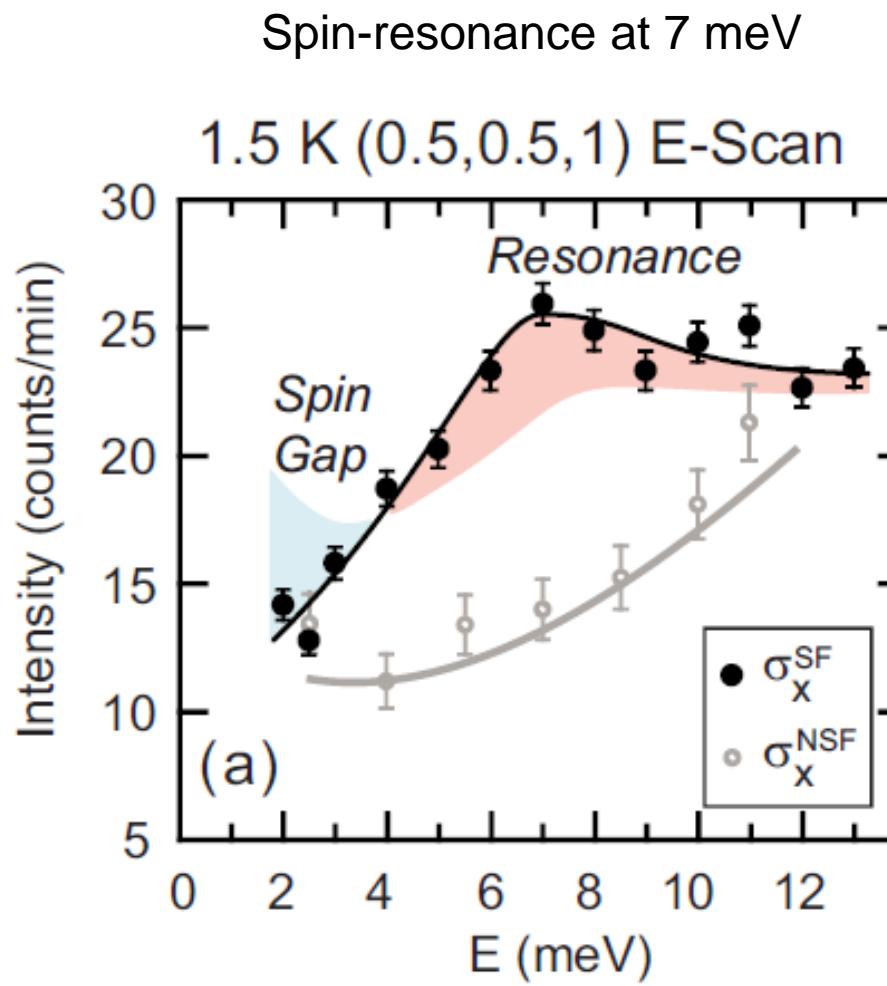
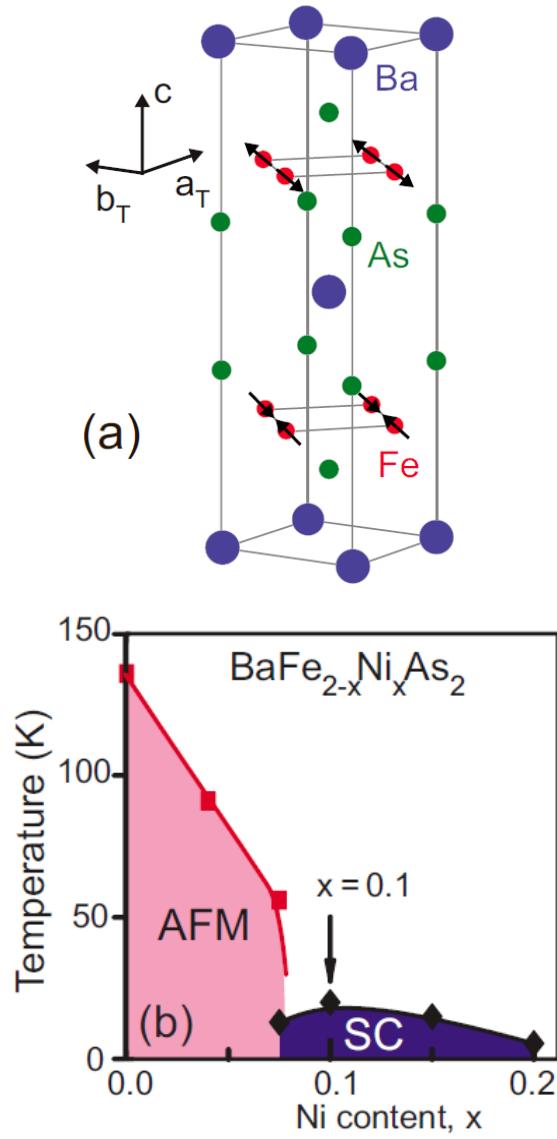
MMK and I. Eremin, PRB 78, 140509(R) (2008)

T.A. Maier and D.J. Scalapino, PRB 78, 020514(R) (2008)

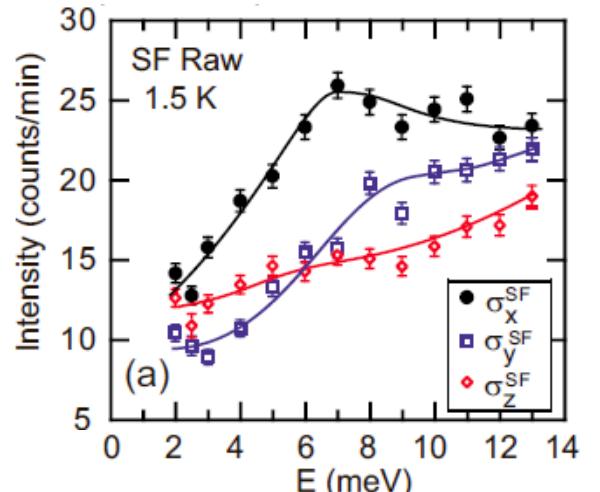
Example of inelastic neutron scattering data on $\text{BaFe}_{1.85}\text{Co}_{0.15}\text{As}_2$



Polarized neutron scattering on $\text{BaFe}_{2-x}\text{Ni}_x\text{As}_2$



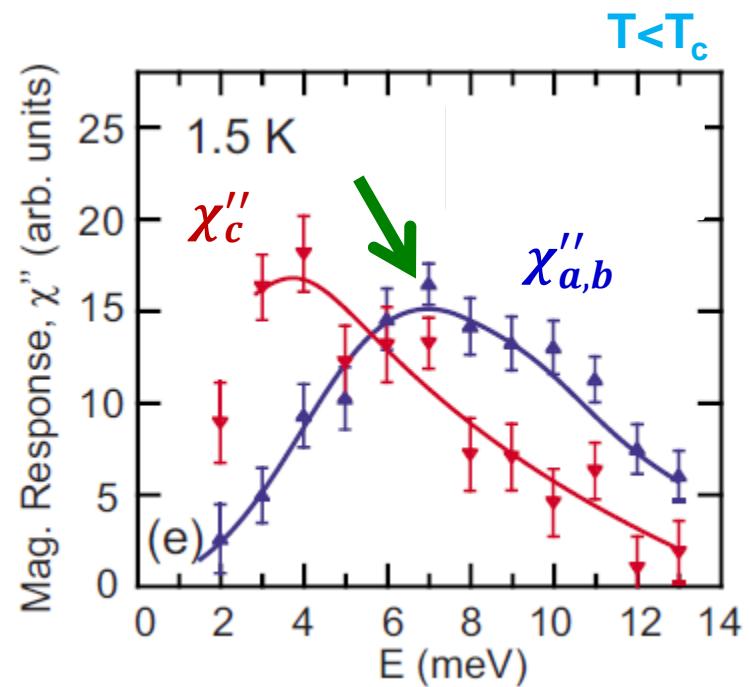
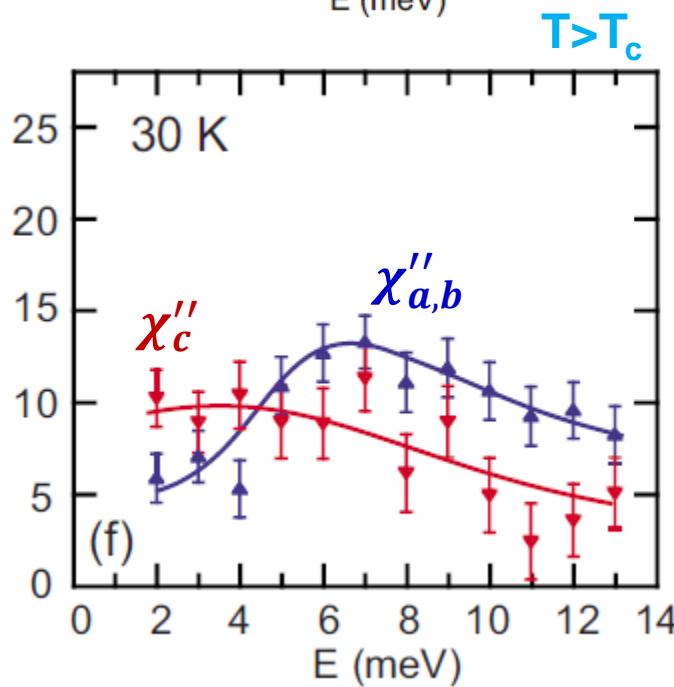
Polarized neutron scattering: spin-resonance anisotropy



