



San Sebastián

Tailoring Surface and Thin Film electronic states with Curved Surfaces

J. Enrique Ortega

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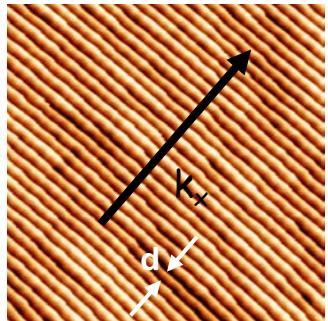
²*Centro de Física de Materiales, Materials Physics Center CSIC-UPV/EHU*

³*Donostia International Physics Center (DIPC)*

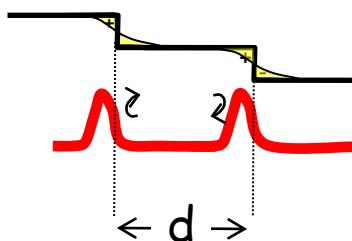
San Sebastián
Spain

Scattering at noble metal superlattices probed with ARPES

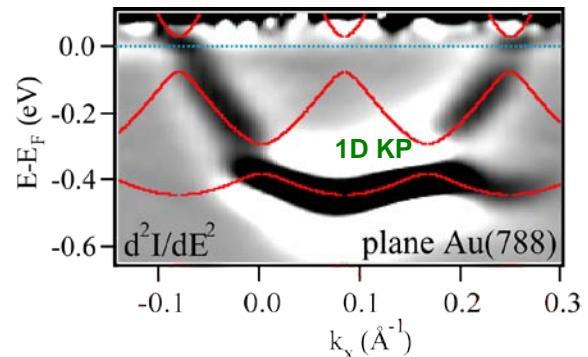
Stepped surfaces



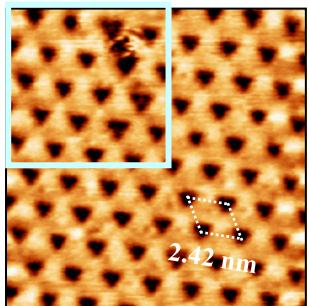
1D-lattice



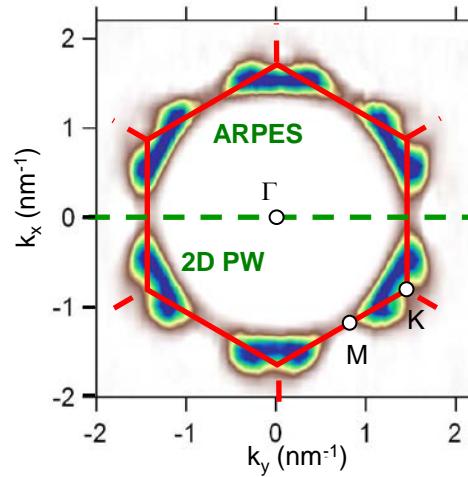
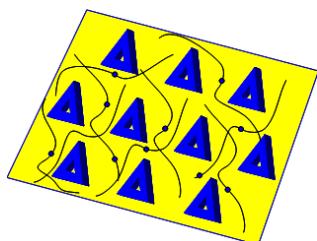
ARPES



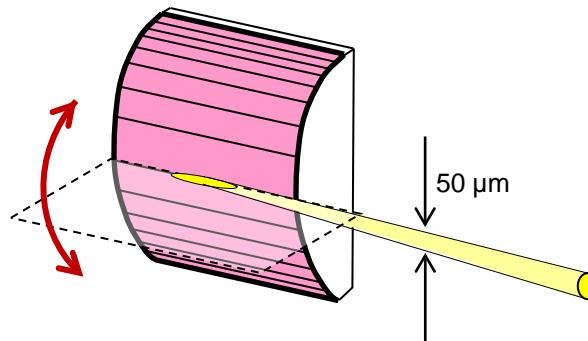
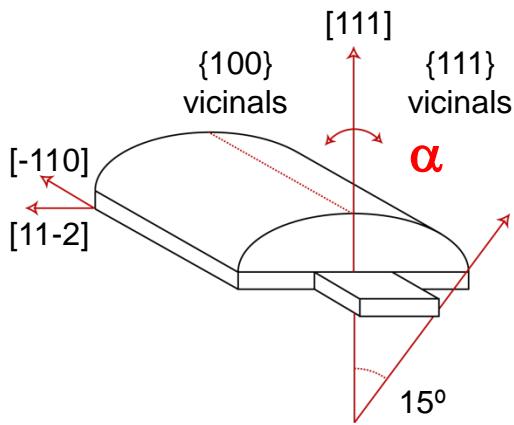
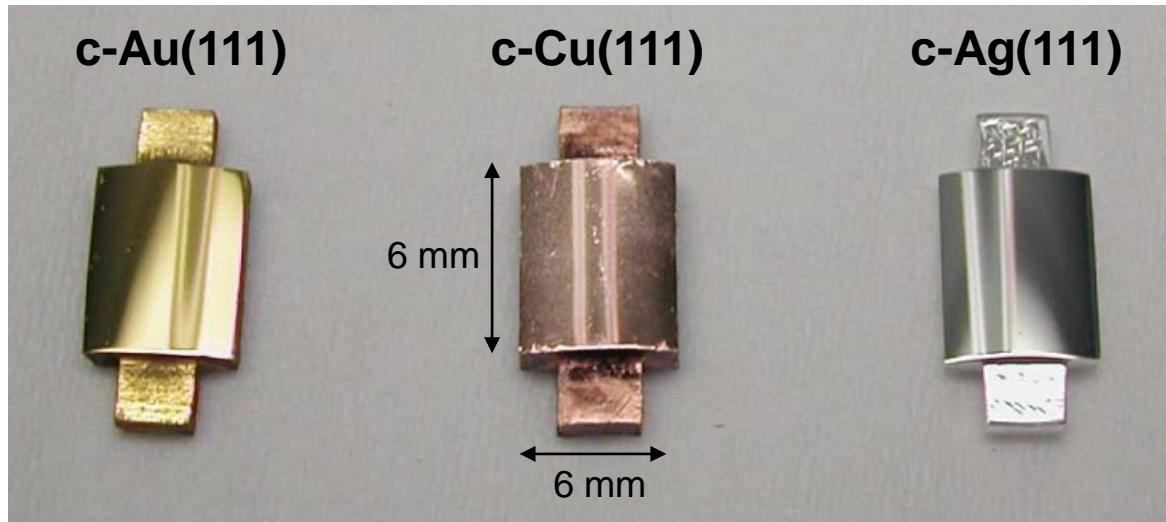
Dislocation networks



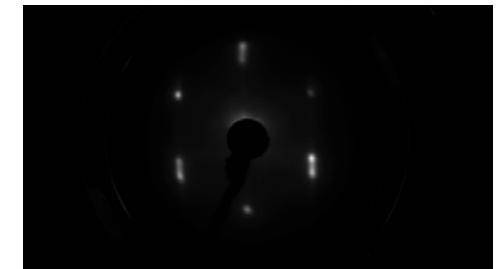
2D-lattice



Tunable step lattices: curved surfaces



c-Ni(111)



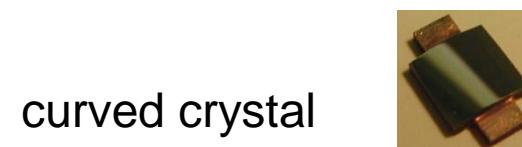
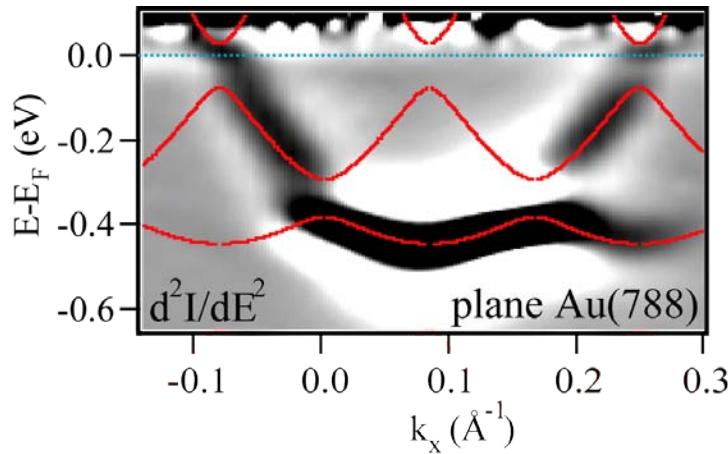
We aim at fabricating samples that can be mounted using standard holders and be prepared through regular cleaning procedures

