

The electronic structure within the norm-conserving cluster perturbation theory

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CDMFT

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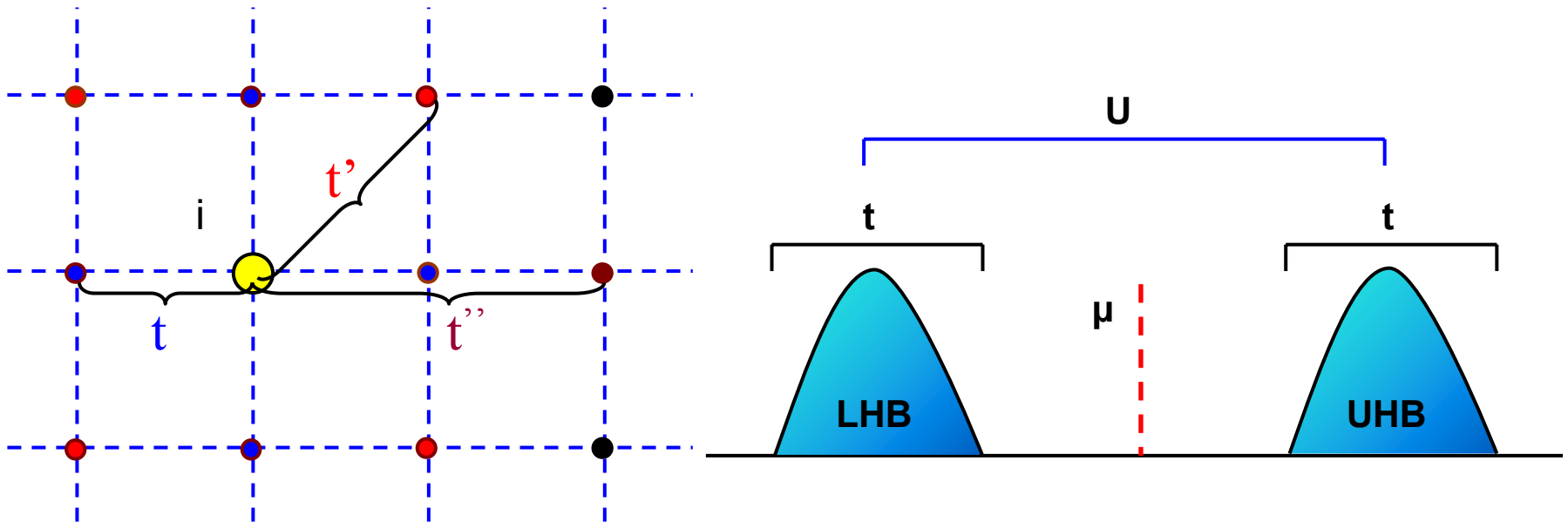
V-CPT

The one-band two-dimensional Hubbard model

$$H = H_0 + H_1, \quad (1)$$

$$H_0 = \sum_{i\sigma} \left\{ (\varepsilon - \mu) n_{i\sigma} + \frac{U}{2} n_{i\sigma} n_{i\bar{\sigma}} \right\}, \quad (2)$$

$$H_1 = \sum_{i \neq j, \sigma} t_{ij} a_{i\sigma}^+ a_{j\sigma}. \quad (3)$$

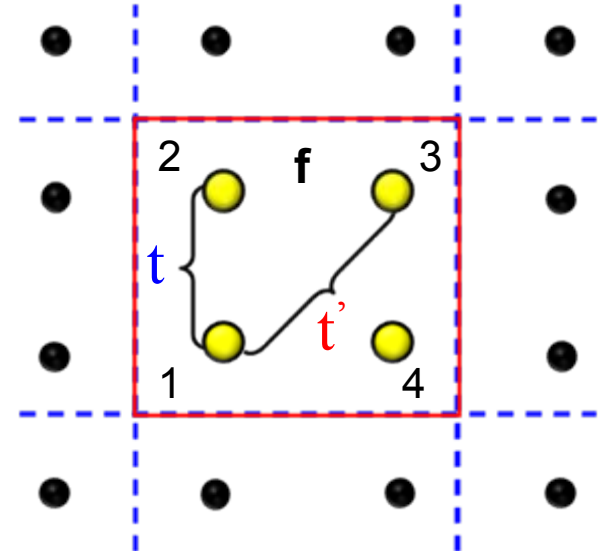


The Hubbard model for cluster 2x2

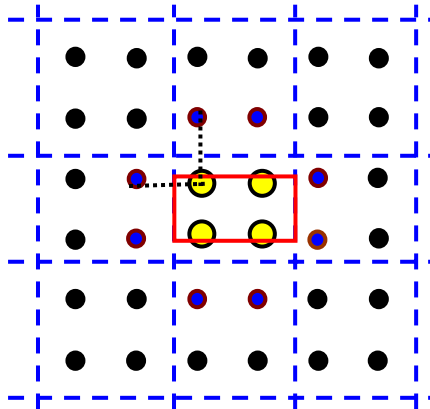
$$H = \sum_f H_0^c(f) + \sum_{f \neq g} H_1^c(f, g), \quad (4)$$

$$H_1^c = H_t^c + H_{t'}^c + H_{t''}^c,$$

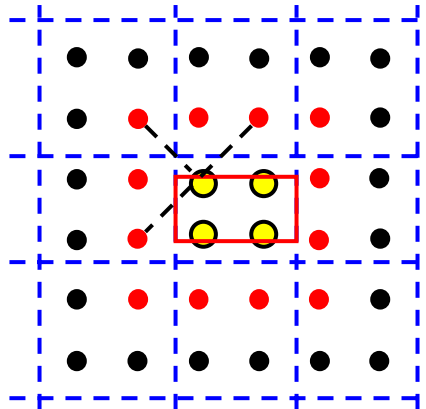
$H_0^c \Rightarrow$



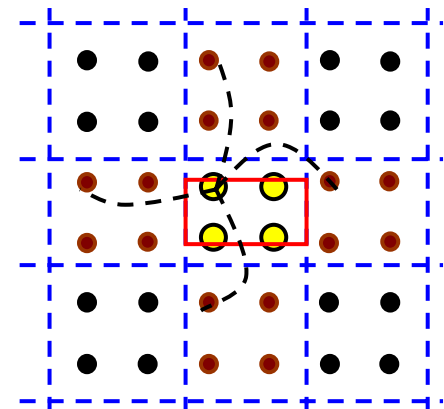
H_t^c



$H_{t'}^c$



$H_{t''}^c$



Exact diagonalization

$$H_0^c(f) \rightarrow \left\{ \varepsilon_p, |p\rangle \equiv |m, N_e\rangle \right\}$$

$$|p\rangle = \sum_n C_n |\varphi_n\rangle \quad (5)$$

Hubbard X-operators

$$X_f^\alpha \equiv X_f^{pq} = |p\rangle\langle q| = |m, N_e\rangle\langle m', N_e'| \quad (6)$$

$$a_{f\sigma i} = \sum_\alpha \gamma_{i\sigma}(\alpha) X_f^\alpha, \quad \gamma_{i\sigma}(\alpha) = \langle m', N_e - 1 | a_{i\sigma} | m, N_e \rangle \quad (7)$$

where f is the cluster index, i is the site index (1, 2, 3, 4), σ is an electron spin at site.

The sum rule

$$[a_{f\sigma i}, a_{f\sigma i}^+]_+ = 1$$

$$\sum_\alpha |\gamma_{i\sigma}(\alpha)|^2 F(\alpha) \equiv f,$$

$$F(\alpha) = \langle X^{pp} \rangle + \langle X^{qq} \rangle$$

$$\implies f > 0.99$$

Hubbard-I


$$D^{-1}(\tilde{k}, \omega) = \left(D^0(\omega) \right)^{-1} - T(\tilde{k}) \quad (8)$$

$$D_{nm}(\tilde{k}, \omega) = \langle\langle X_{\tilde{k}}^n | X_{\tilde{k}}^{m+} \rangle\rangle_{\omega}$$