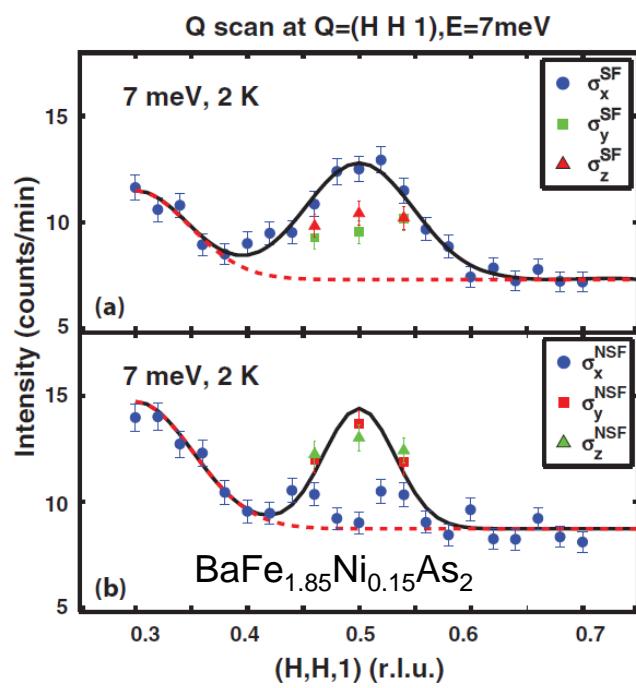
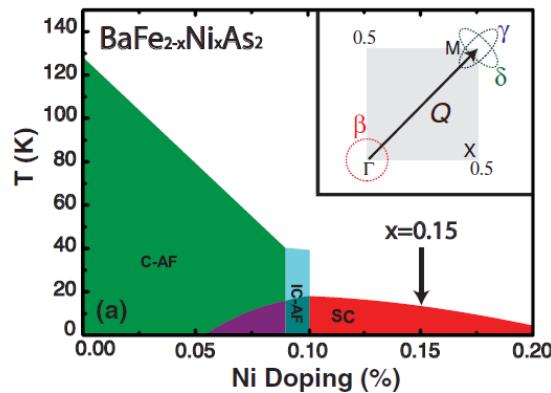
$$\langle S^x S^x \rangle = \langle S^y S^y \rangle = \langle S^z S^z \rangle$$

$$\chi_{+-} = \chi_{xx} + \chi_{yy} = 2\chi_{zz}$$

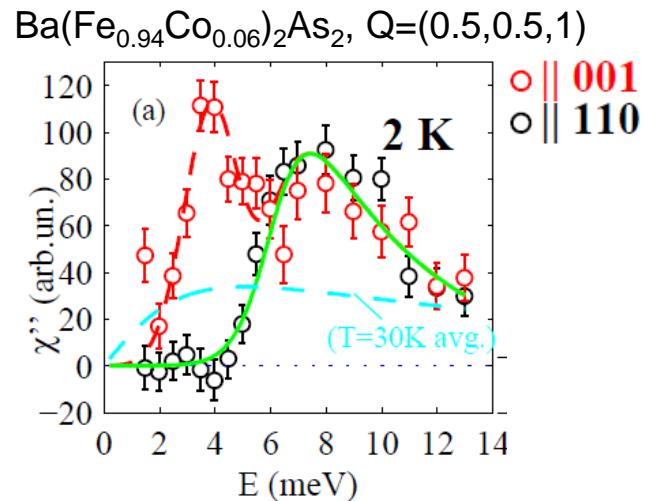
$\chi''_{a,b} \neq \chi''_c$: Spin-rotational invariance is broken!



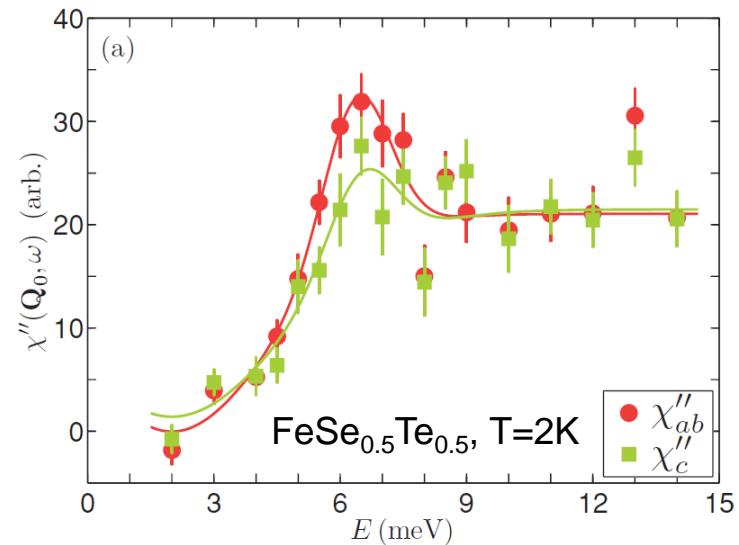
Spin-resonance anisotropy: other studies



M. Liu et al., PRB 85, 214516 (2012)



P. Steffens et al., arXiv:1210.6386v1



P. Babkevich et al., PRB 83, 180506(R) (2011)

NMR above T_c : $1/T_1$ anisotropy in $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$

$x=0.32, T_c = 38.5 \text{ K}$ and $x=0.28, T_c = 31.5 \text{ K}$

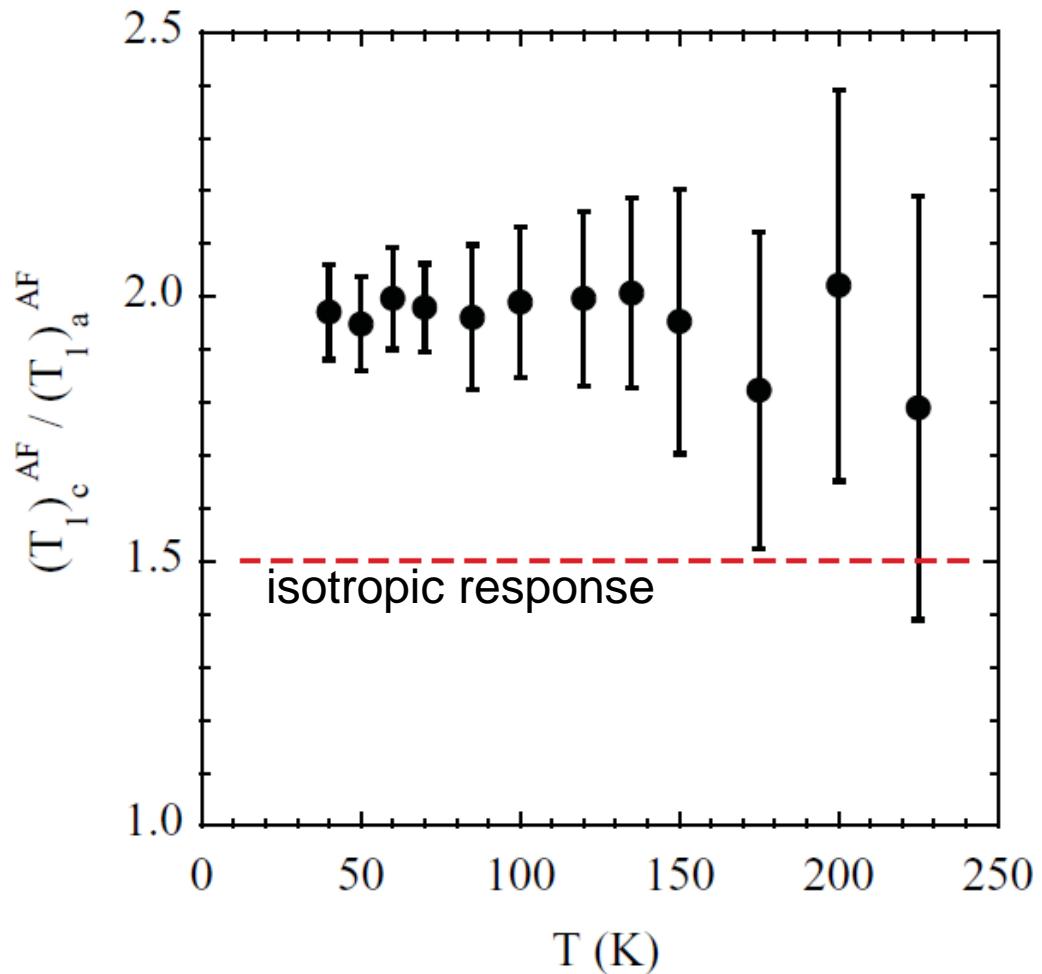
$^{75}\text{As-NMR}$

$$R = \frac{(1/T_1)_a}{(1/T_1)_c} = \frac{\chi_a'' + \chi_b''}{2\chi_c''} + \frac{1}{2}$$