Flat versus curved crystals in ARPES

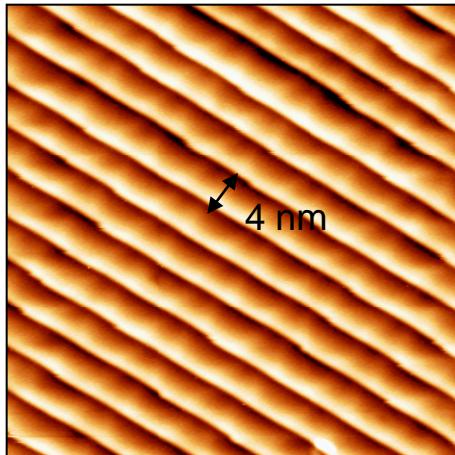
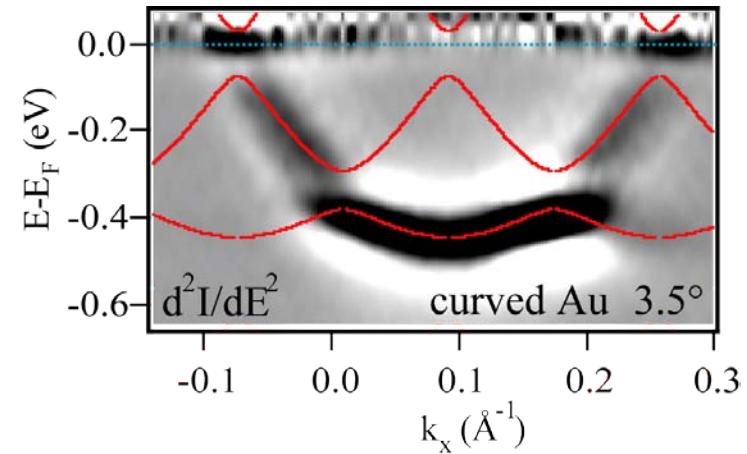


Au(788) $\alpha=3.5$

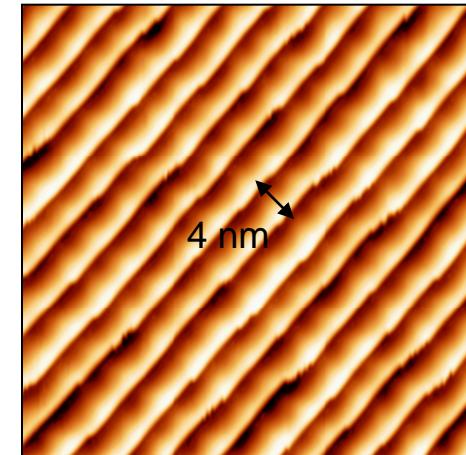
flat crystal



curved crystal

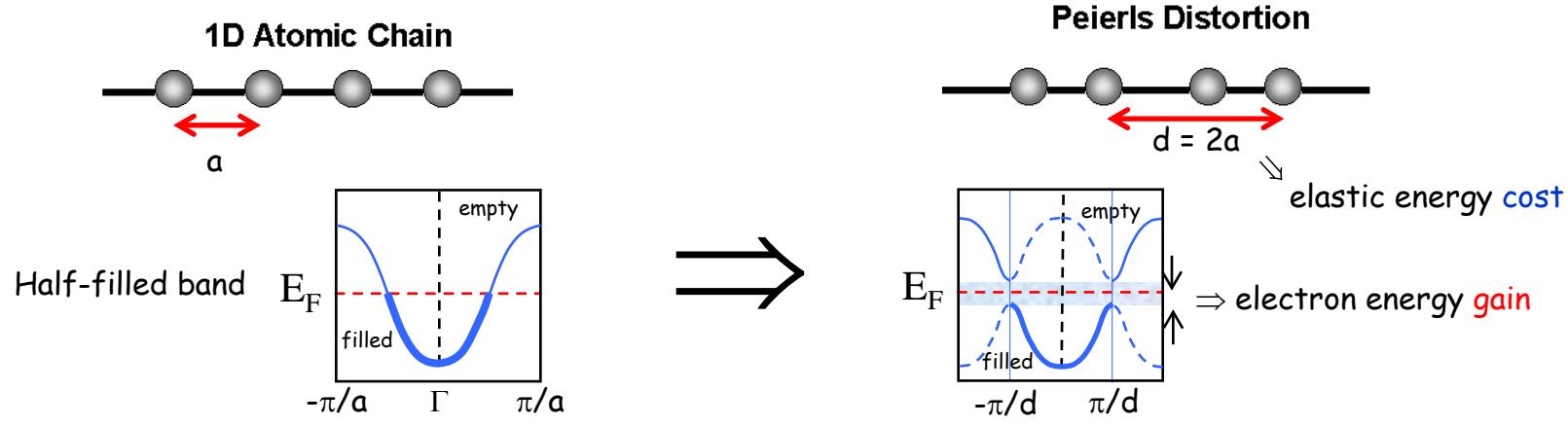


ARPES measurements from flat and curved crystals with the same crystallographic orientations are completely analogous

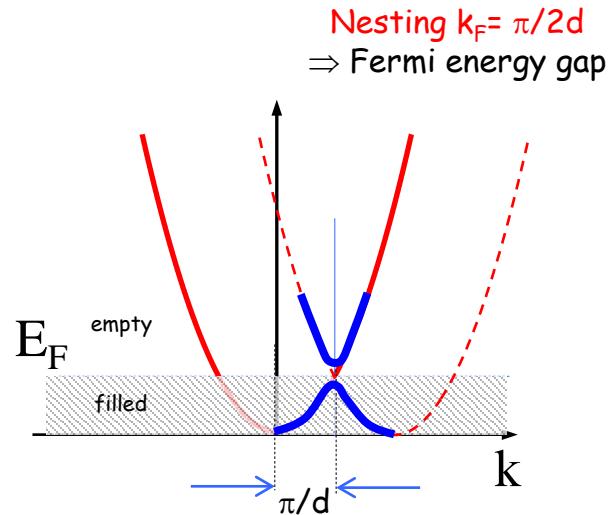


Curved surfaces for what?

Noble metals: structure/electronic interplay?



What kind of structural instabilities arise by approaching in 1D nesting of the Shockley state in a step superlattice?

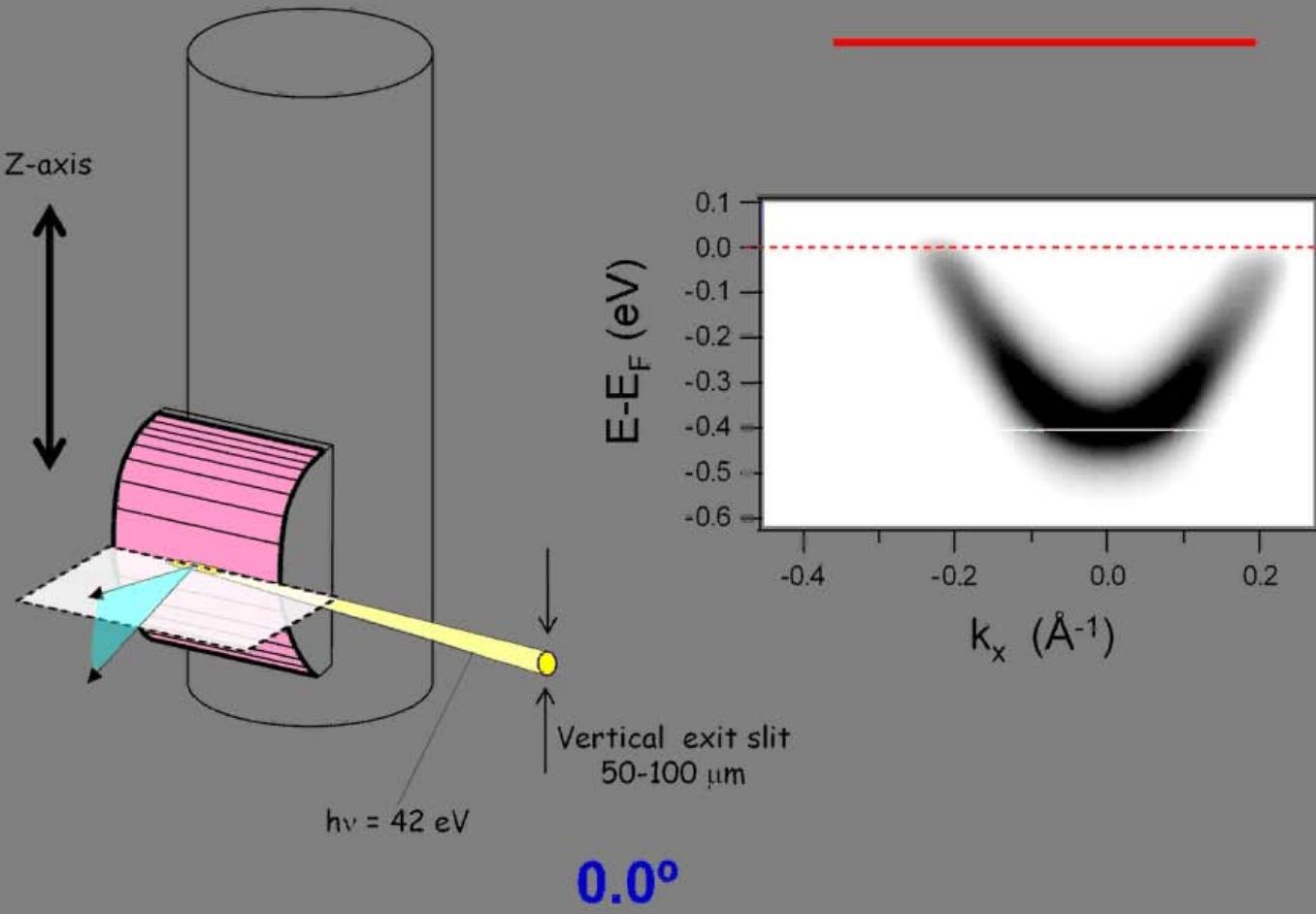


Outline

- ARPES from noble metal curved surfaces:
 - ✓ *tuning umklapps and gaps*
 - ✓ *revisiting step barrier strength*
- Thin films on curved surfaces:
 - ✓ Ag/c-Au(111): *lateral scattering in quantum-wells*
 - ✓ $BiAg_2$ /c-Ag(111): *SOC bands*