$$D_{nm}^0(\omega) = \frac{F_n}{\omega - \Omega_n} \delta_{n,m}$$

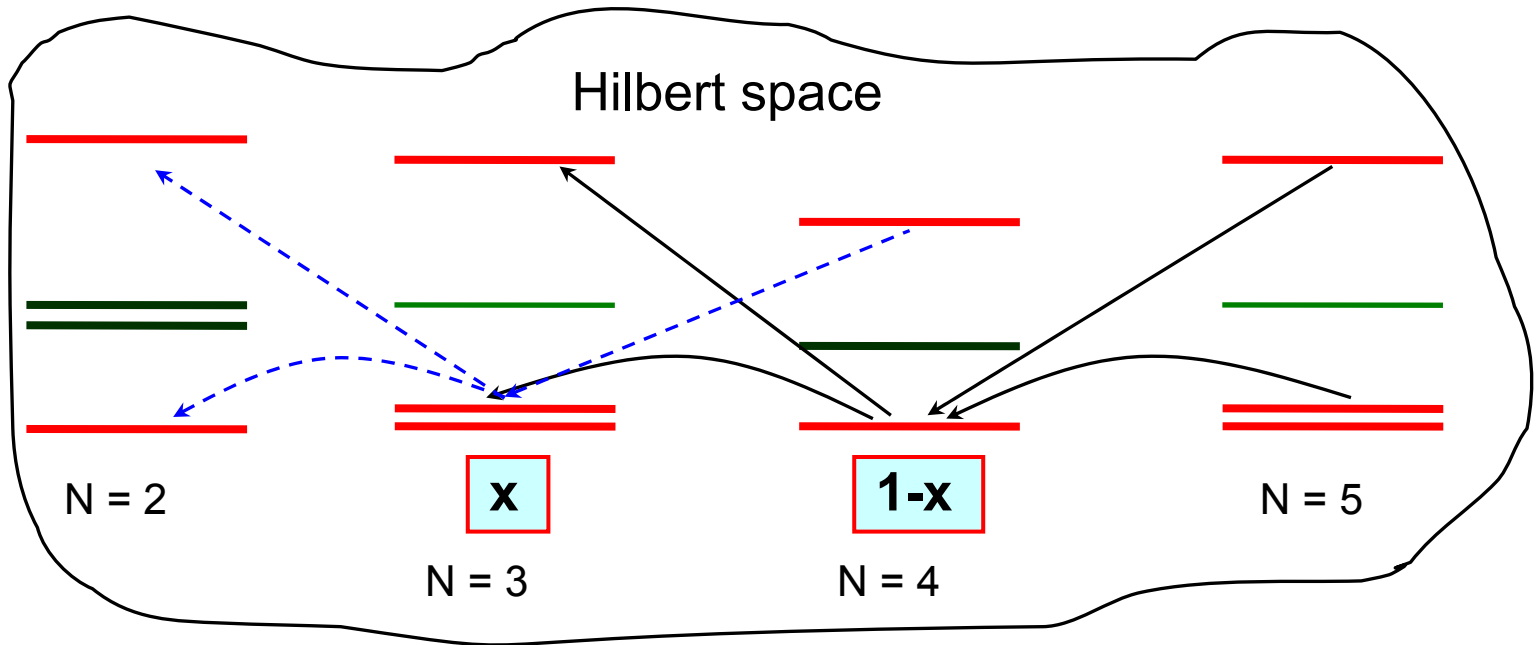
$$F_n \equiv F(pq) = \langle X^{pp} \rangle + \langle X^{qq} \rangle$$

$$G_{\sigma}(\mathbf{k}, \omega) = \frac{1}{N_c} \sum_{nm} \sum_{i,j=1}^{N_c} \gamma_{\sigma_i}(n) \gamma_{\sigma_j}(m) e^{-i\mathbf{k}(r_i - r_j)} D_{nm}(\mathbf{k}, \omega) \quad (9)$$

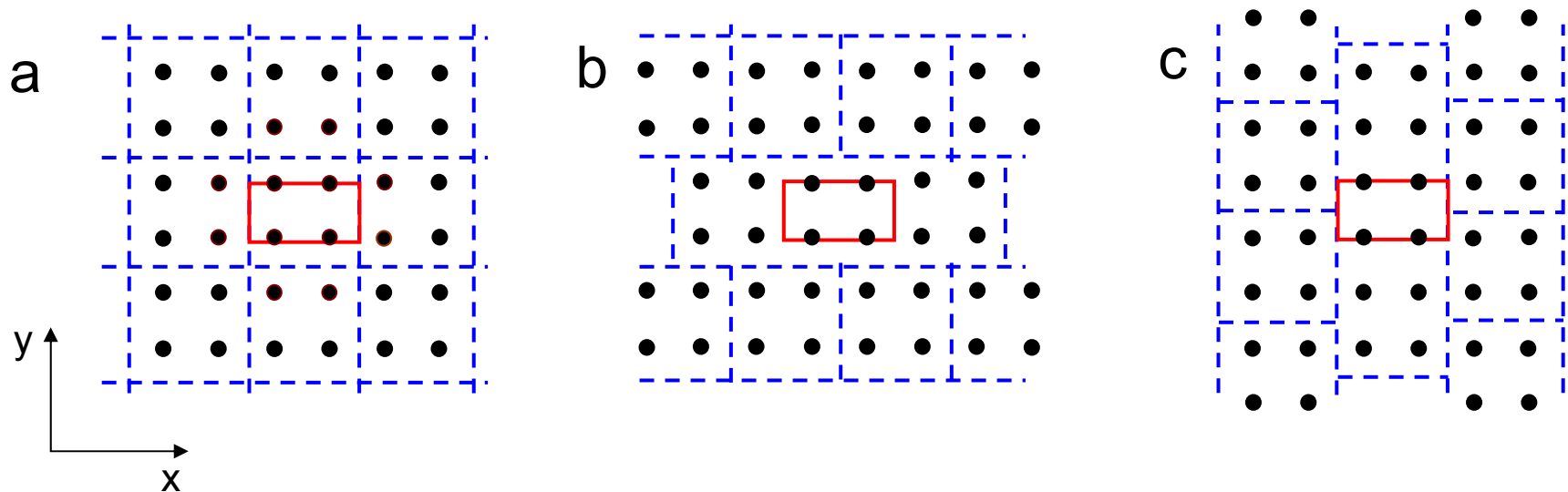
$$A_{\sigma}(k, \omega) = -\frac{1}{\pi} \lim_{\eta \rightarrow +0} (\text{Im} G_{\sigma}(k, \omega + i\eta + \mu)) \quad (10)$$

$$\int d\omega A_{\sigma}(k, \omega) = \langle [a_{k\sigma}, a_{k\sigma}^+]_{+} \rangle = f \quad (11)$$

The scheme of doping

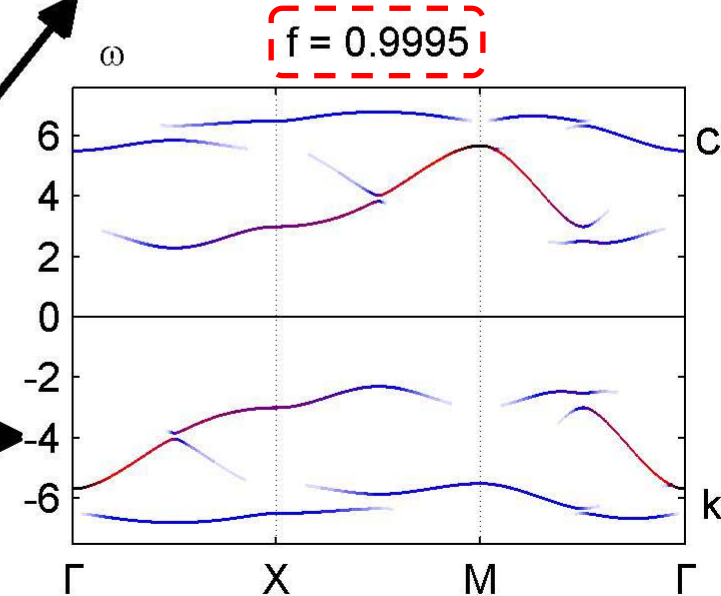
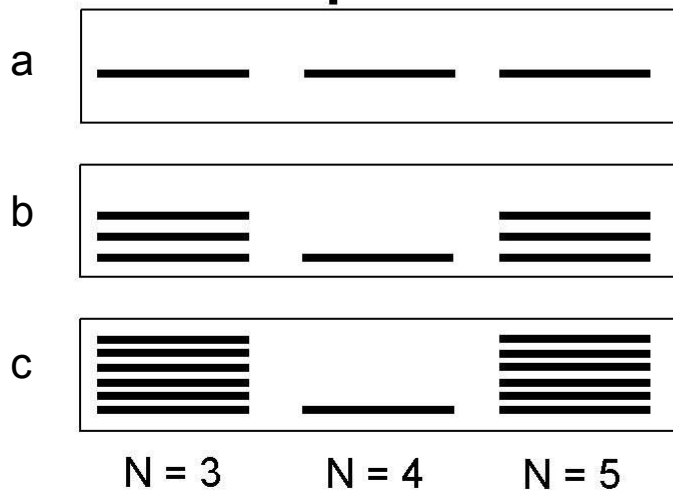
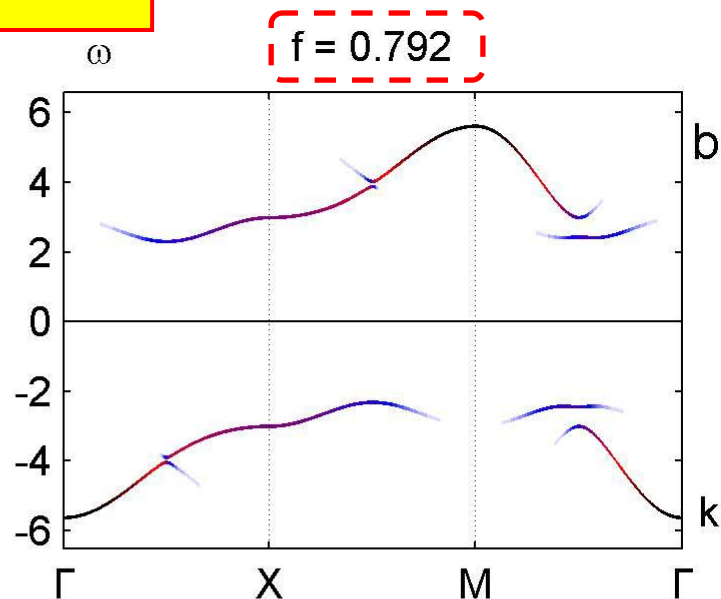
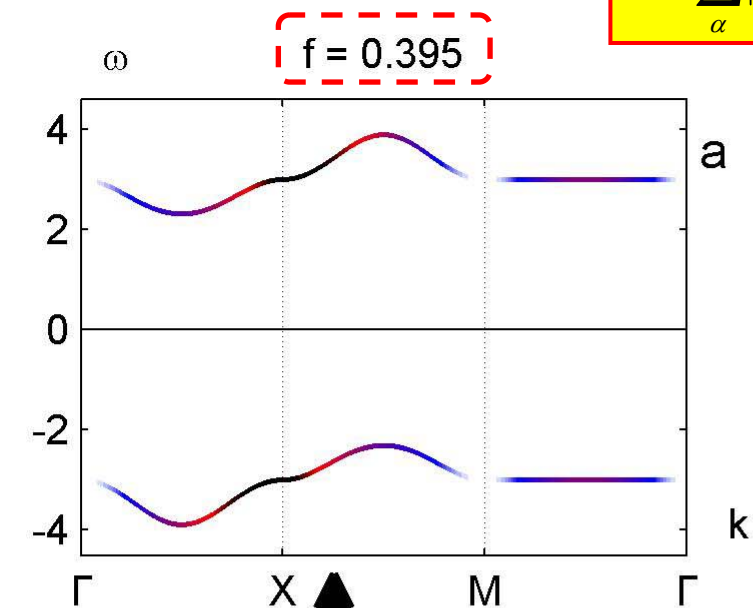