$$R^{\text{iso}} = 1.5$$

$$\frac{1}{2}\chi_{+-} > \chi_{zz}$$

by about 50%



K. Matano et al., EPL 87, 27012 (2009)

Z. Li et al., PRB 83, 140506(R) (2011)

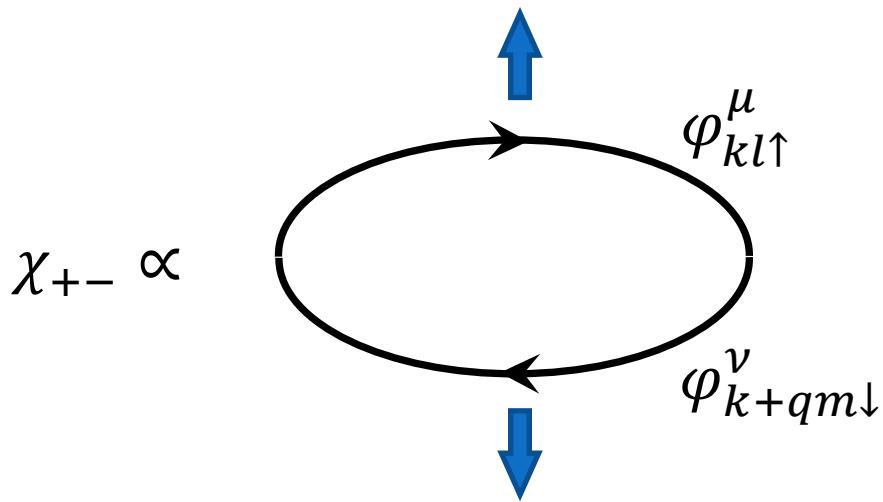
What can cause violation of spin-rotational invariance?



Spin-orbit coupling:

$$H_{SO} = \lambda \sum_f \mathbf{L}_f \cdot \mathbf{S}_f$$

Spin-orbit coupling affects the spin susceptibility

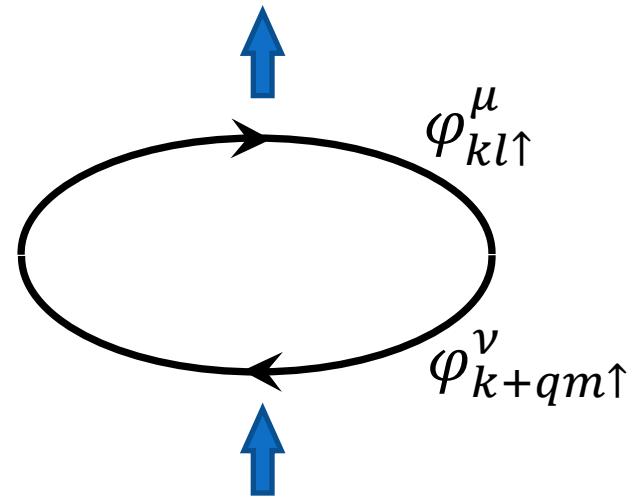


Matrix elements (orbitals l – bands μ):

$$d_{kl\sigma} = \sum_\mu \varphi_{kl\sigma}^\mu b_{k\mu\sigma}$$

$$\hat{H}_{SO} = \frac{\lambda}{2} \begin{pmatrix} \widehat{L_z S_z} & \widehat{L_- S_+} \\ \widehat{L_+ S_-} & -\widehat{L_z S_z} \end{pmatrix}$$

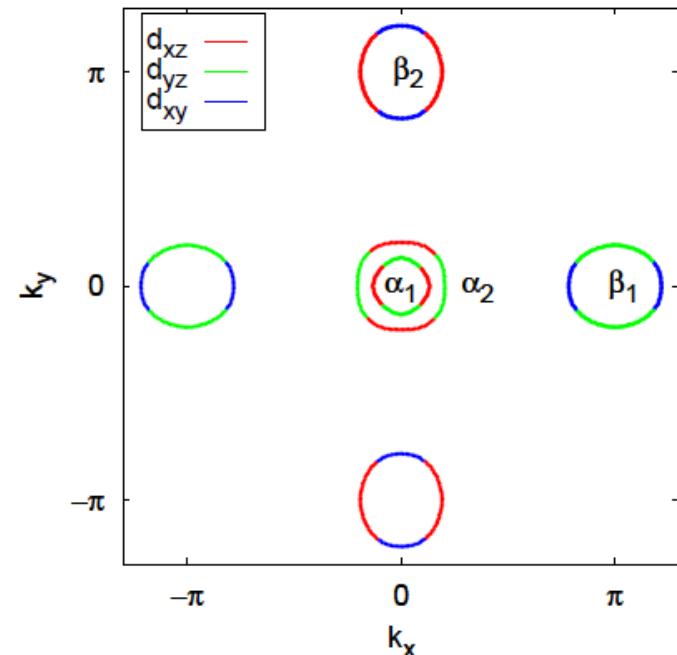
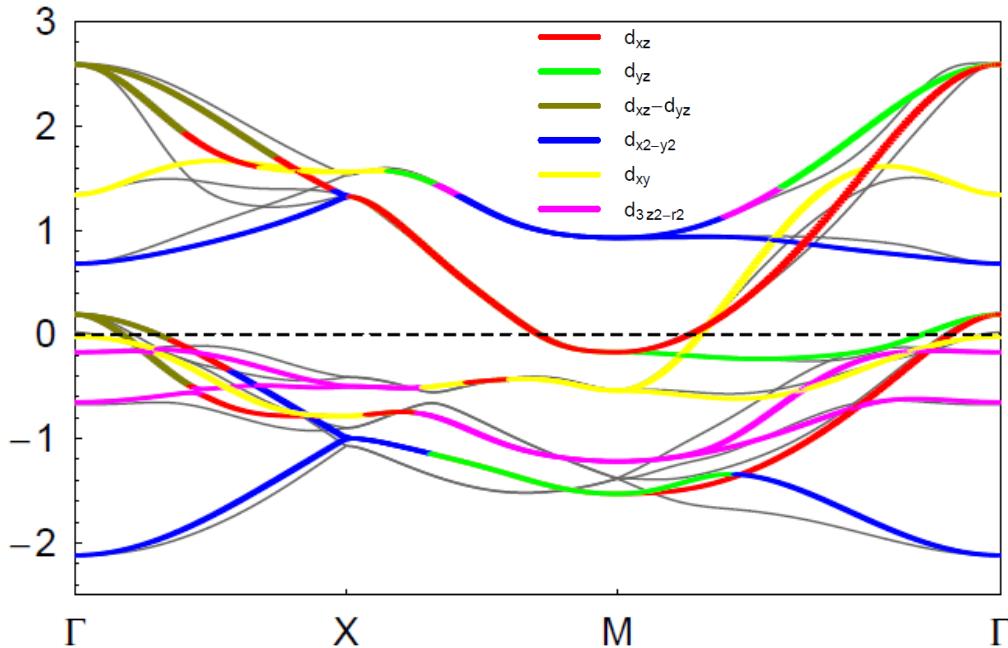
$\chi_{zz} \propto$