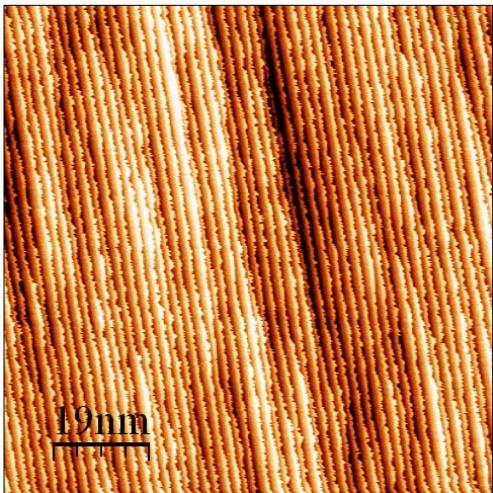
c-Cu(111)

Scanning the light beam

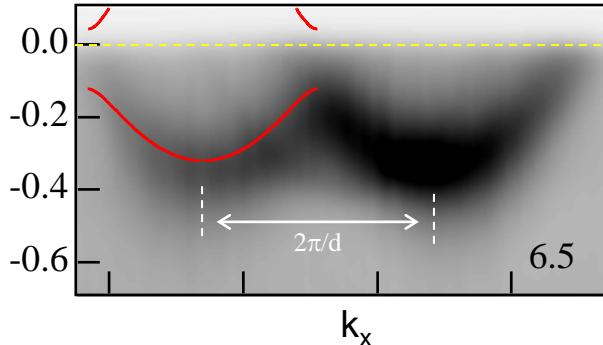


Step edge roughening in c-Cu(111)

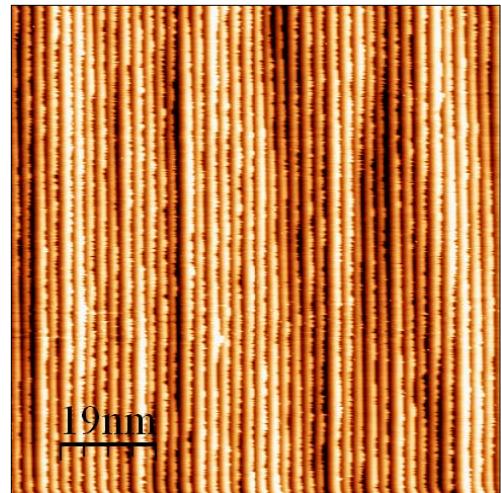
$\alpha=7.1$



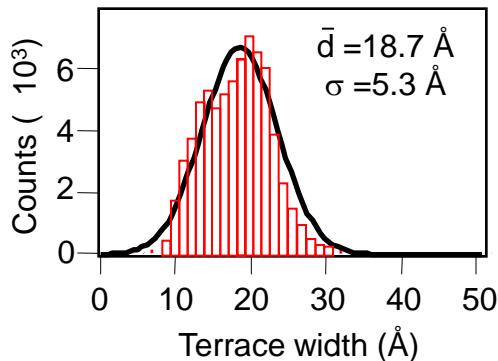
$\alpha=6.5$, $d=\lambda_F/2$
(Fermi gap)



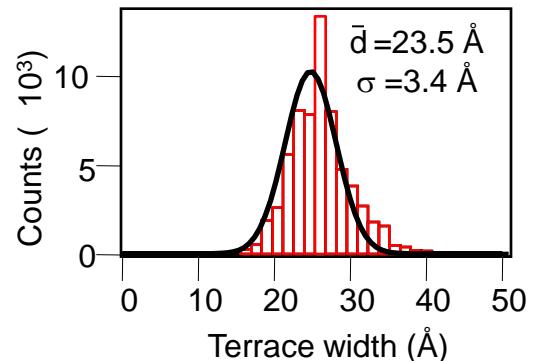
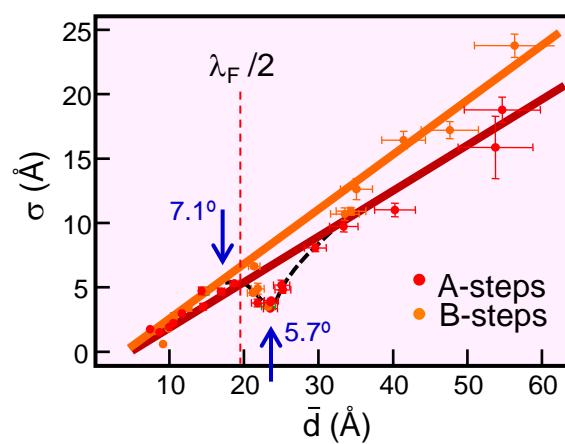
$\alpha=5.7$



J. E. Ortega *et al.*, Phys. Rev. B
83, 085411 (2011)

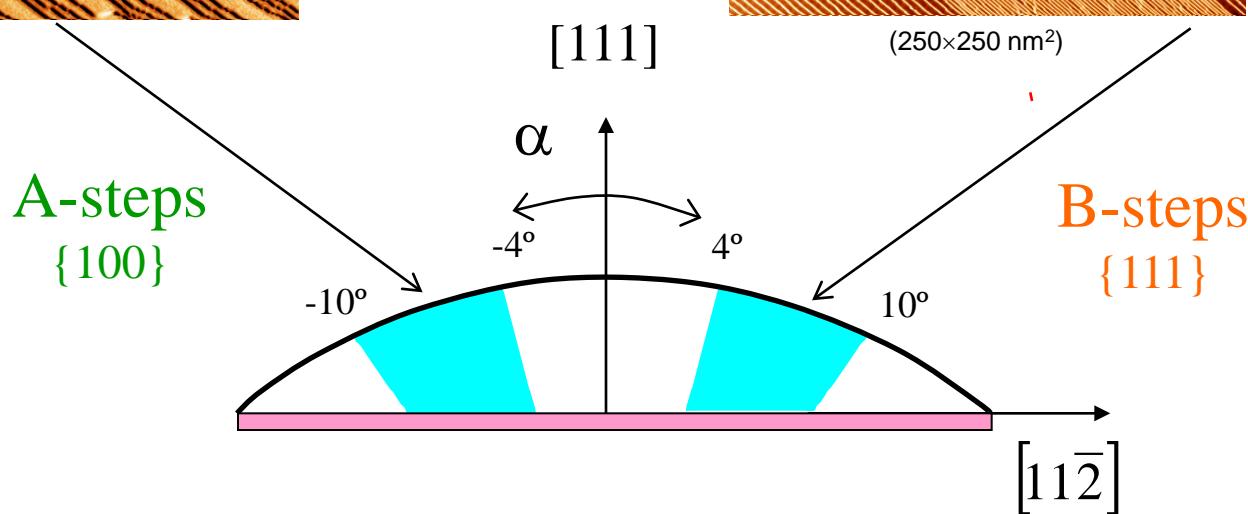
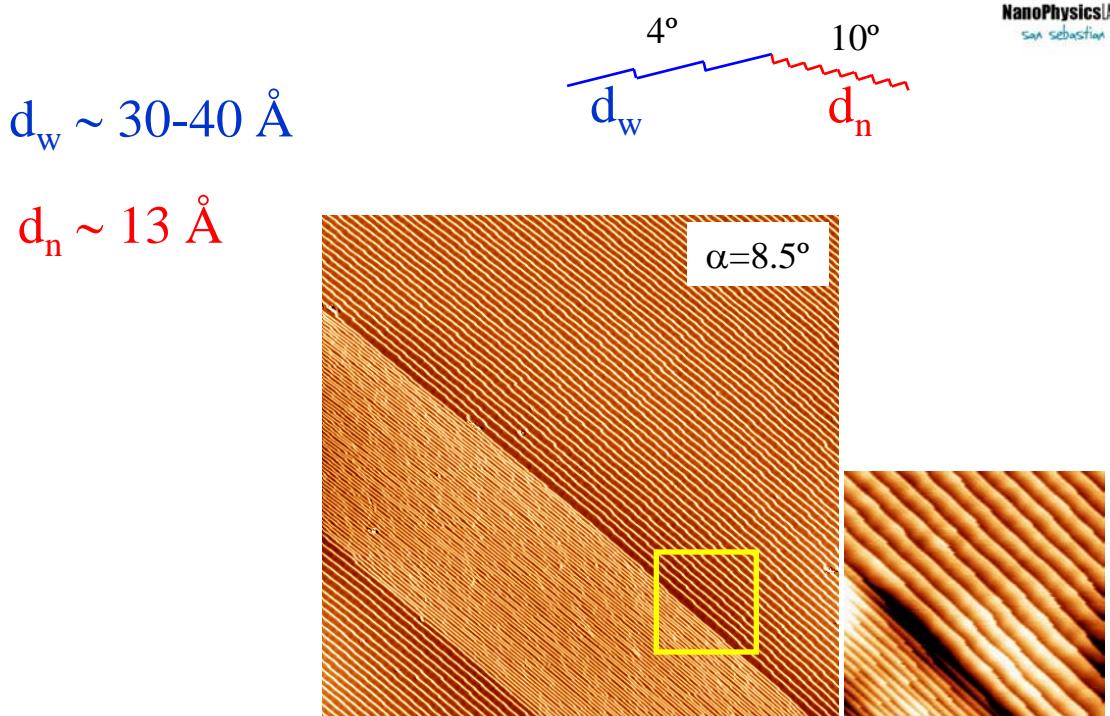
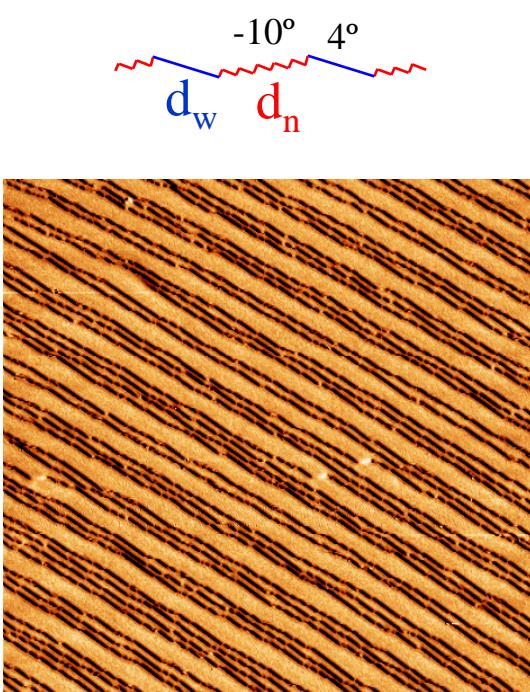