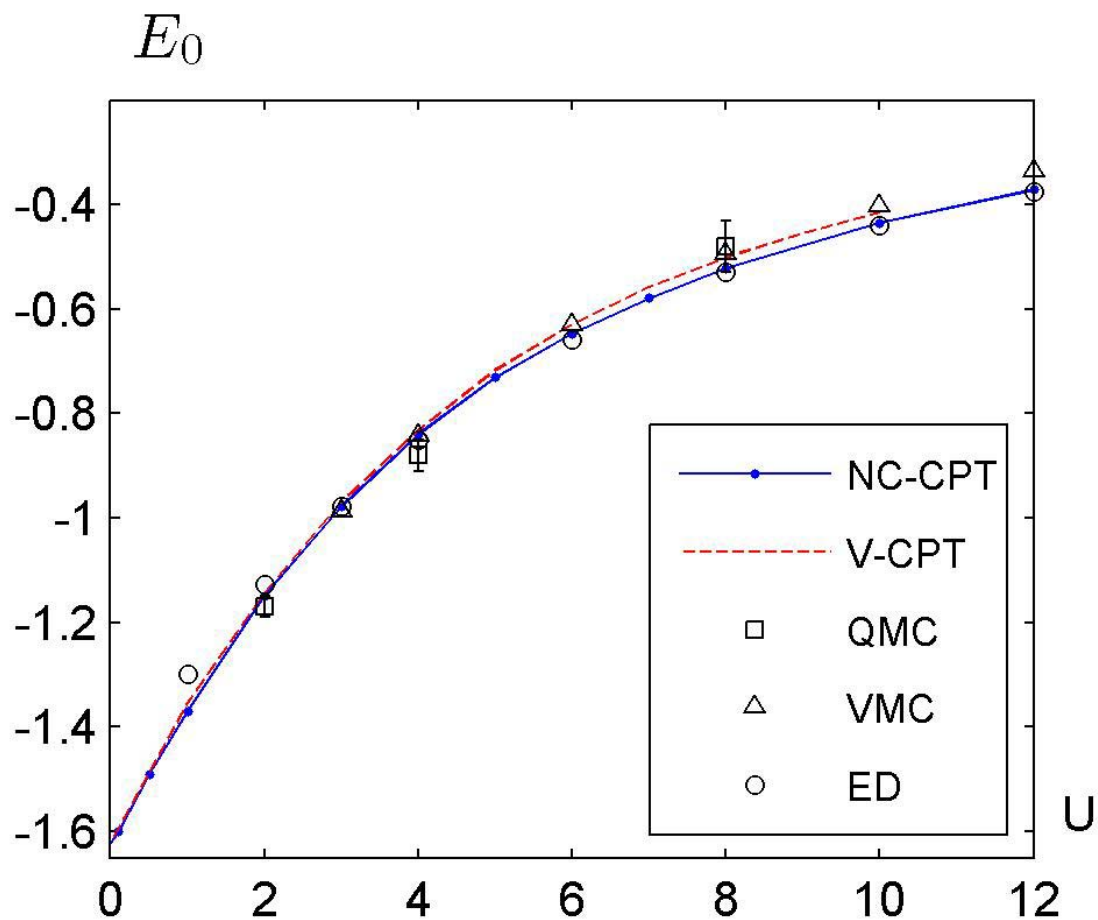
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$$T_{\alpha\beta}(\mathbf{k}) = \frac{1}{2} \left(T_{\alpha\beta}^b(\mathbf{k}) + T_{\alpha\beta}^c(\mathbf{k}) \right), \quad (12)$$