Hamiltonian and interactions

$$H = H_0 + H_{int}$$

H_0 = 5-orbital tight-binding model

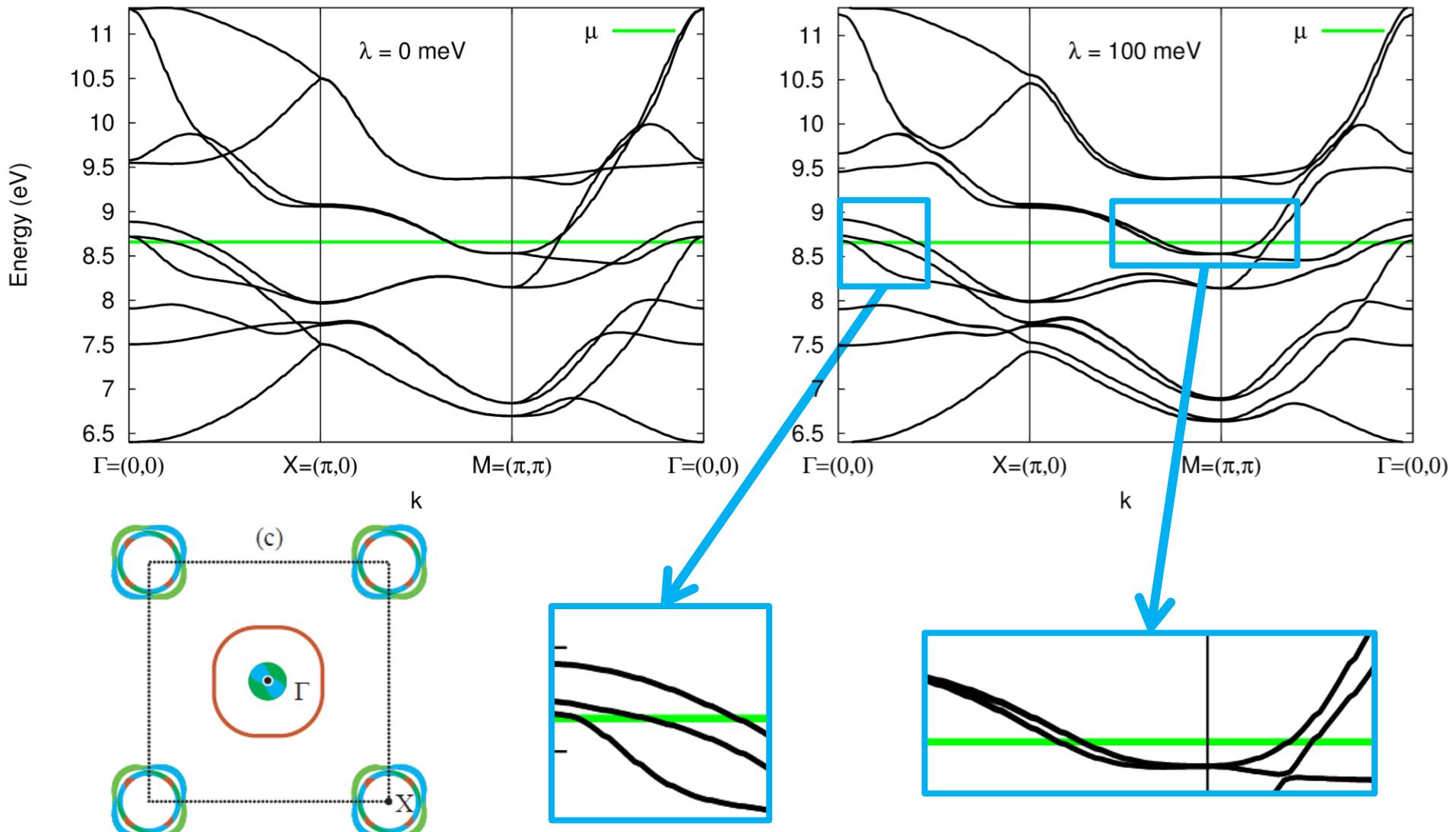


Matrix elements (orbitals l – bands μ): $d_{kl\sigma} = \sum_{\mu} \varphi_{kl\sigma}^{\mu} b_{k\mu\sigma}$

Most general 2-body Hamiltonian with [introsite](#) interactions only

$$H_{int} = U \sum_{i,m} n_{im\uparrow} n_{im\downarrow} + J \sum_{i,m < n} \sum_{\sigma, \sigma'} d_{in\sigma}^{\dagger} d_{im\sigma'}^{\dagger} d_{in\sigma'} d_{im\sigma} + U' \sum_{i,m < n} n_{in} n_{im} + J' \sum_{i,m \neq n} d_{in\uparrow}^{\dagger} d_{in\downarrow}^{\dagger} d_{im\downarrow} d_{im\uparrow}$$

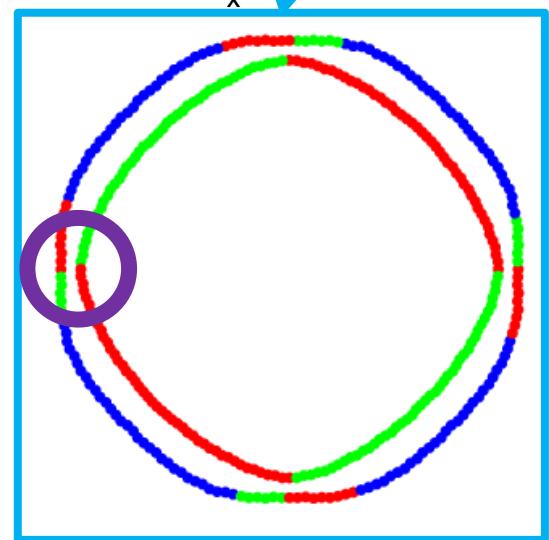
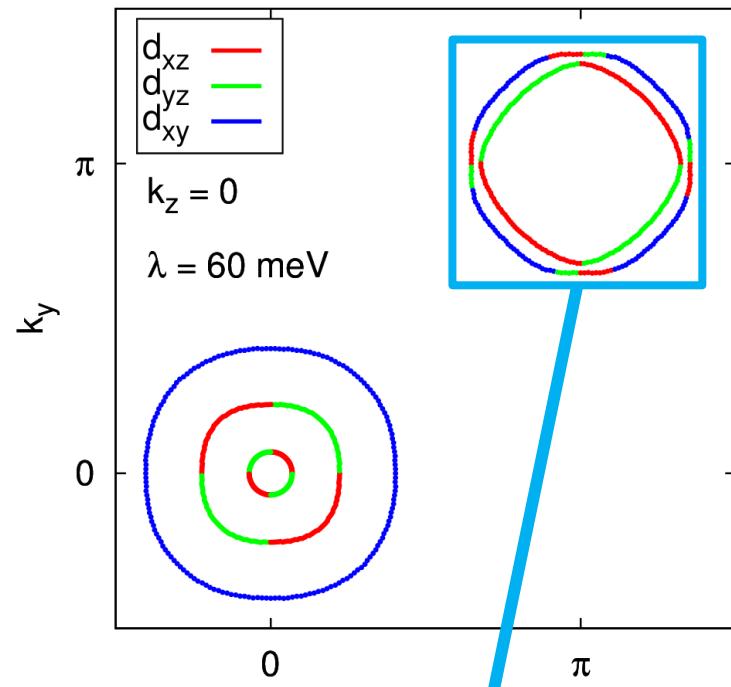
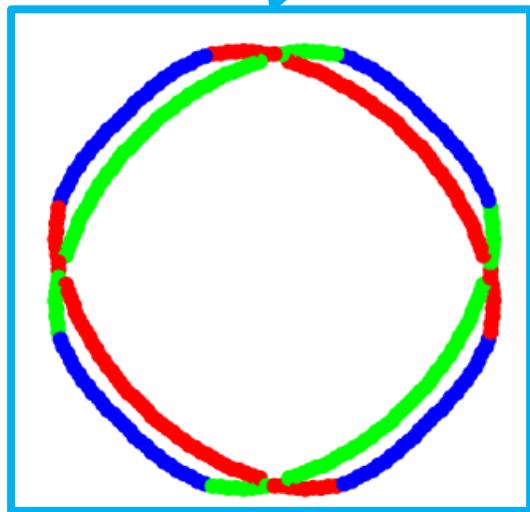
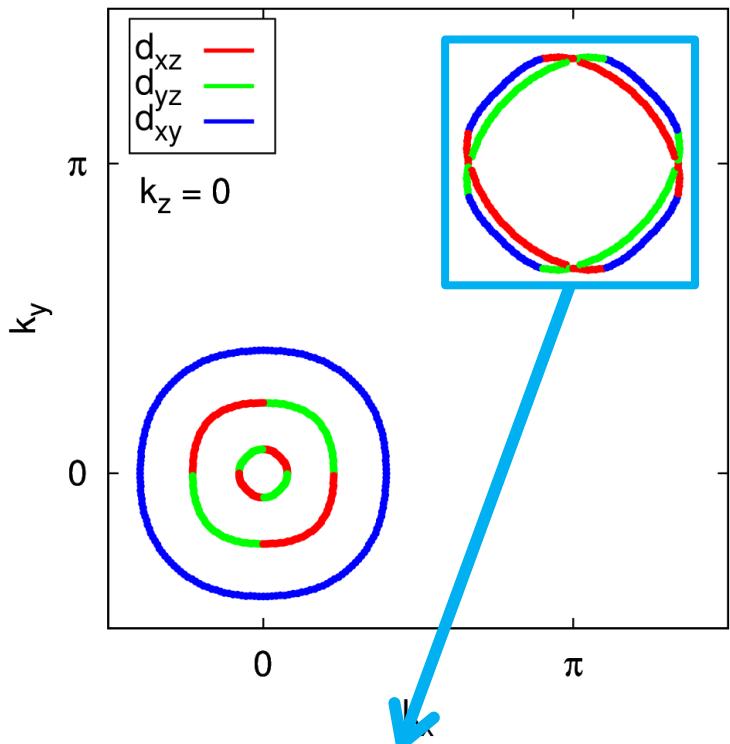
Band structure changes due to the SO coupling: 10-orbital model for LiFeAs



From A.A. Kordyuk review,
Low Temp. Phys. (2012)

SO coupling – splitting of many band crossings

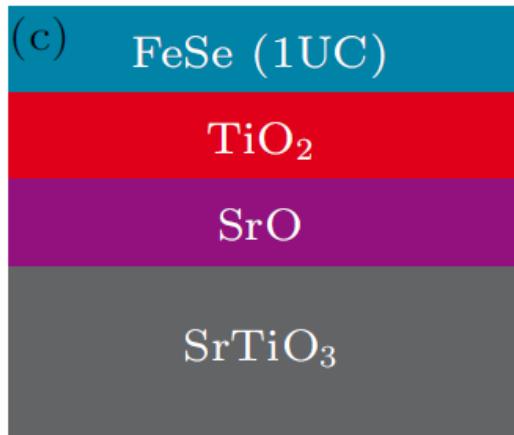
Fermi surface topology changes due to the SO coupling