Rough step edges



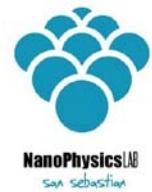
Sharp step edges

Au(111) faceting

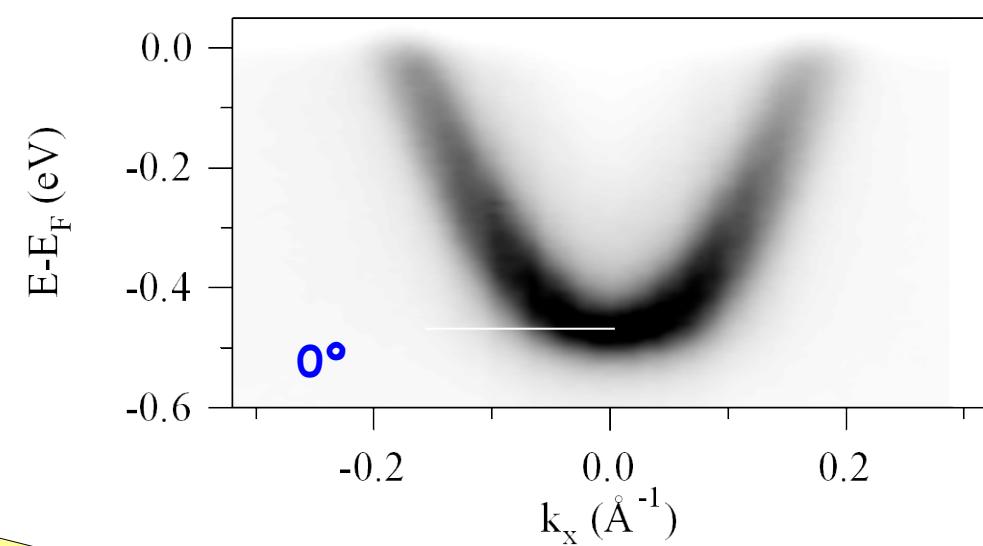
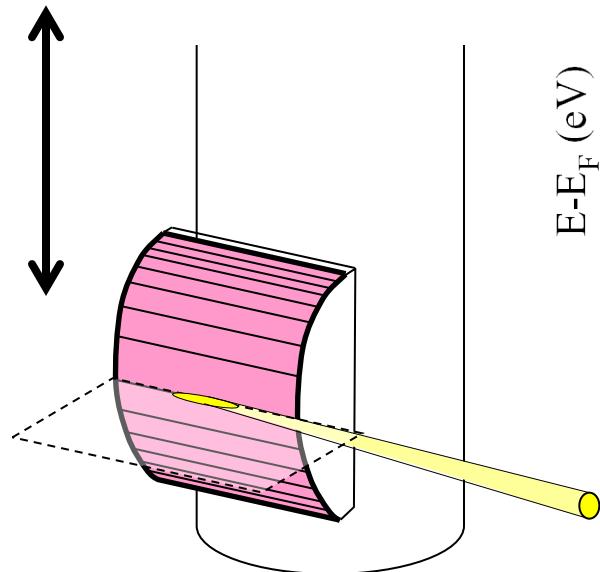


Electronic states at the faceting transition in Au(111)

M. Corso *et al.*, J. Phys. Condens. Matter **21**, 353001 (2009).

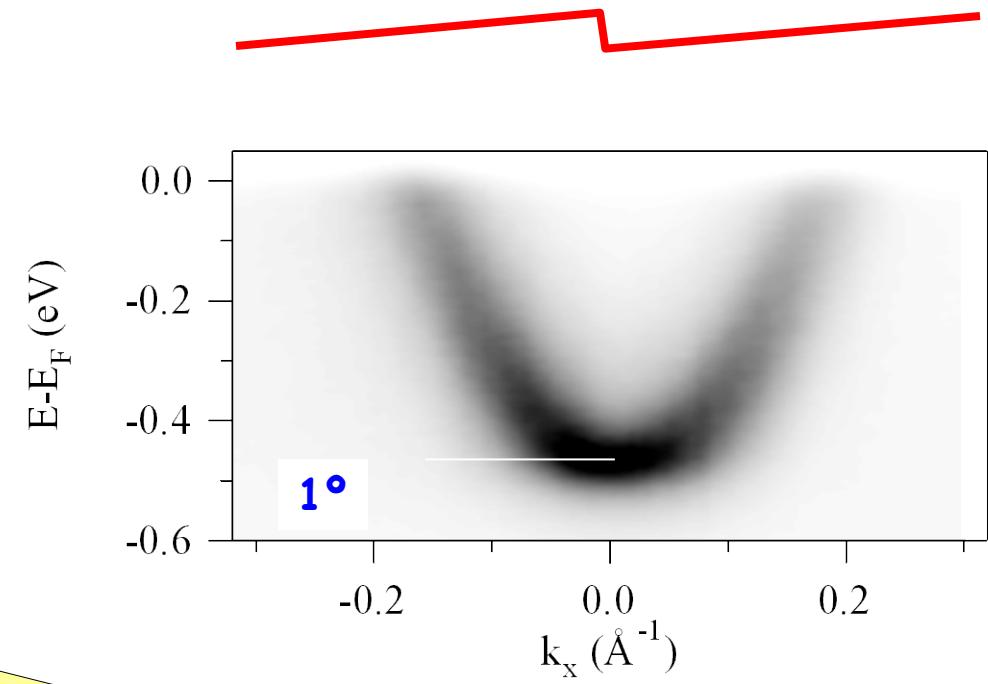
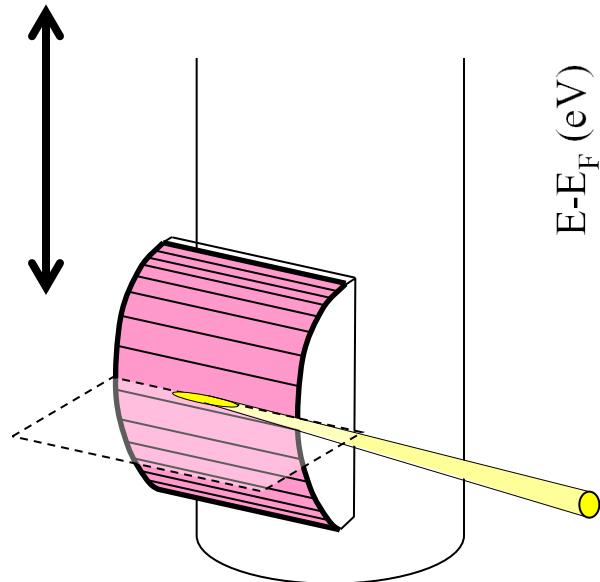