$$f = \sum_{\alpha} |\gamma_{i\sigma}(\alpha)|^2 F(\alpha)$$



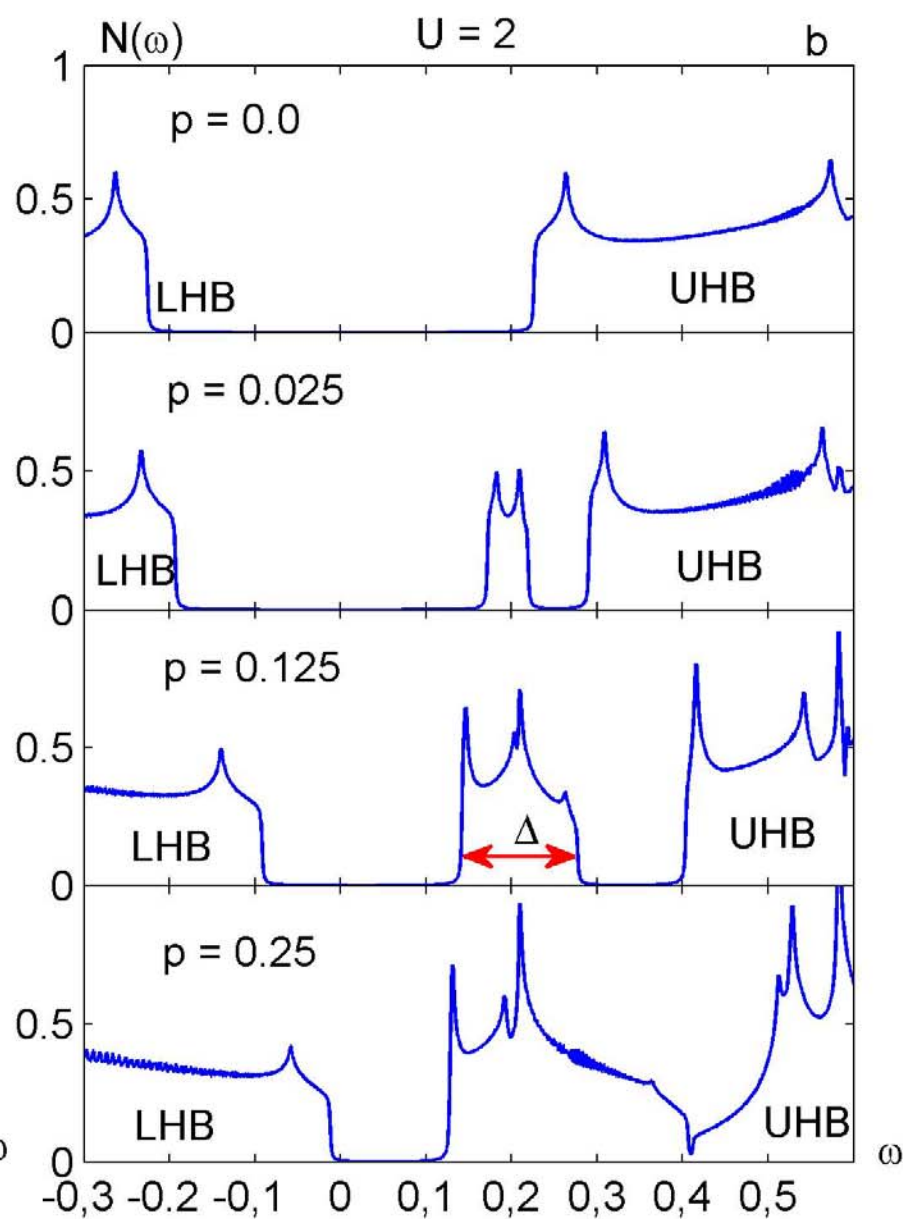
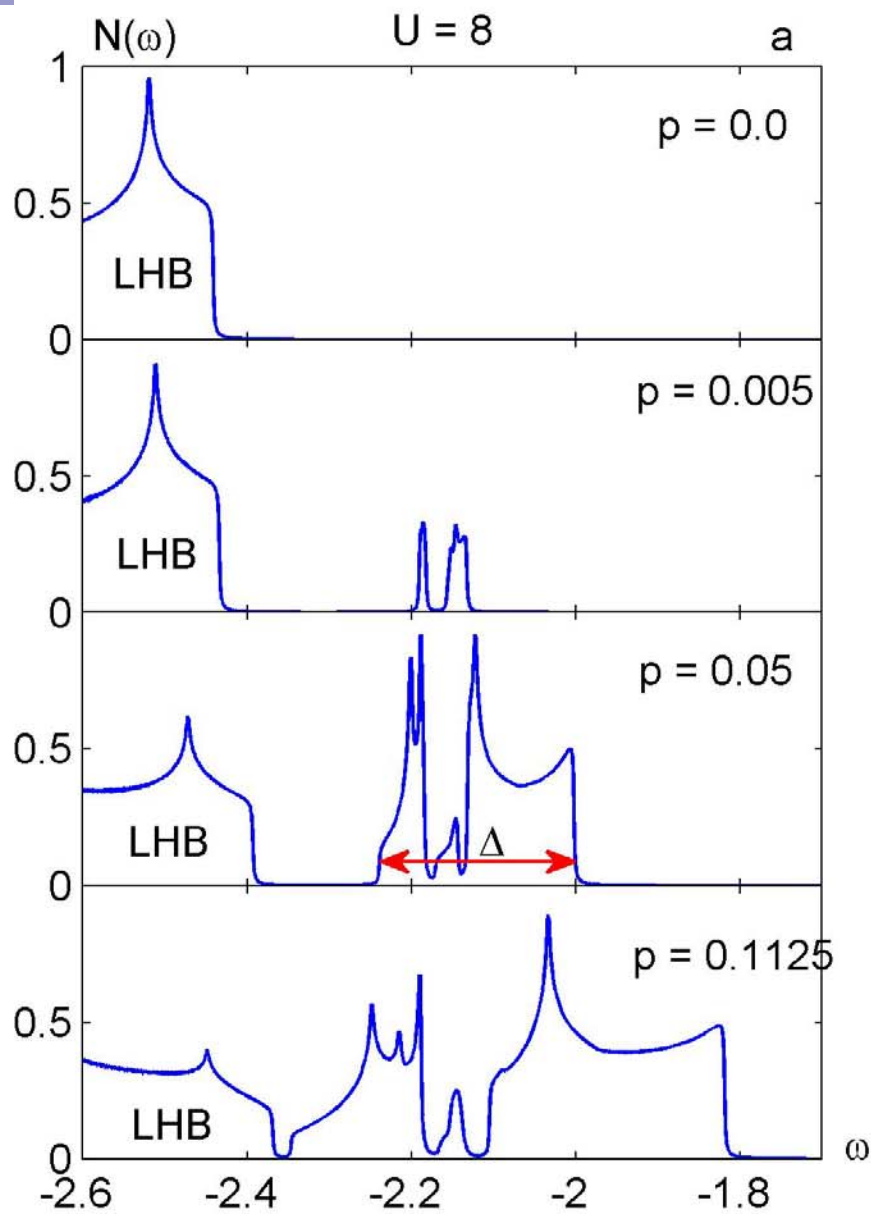