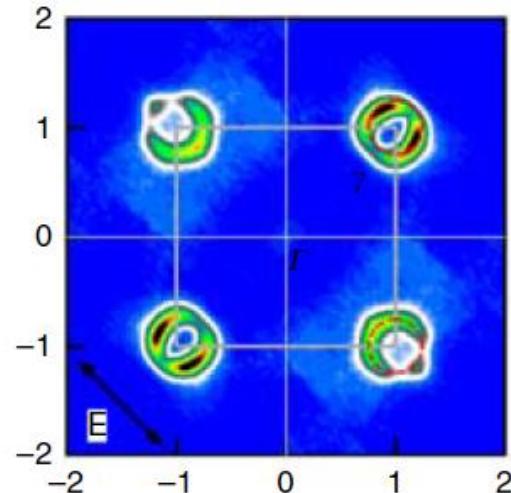
SO coupling “splits”
the Fermi surfaces

MMK et al., unpubl

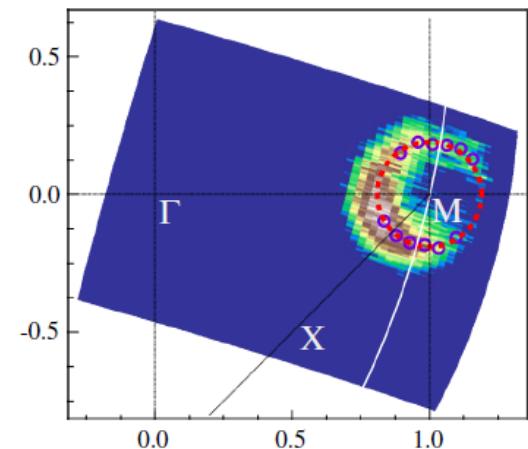
A Fe_2Se_2 and FeSe monolayers



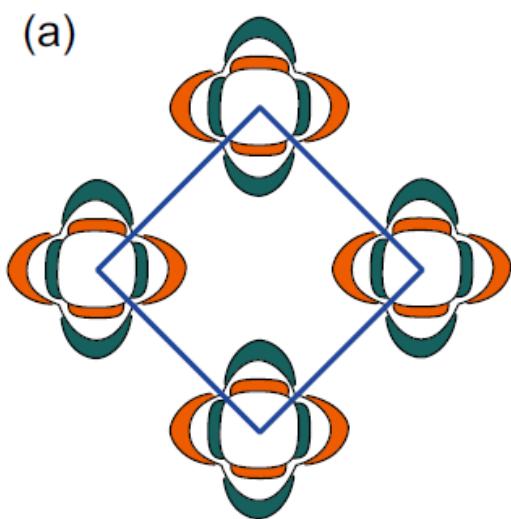
Wang et al, Chin. Phys. Lett.
29, 037402 (2012)



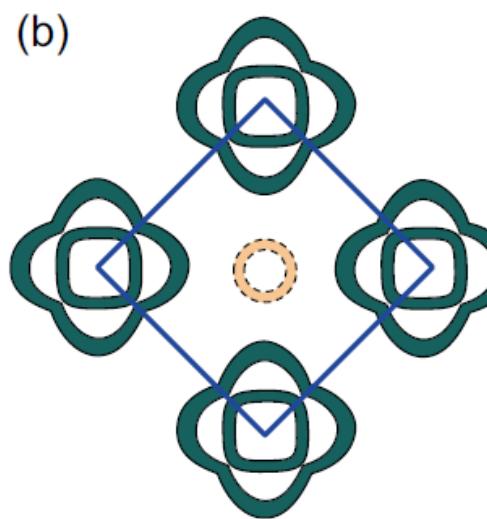
FeSe: Liu et al, Nat.
Comm. 3, 931 (2012)



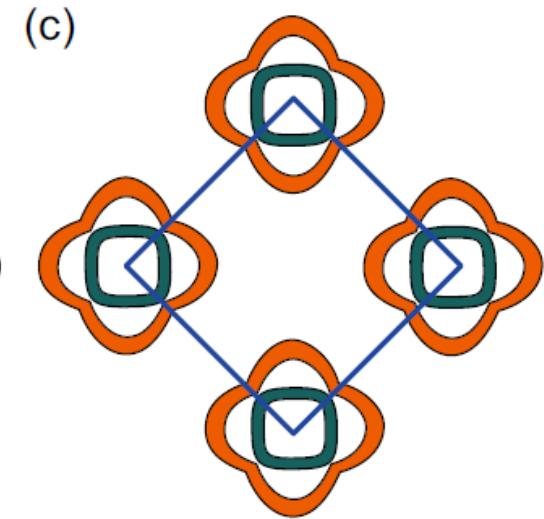
$\text{K}_{0.8}\text{Fe}_{1.7}\text{Se}_2$: Qian et al.,
PRL 106, 187001 (2011)