Z-axis



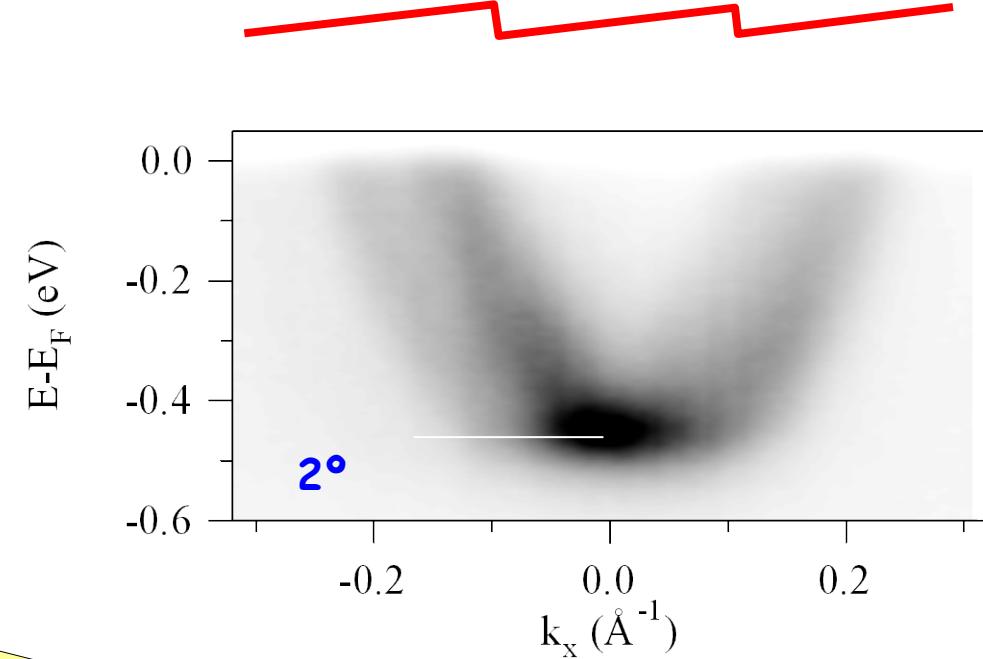
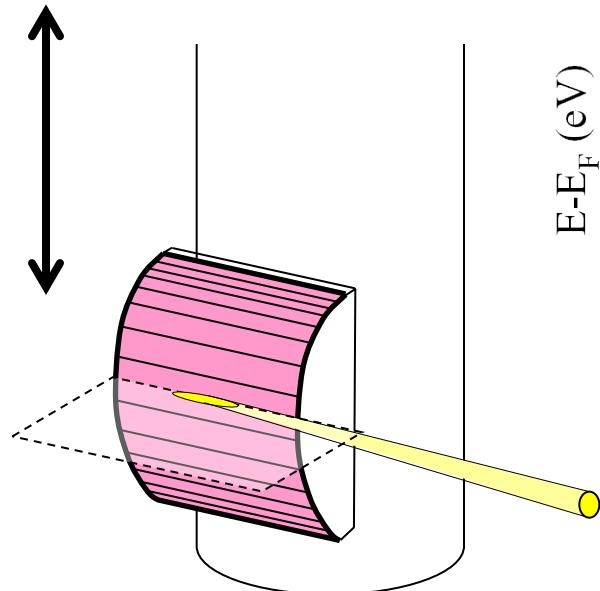
Electronic states at the faceting transition in Au(111)

Z-axis

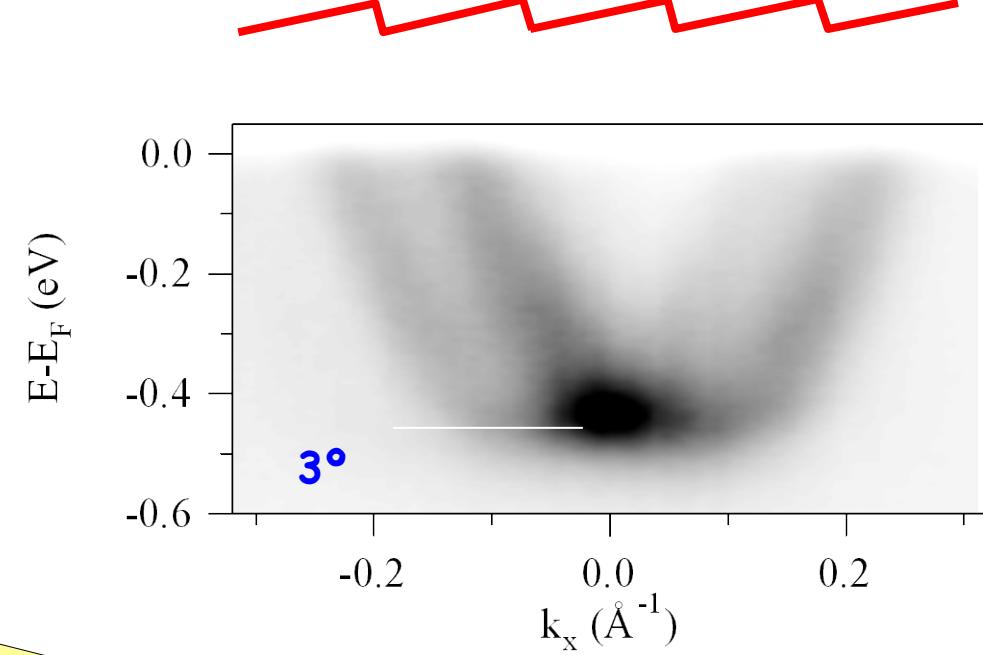
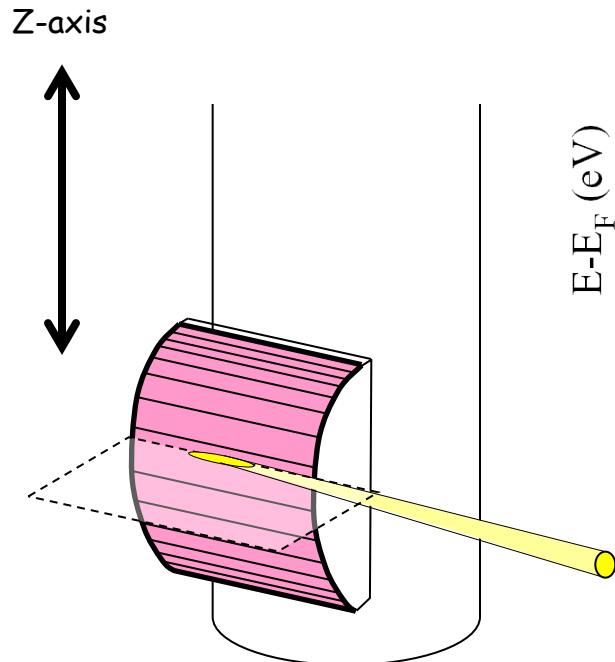


Electronic states at the faceting transition in Au(111)

Z-axis



Electronic states at the faceting transition in Au(111)

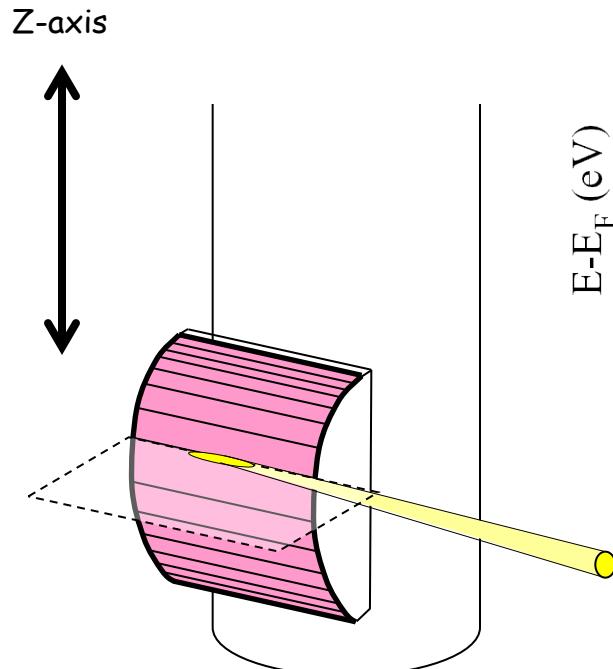


Electronic states at the faceting transition in Au(111)