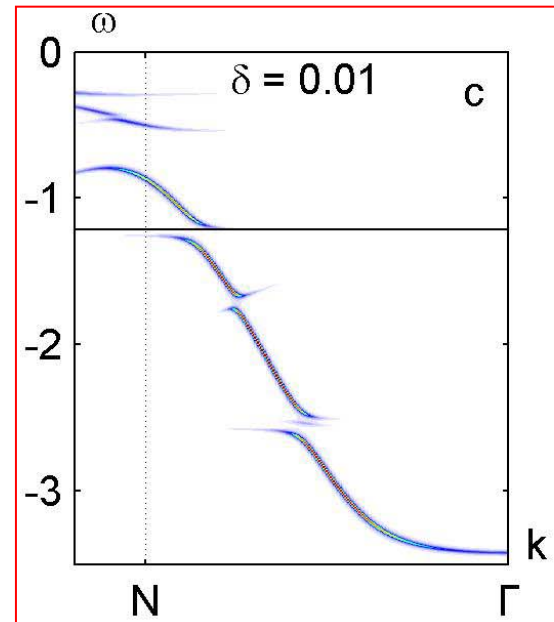
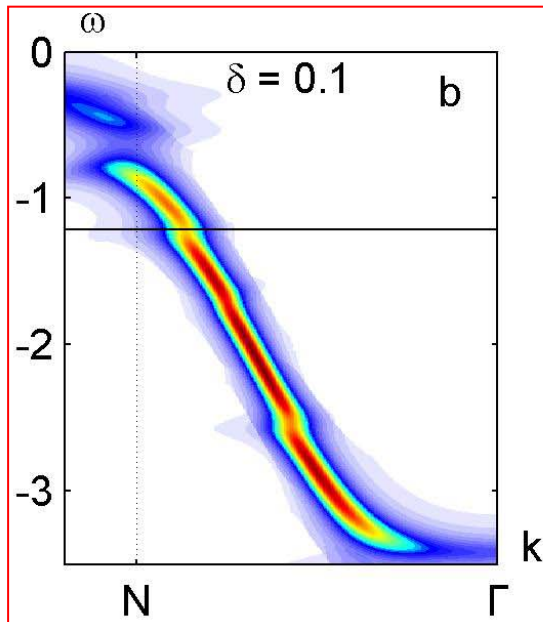
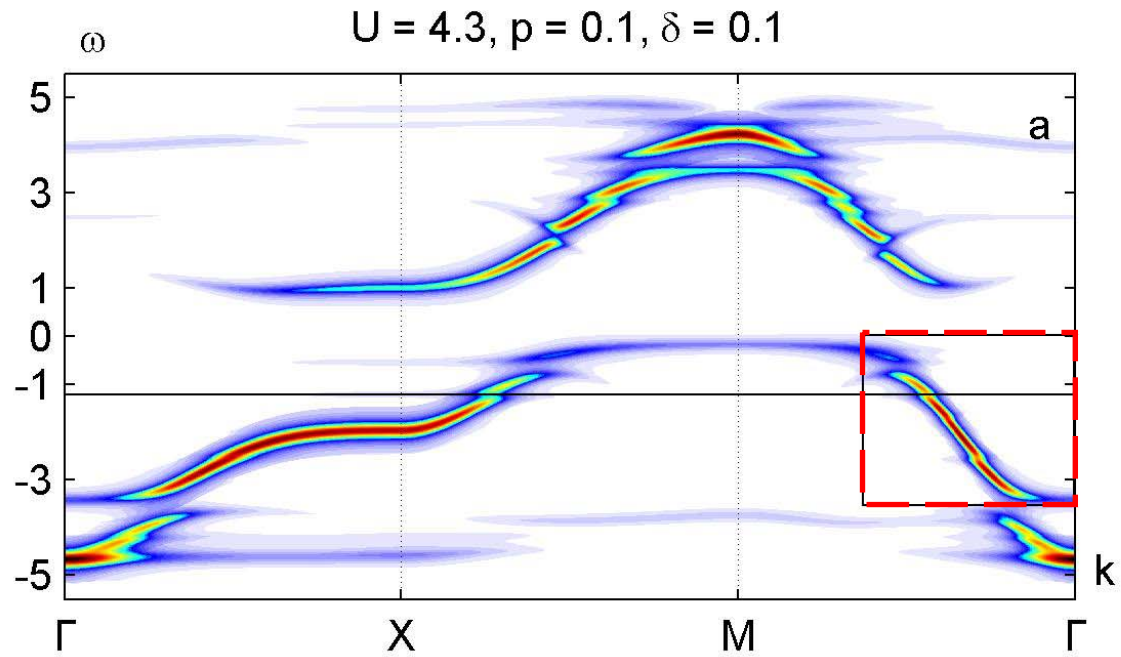
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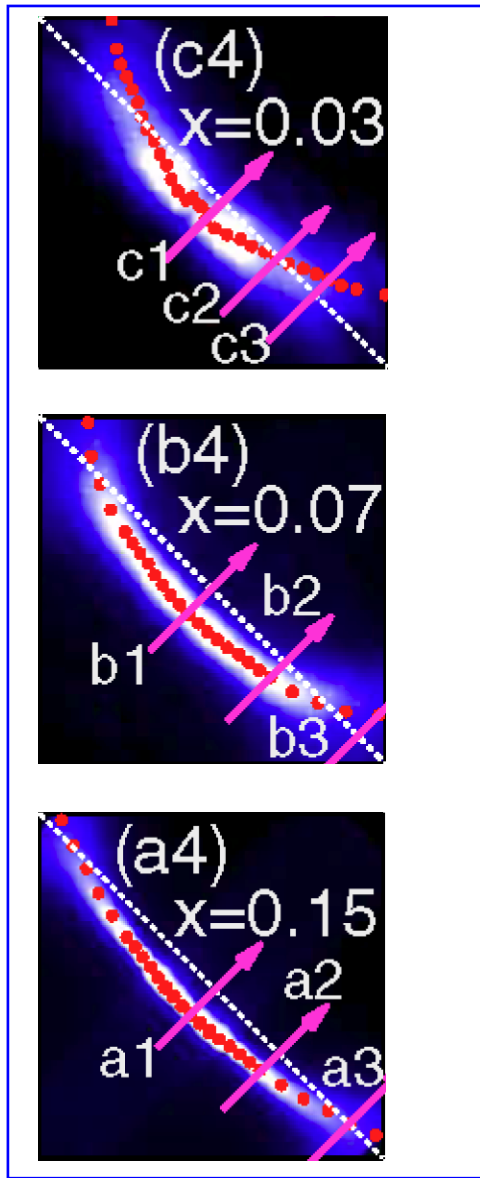
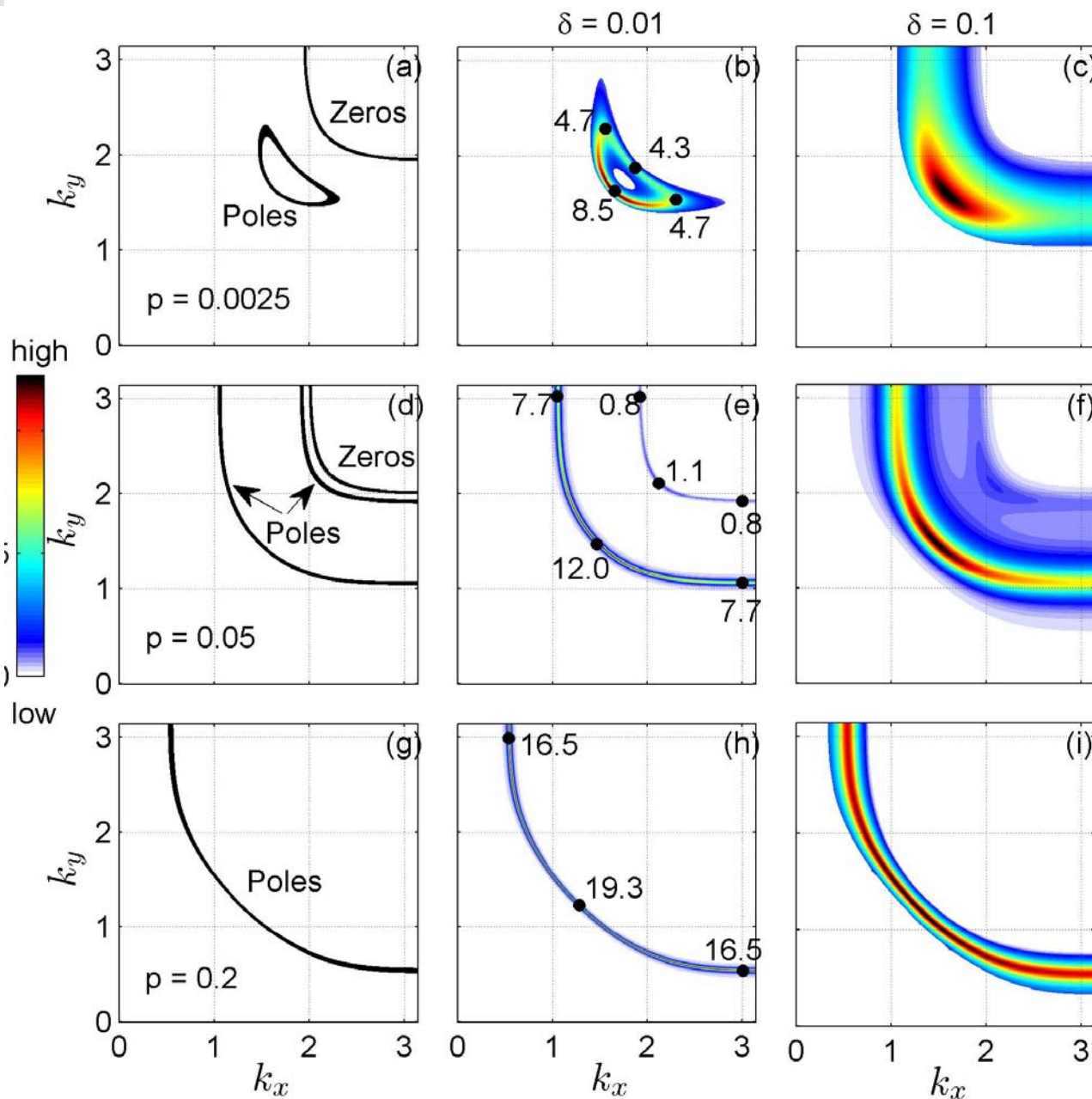
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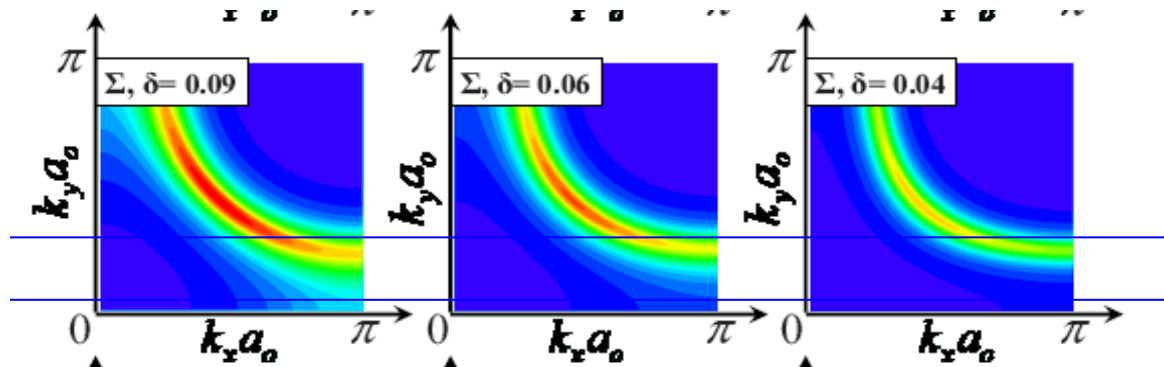
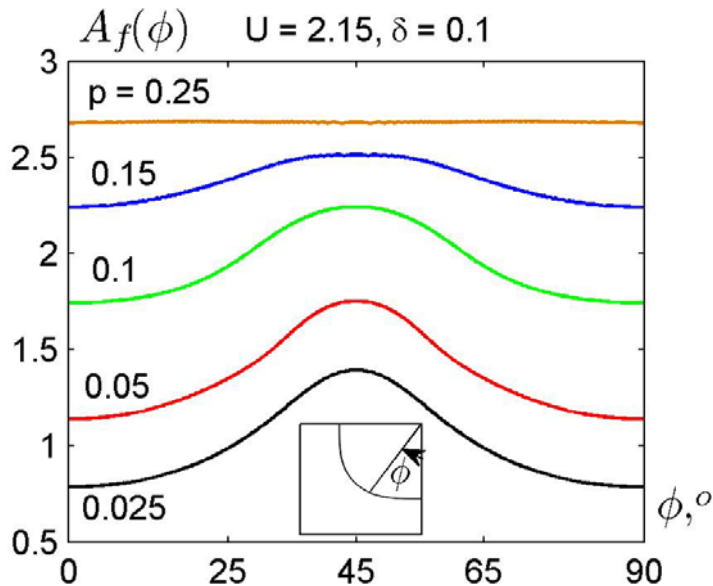
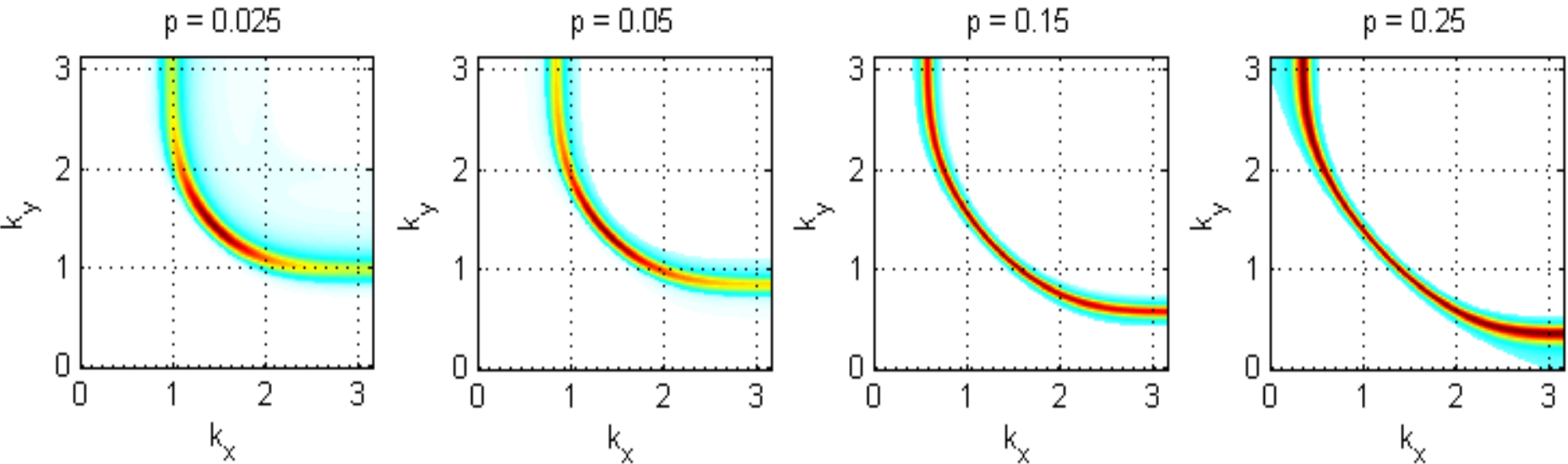