“quasi-nodeless” d



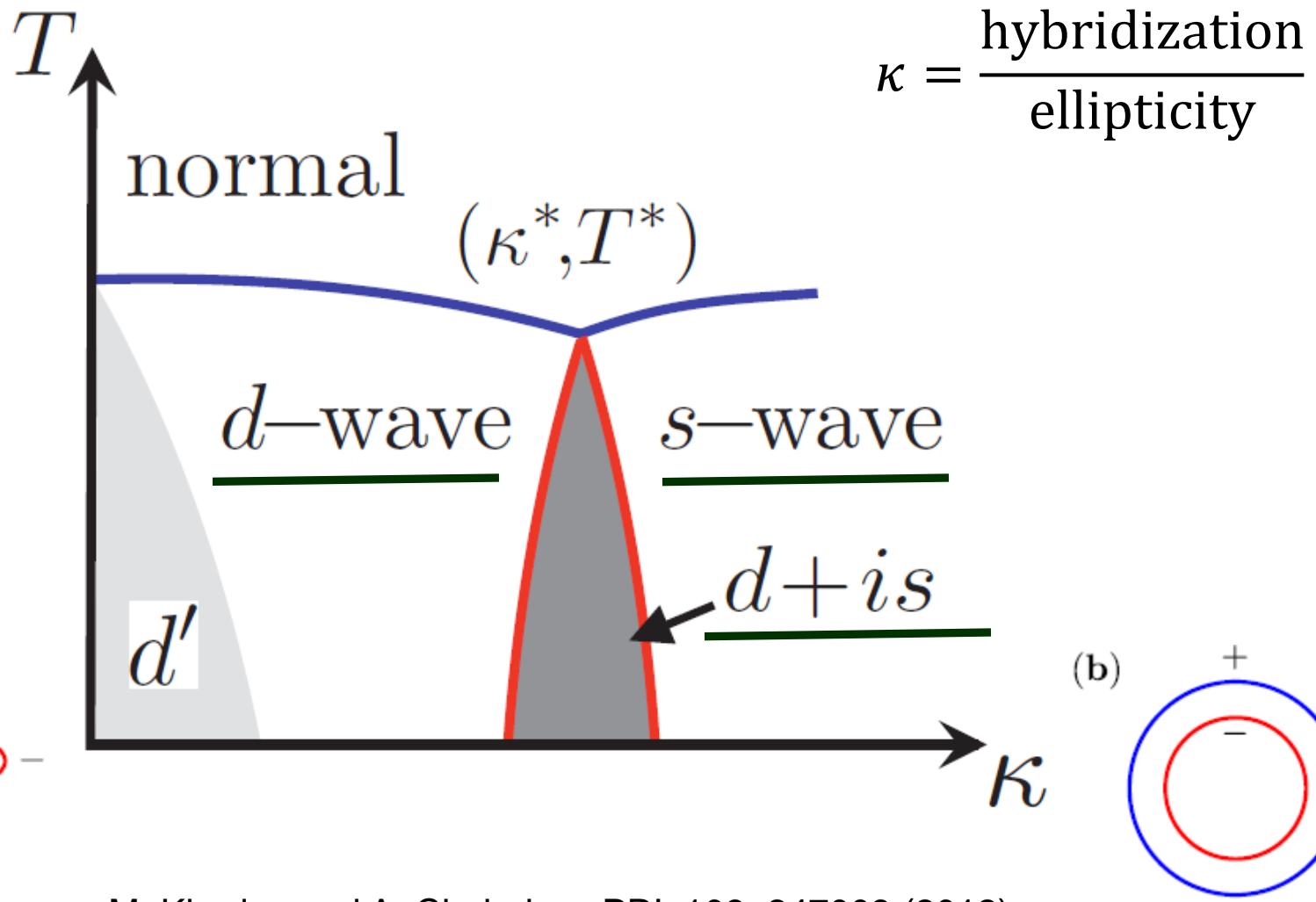
“incipient” s_{\pm}



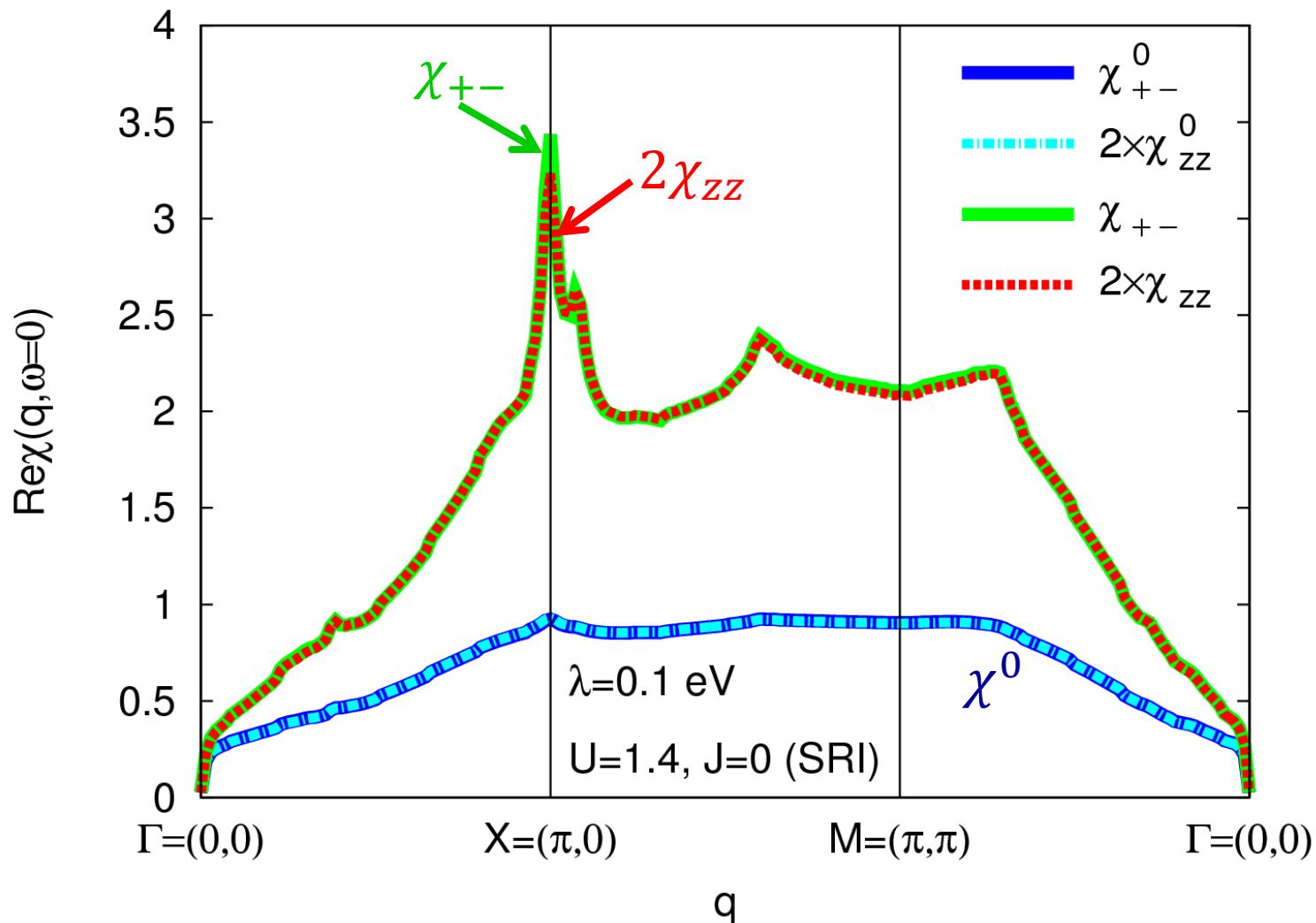
“bonding – antibonding” s_{\pm}

Role of hybridization and ellipticity (AFe_2Se_2)

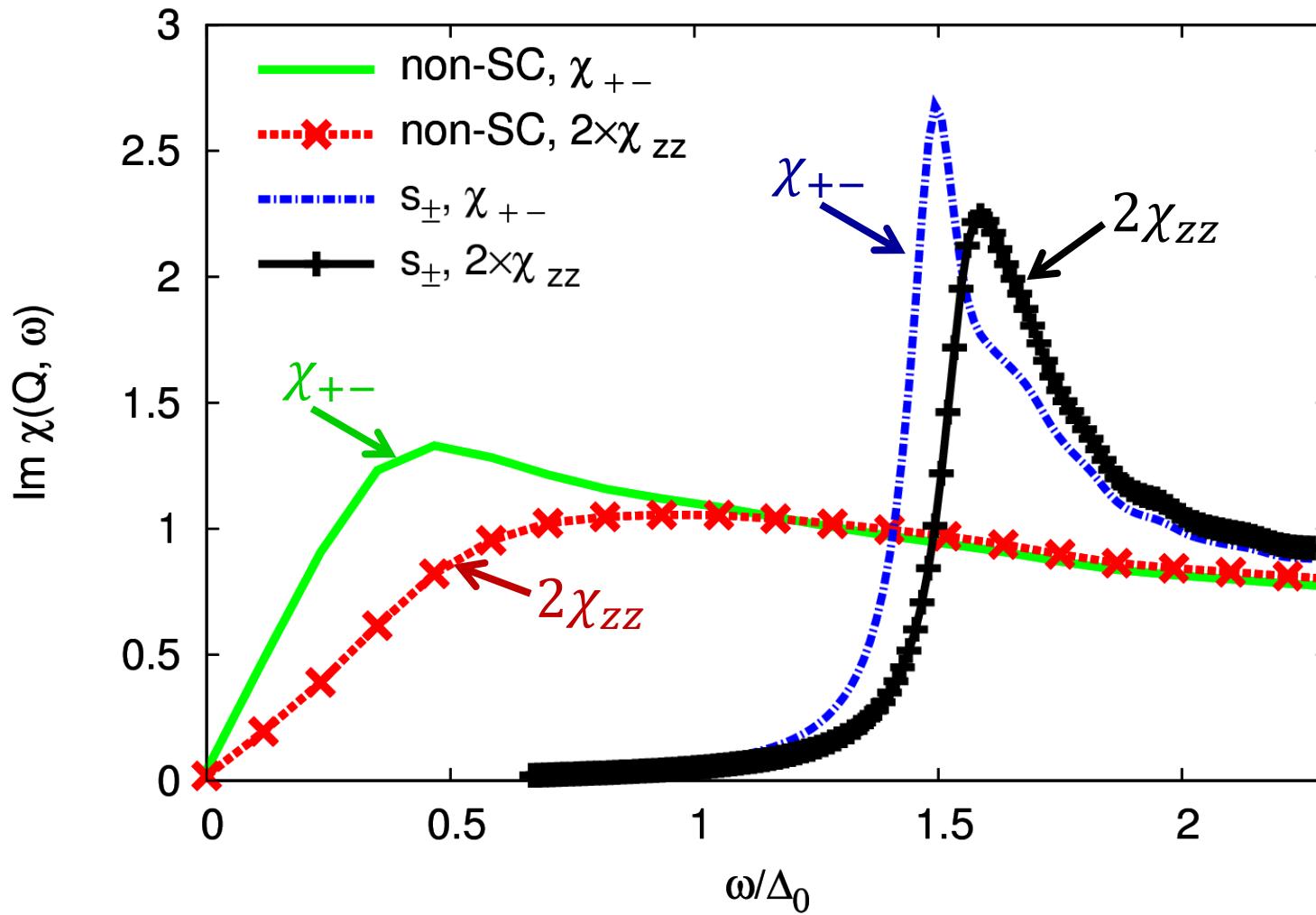
Two parameters: hybridization and the ellipticity of electron pockets