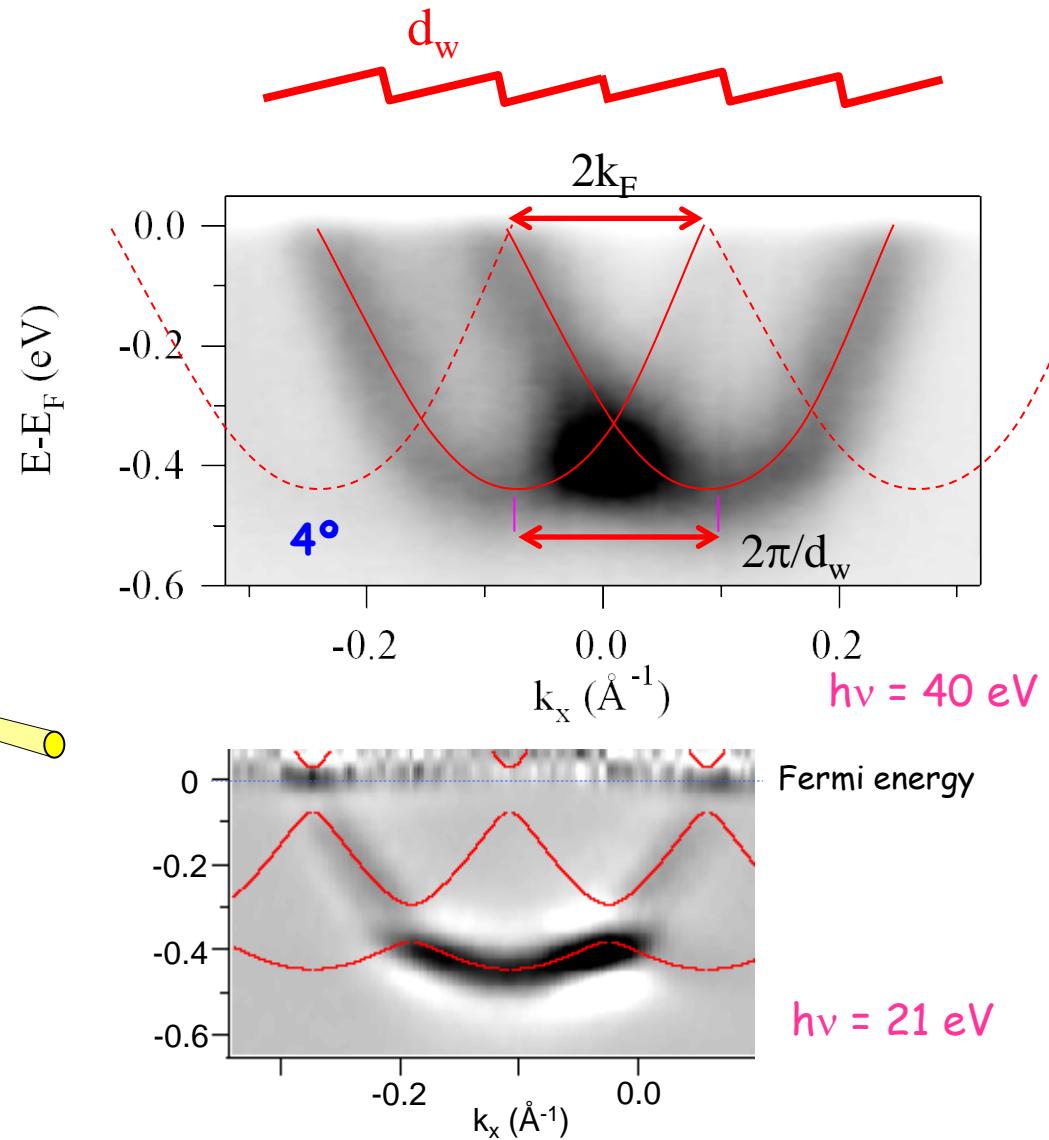
M. Corso *et al.*, J. Phys. Condens. Matter **21**, 353001 (2009).



Lower onset of faceting

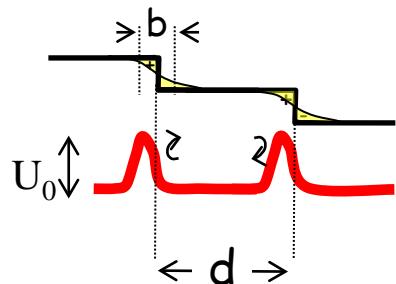


A step lattice Fermi gap opens up at the lower onset of faceting



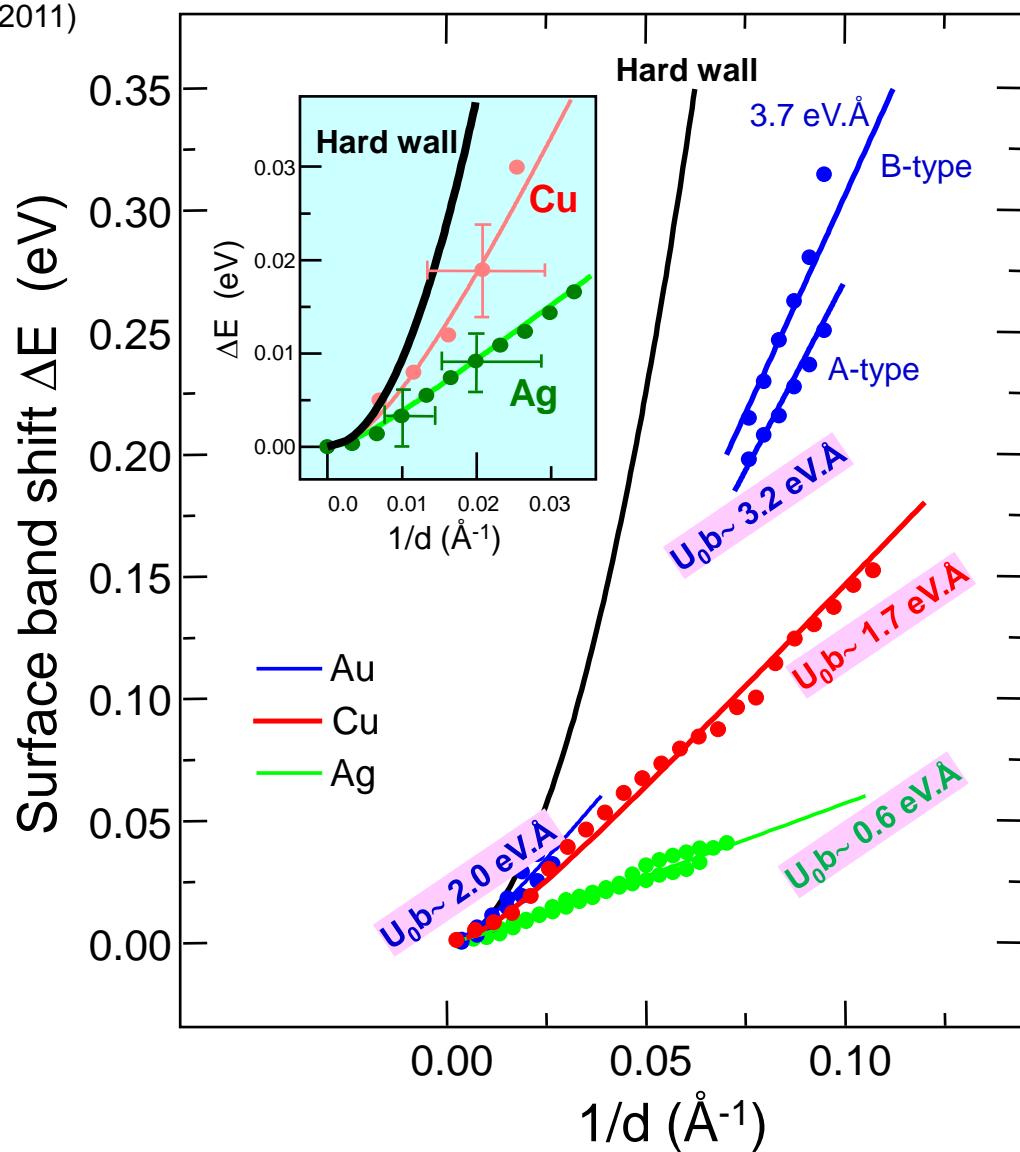
Step scattering through curved surfaces I: terrace-size effect