$\delta = 20 \text{ meV} / 250 = 0.08 t$



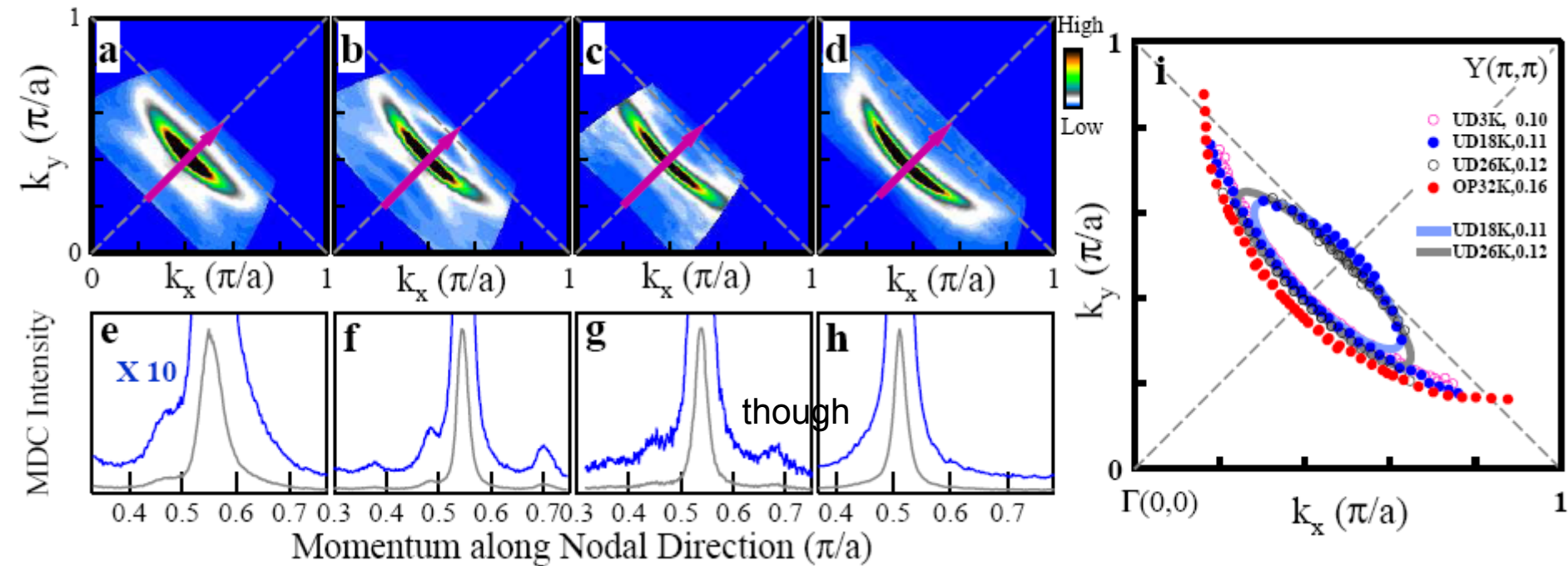
$$\delta = 0.1$$



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CDMFT

Doping evolution of Fermi surface topology in La-Bi2201

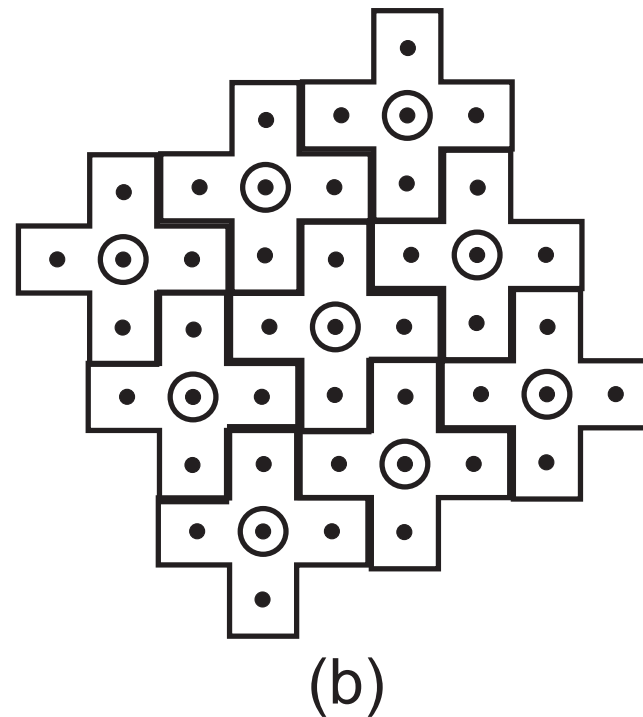
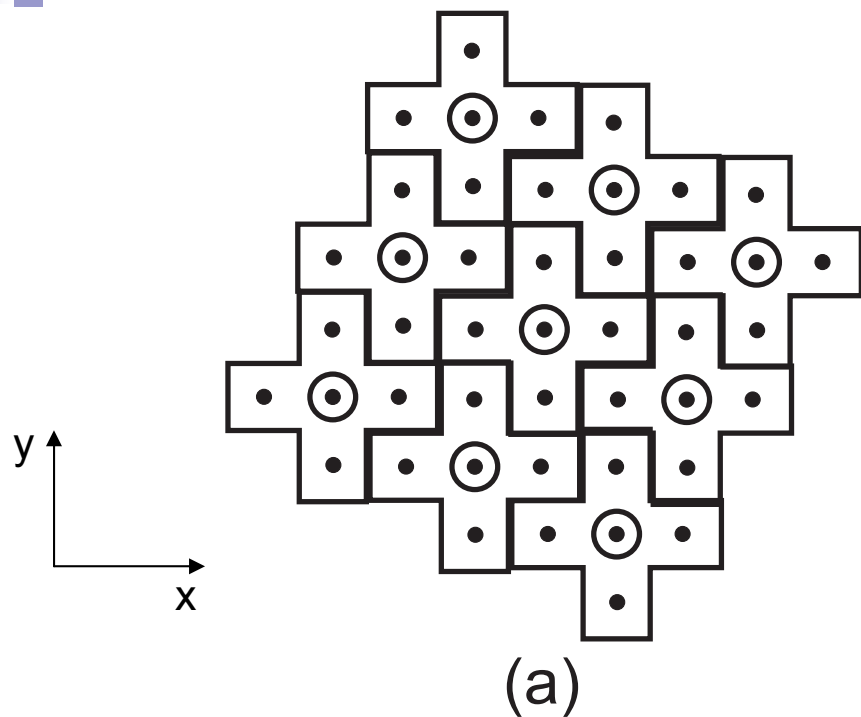


(a-d) show Fermi surface mapping of UD3K (underdoped, T_c 3 K), UD18K, UD26K (underdoped, T_c = 26 K) and OP32K (optimally-doped, T_c = 32 K) samples, respectively.