Spin susceptibility at $\omega = 0$ in the 5-orbital model with the SO



Spin-resonance at $\mathbf{q} = \mathbf{Q}$ in the 5-orbital model with the SO

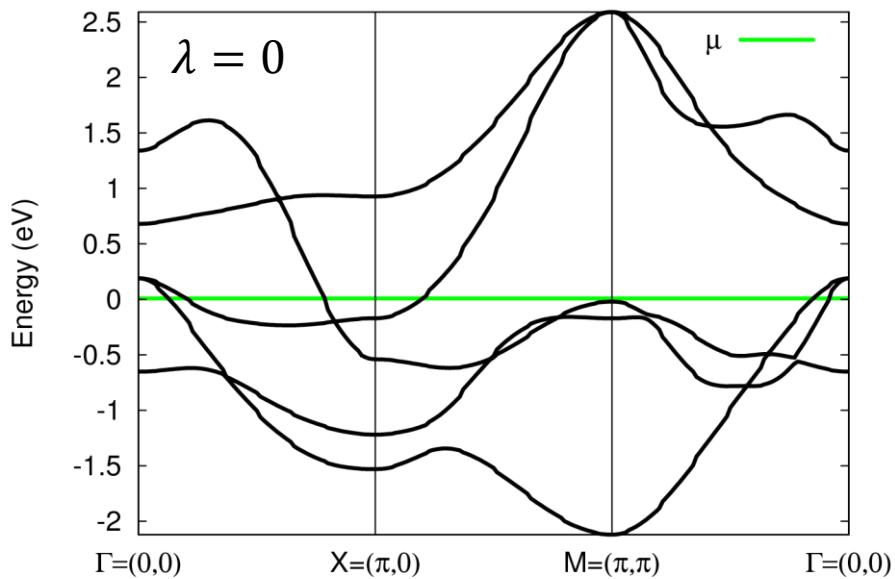


Spin-orbit coupling provides both asymmetry and frequency shift

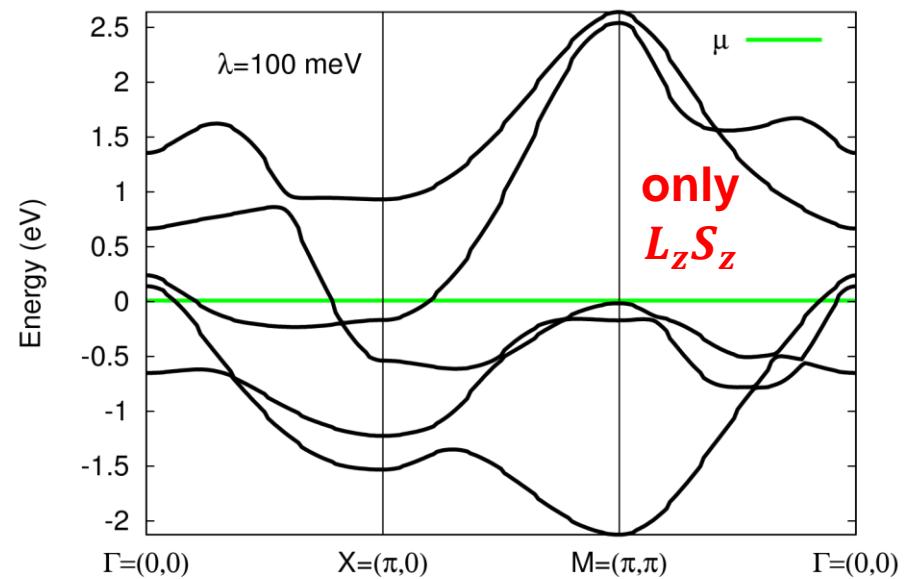
Conclusions

- **Spin-orbit (SO) coupling** provides momentum-dependent disparity between $+ -$ and zz components of the spin susceptibility. This disparity is maximal at SDW wave vector and present in both normal and SC states.
- **$+ -$ wins over zz component in the normal state** that is consistent with the observation of in-plane orientations of spins in the magnetic SDW state.
- **SO coupling gives qualitative understanding of the spin-space asymmetry** in neutron scattering and NMR data.
- **Spin-resonance for s_{\pm} gap symmetry is partially suppressed in the out-of-plane (zz) component** as compared to the in-plane one.
- **Band structure and Fermi surface are both affected by the SO coupling.**

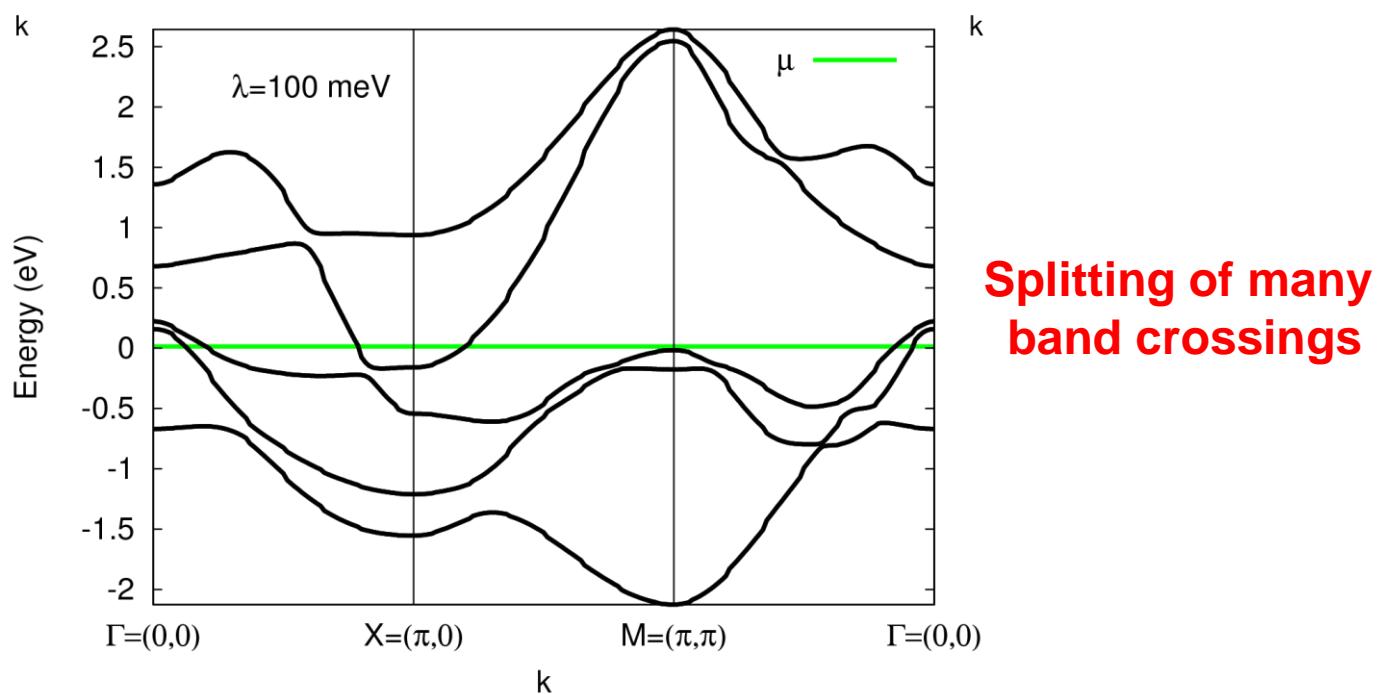
Band structure changes with the finite SO coupling