J. E. Ortega *et al.*, Phys. Rev. B **83**, 085411 (2011)



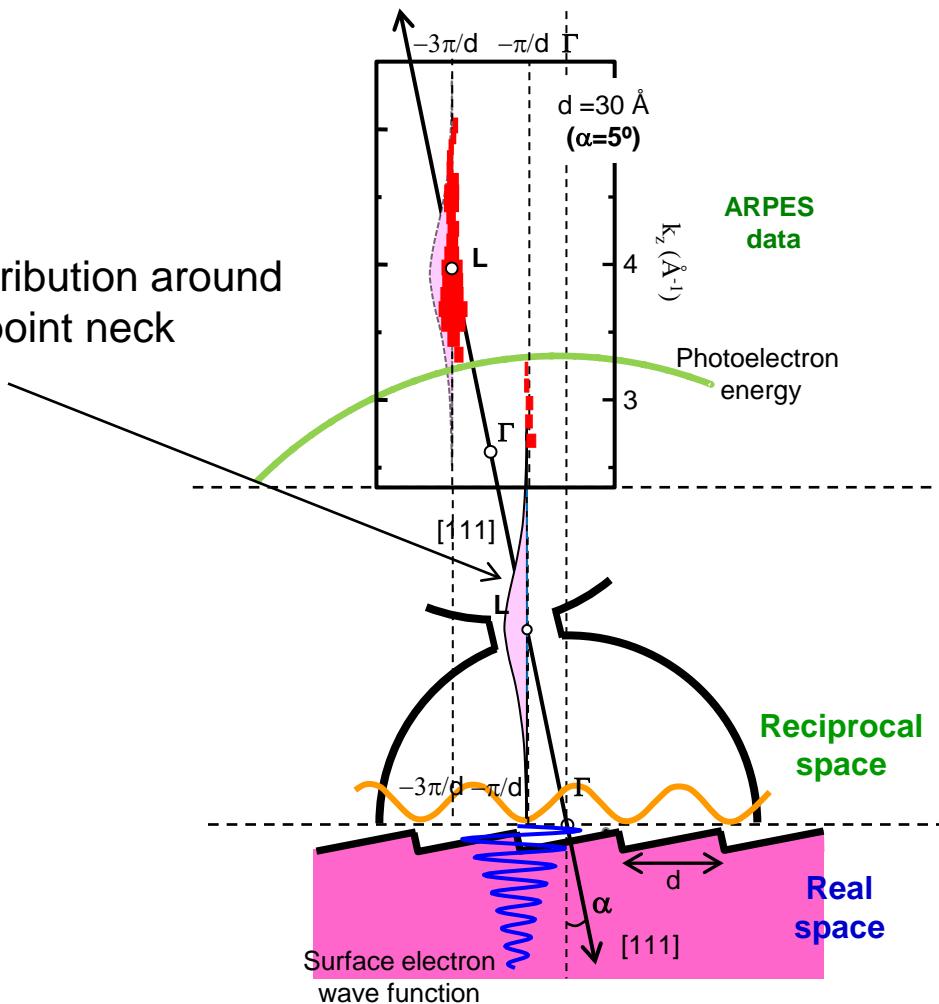
1D Kronig-Penney model

$$\Delta E = \frac{2\hbar^2}{m^*} \frac{1}{d^2} \left(\operatorname{tg}^{-1} \left(\frac{U_0 b}{\sqrt{\frac{2\hbar^2}{m^*} \Delta E}} \right) \right)^2$$

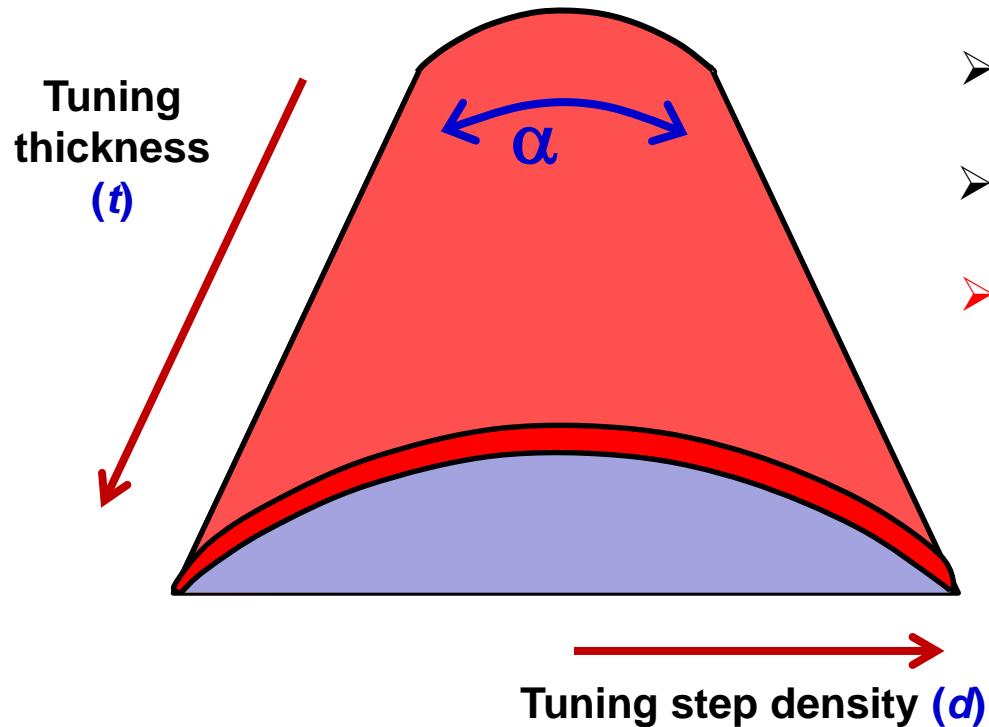


Step scattering through curved surfaces II: wave function plane

Spectral distribution around
the L-point neck

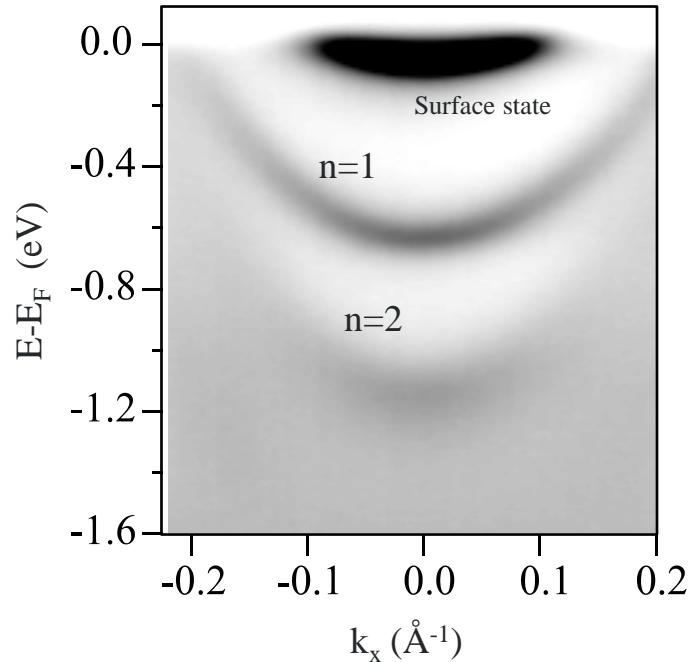
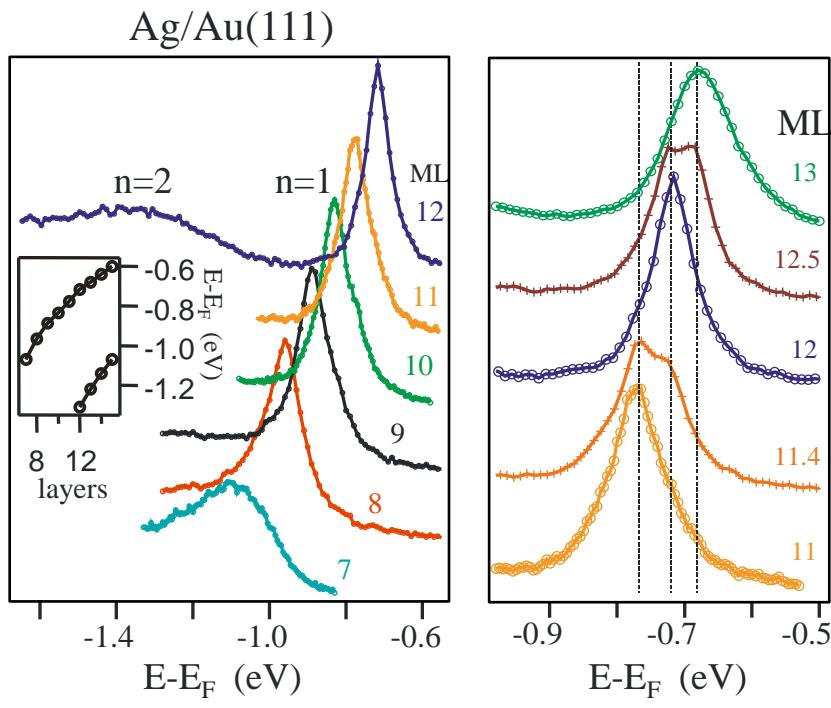
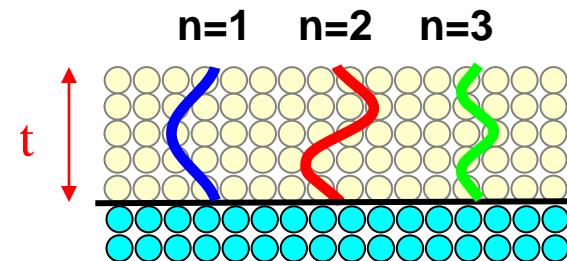
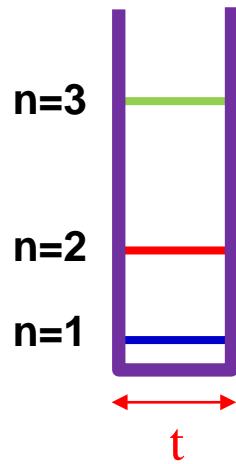
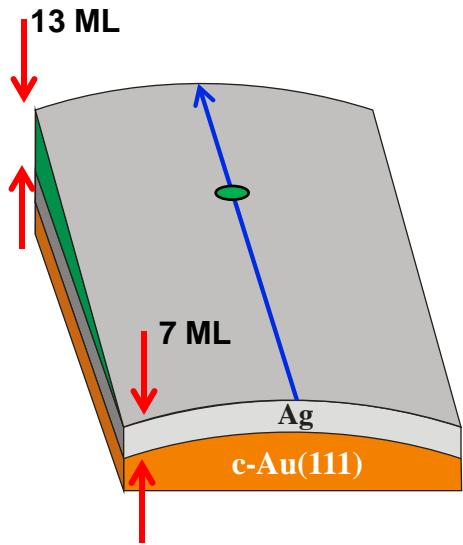