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Our calculations of the spectral weight maps have found the strong dependence on the spectral line width (the broadening parameter) that is related to the experimental resolution. The spectral weight maps with better resolution can give full details about the Fermi surface

Thank you for your attention!



$$T_{nm}(\mathbf{k}) = \frac{1}{2} \left(T_{nm}^{(a)}(\mathbf{k}) + T_{nm}^{(b)}(\mathbf{k}) \right)$$

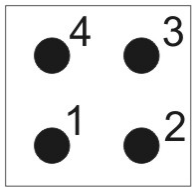
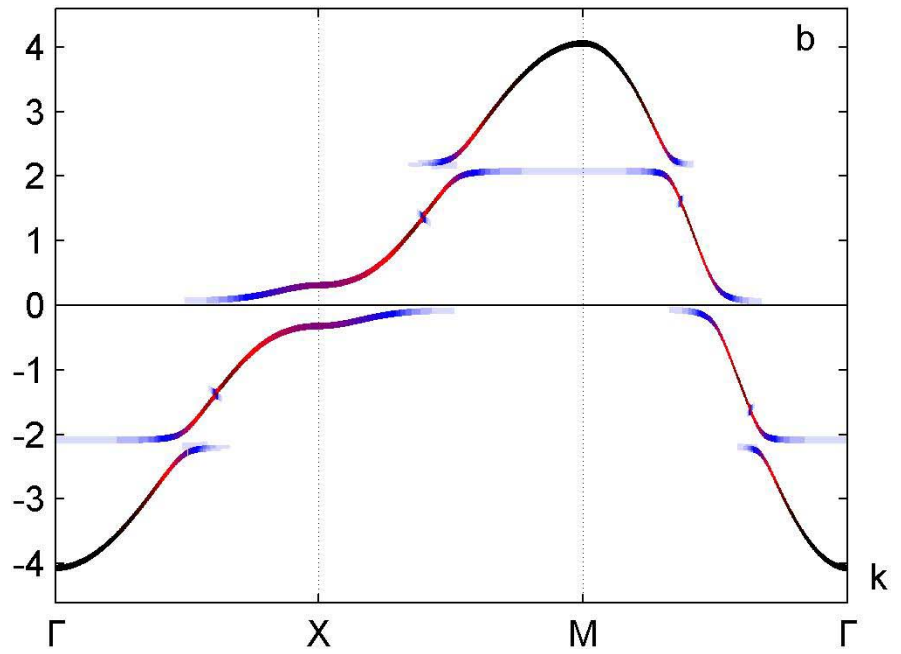
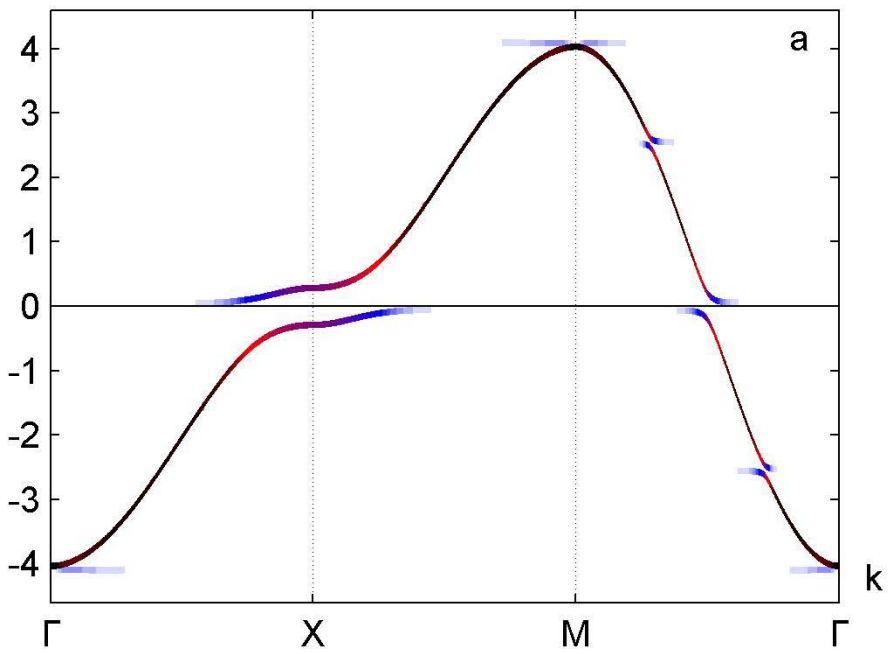
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The band structure

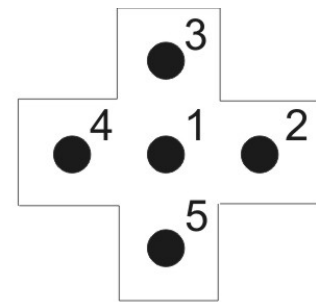
$U = 1$

$N_c = 4$

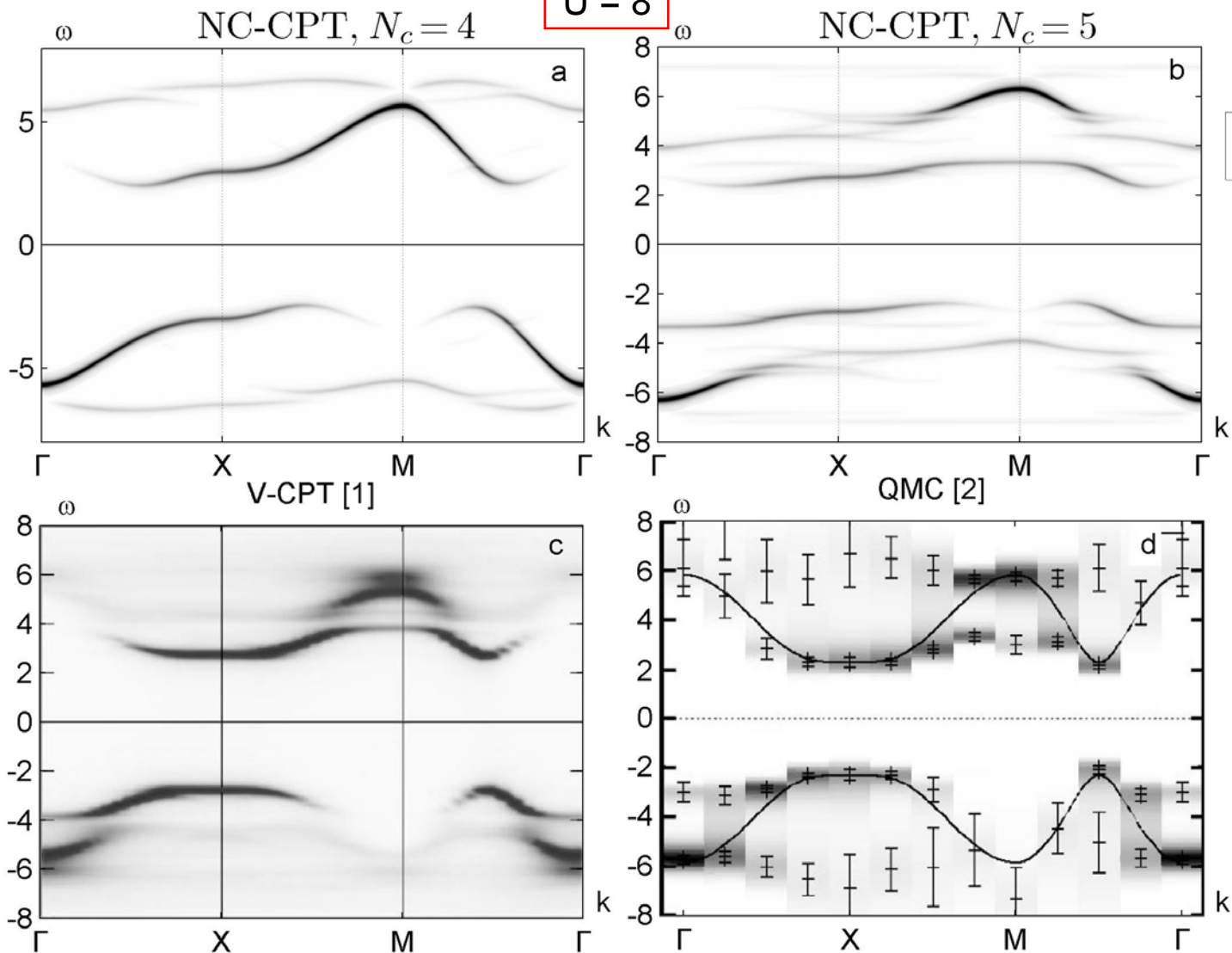
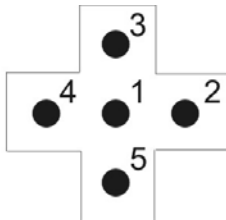
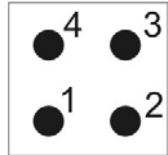
$N_c = 5$