$\lambda = 0$



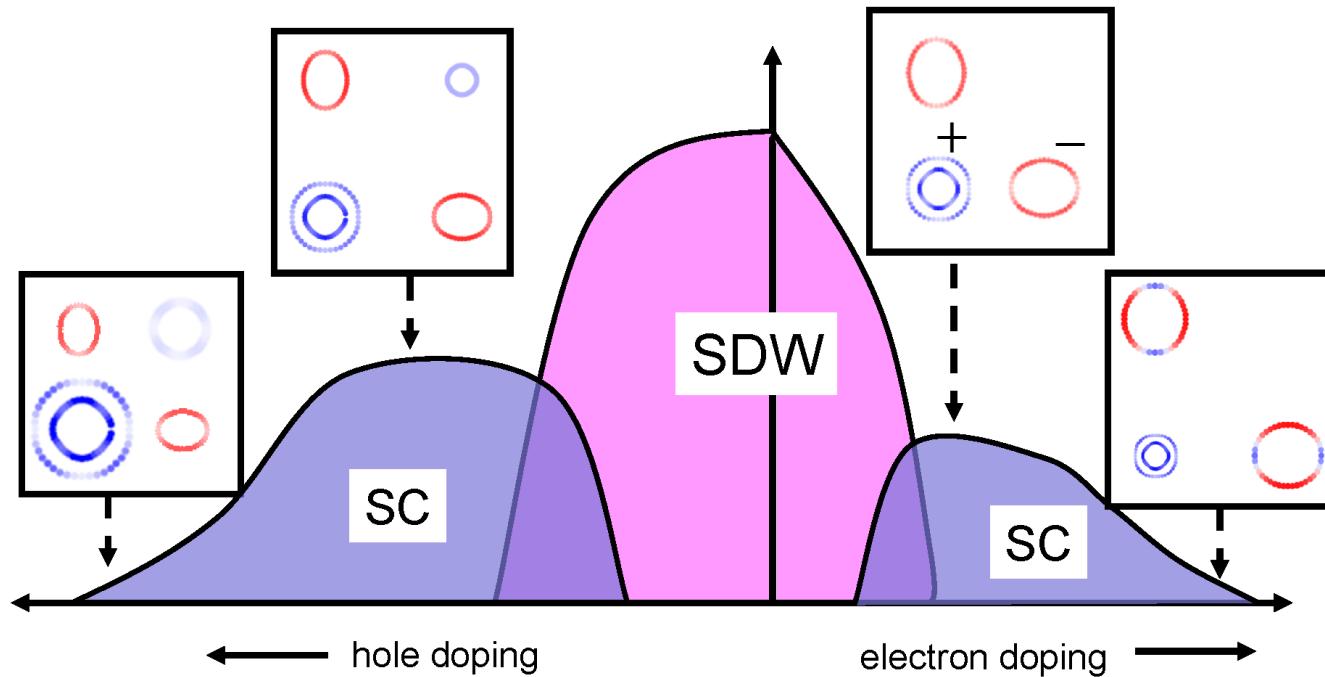
$\lambda=100$ meV
only
 $L_z S_z$



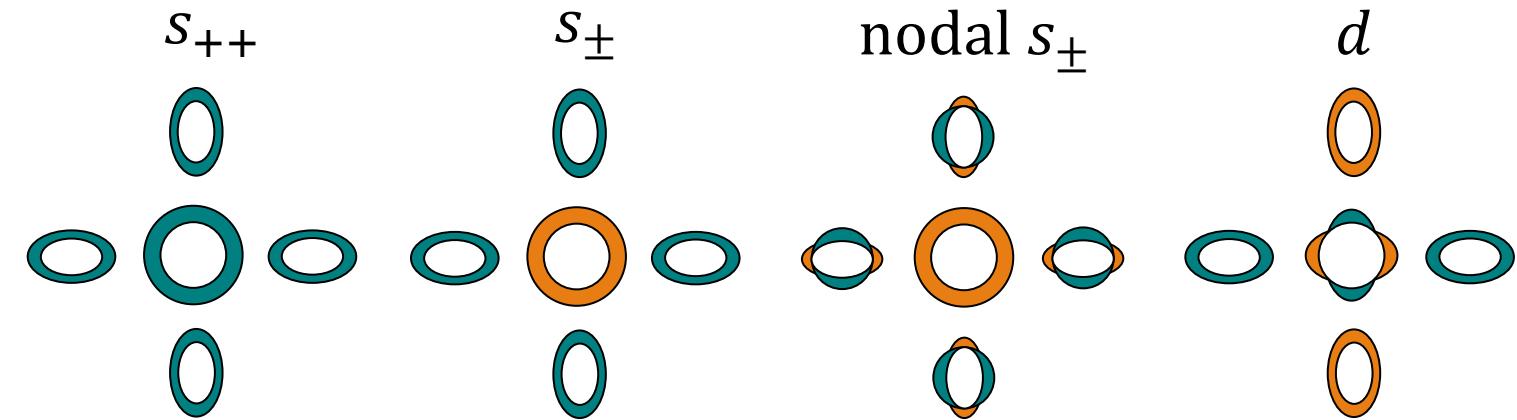
Full SO coupling

Splitting of many band crossings

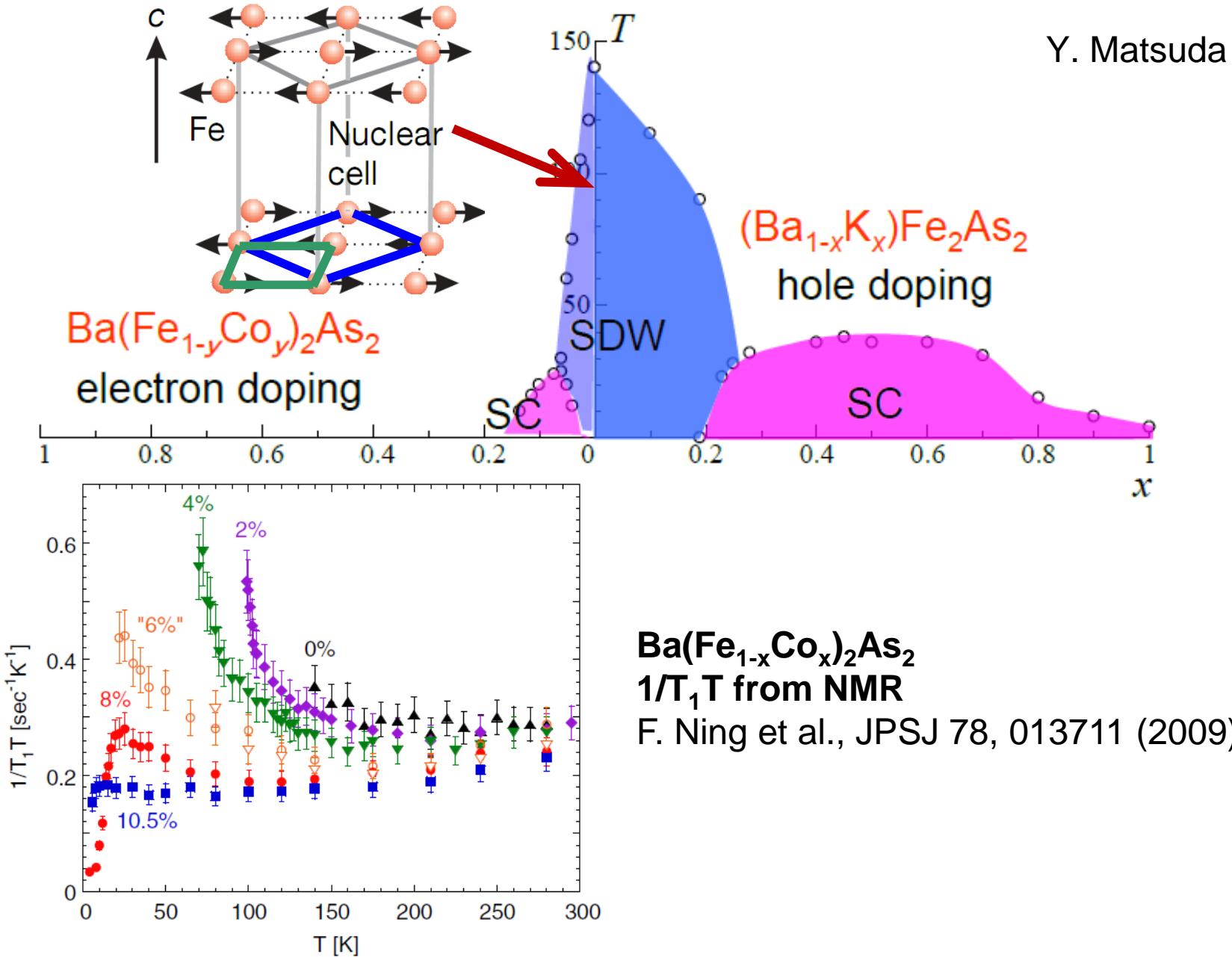
2D spin-fluctuation pairing expectations in FeBS



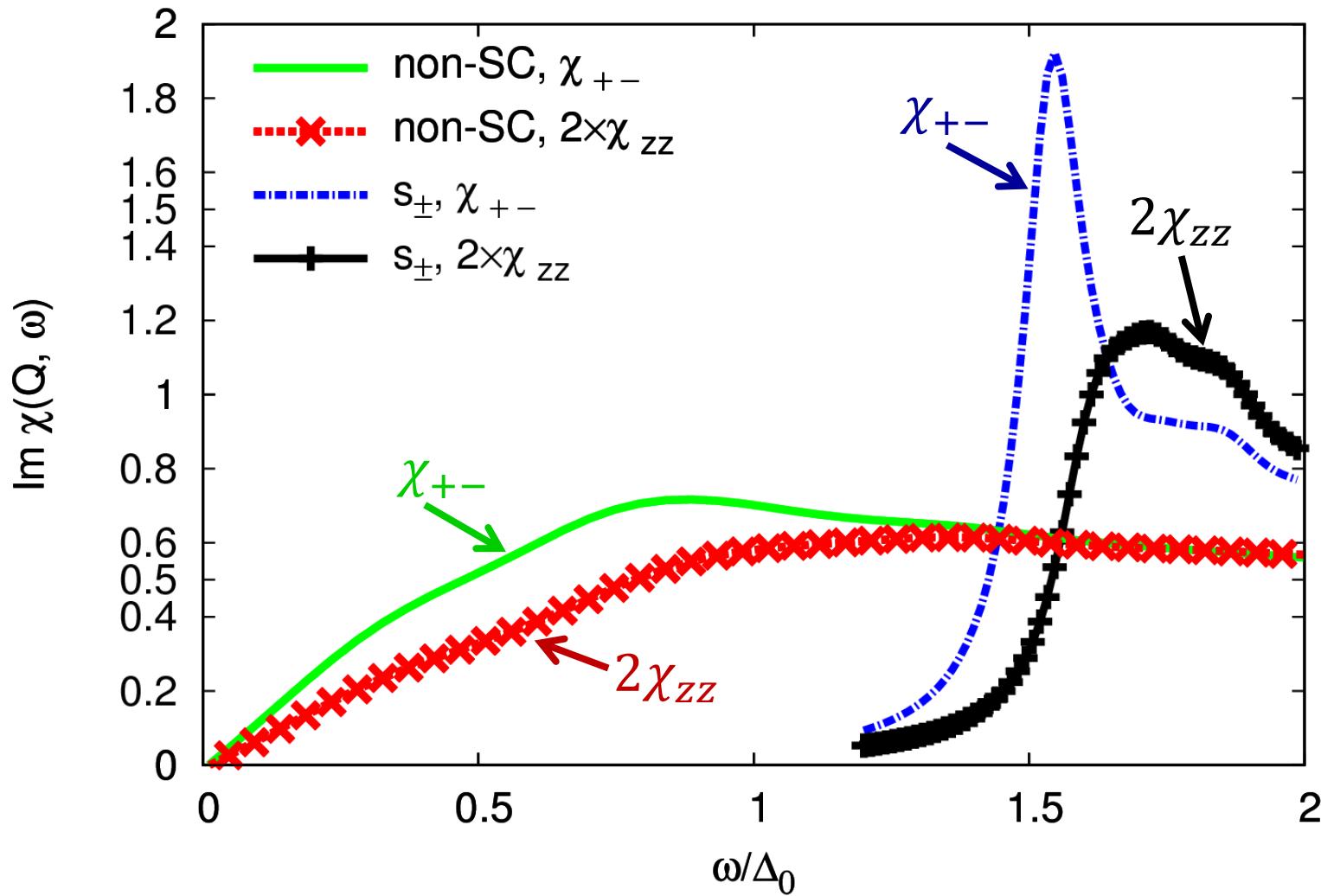
Other symmetry
proposals for FeBS



Phase diagram of Ba-122 system



Spin-resonance at $\mathbf{q} = \mathbf{Q}$ in the 5-orbital model: only $L_z S_z$ -component of SO



Spin-resonance anisotropy if further enhanced

$\lambda = 100 \text{ meV}, U = 1.4, J = 0$ (SRI), $\Delta_0 = 15 \text{ meV}$