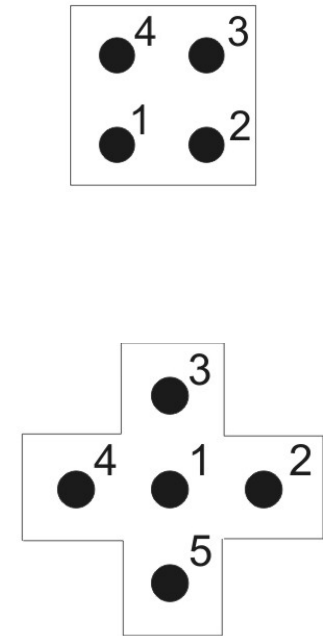
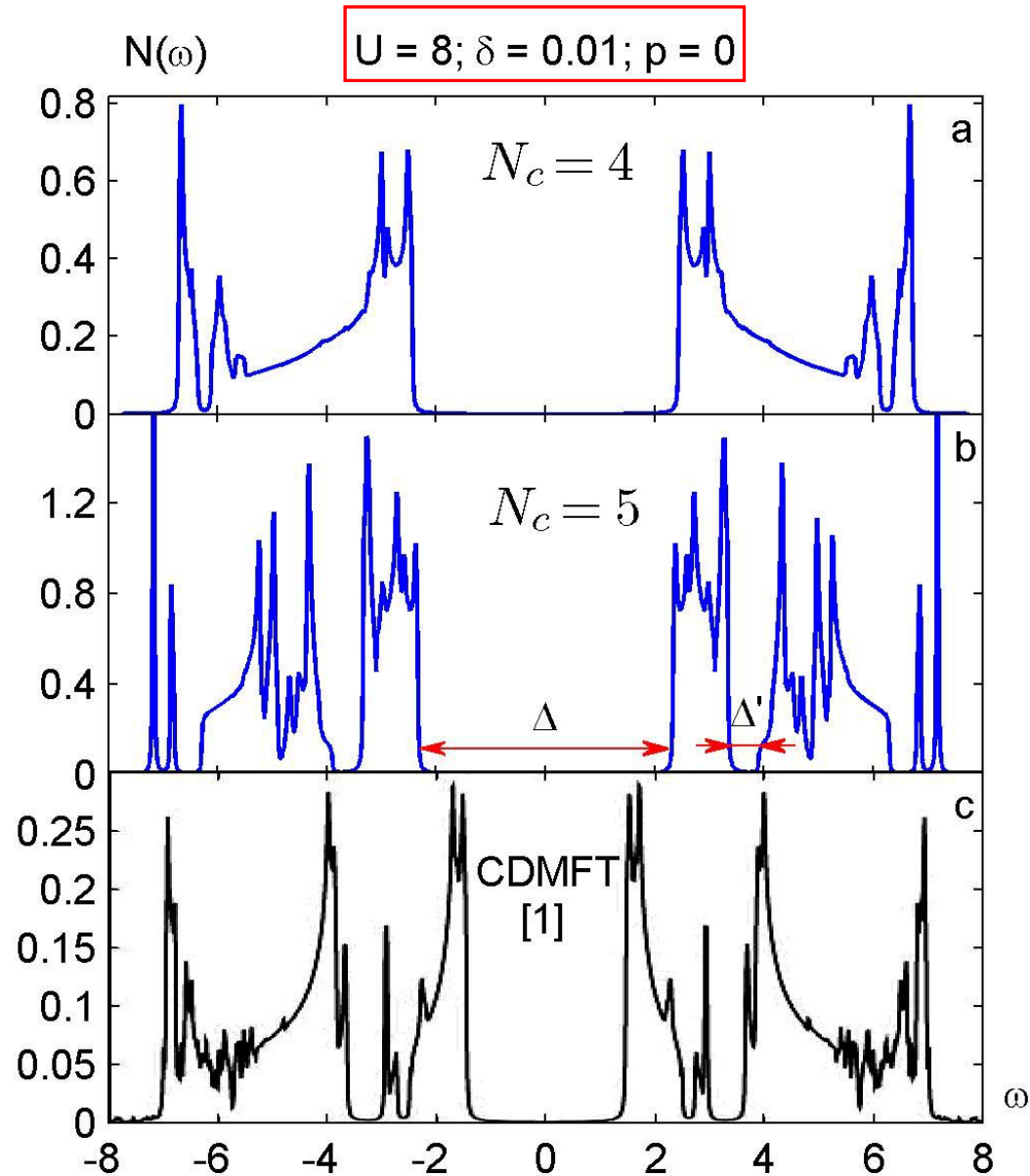
Size effect

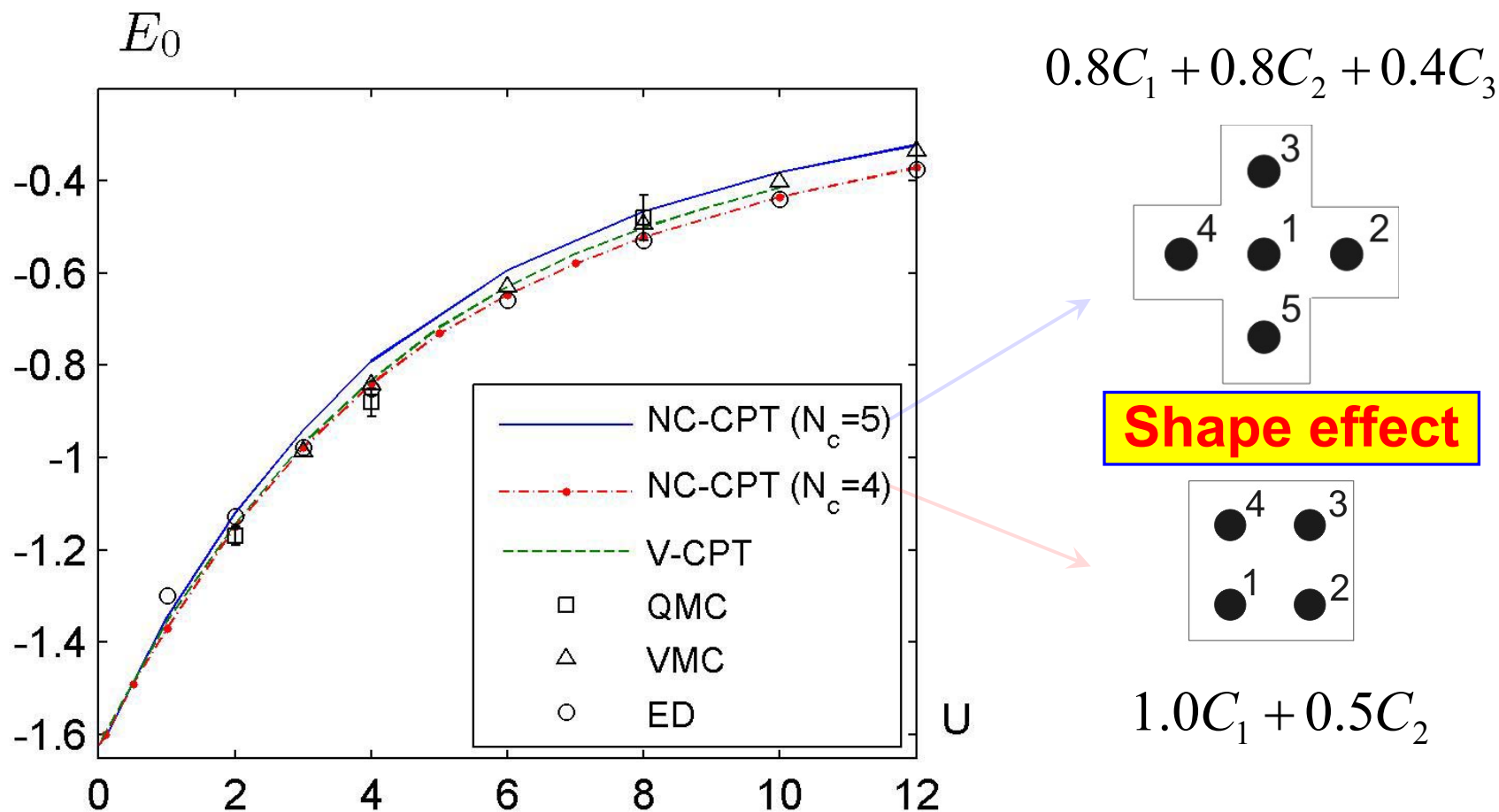


$U = 8$



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The doping transformation of the Fermi surface