Thin films on curved surfaces

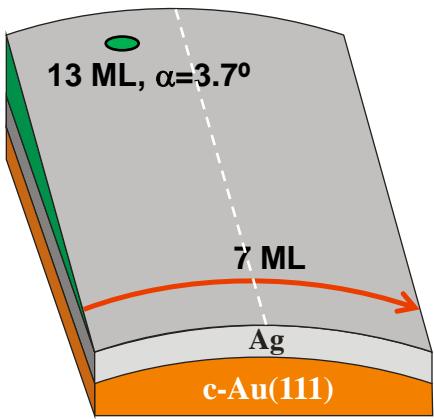


- Ag films on c-Au(111)
- BiAg₂ monolayer on c-Ag(111)
- *Ag monolayer on c-Cu(111)*

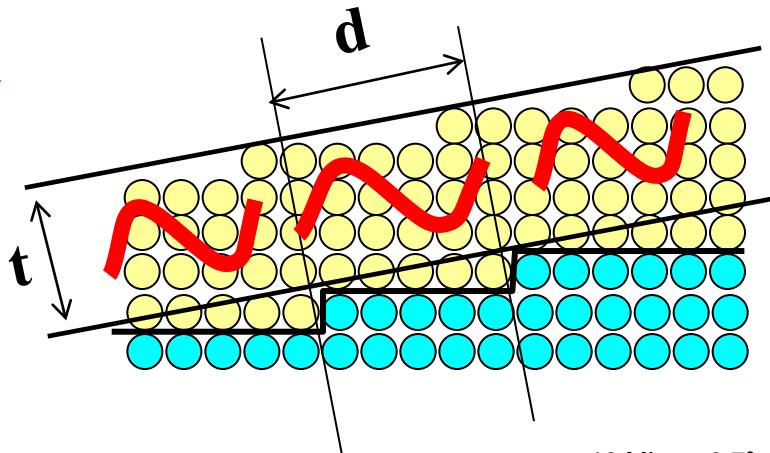
Quantum wells



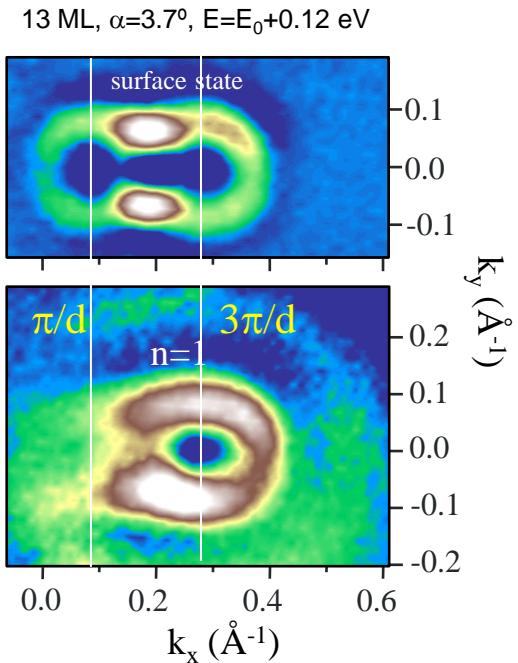
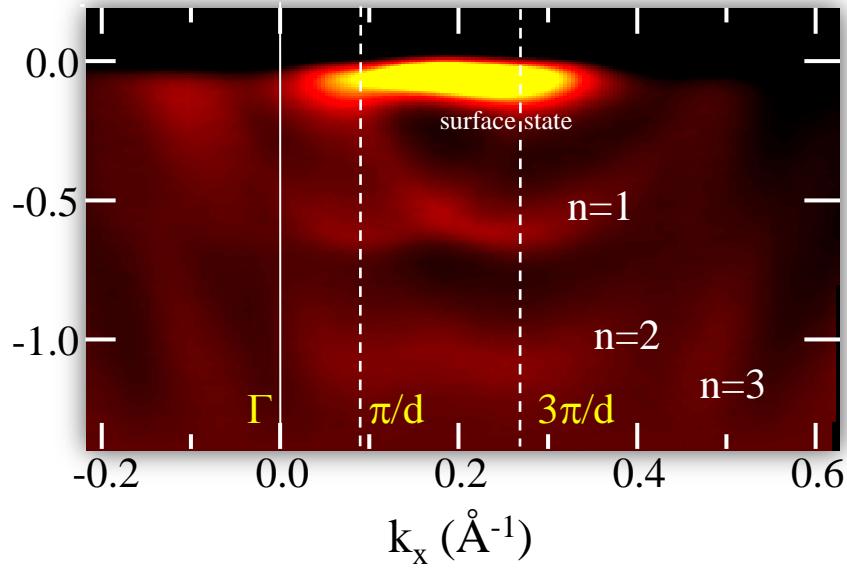
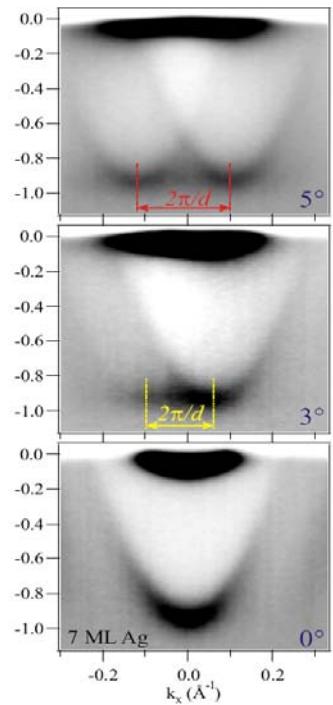
Stepped quantum wells



$h\nu=41$ eV

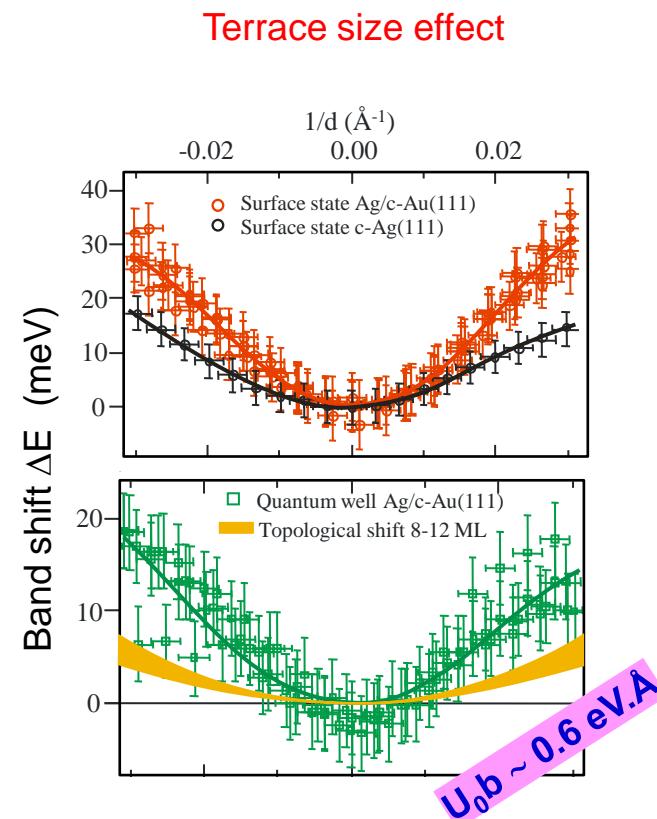
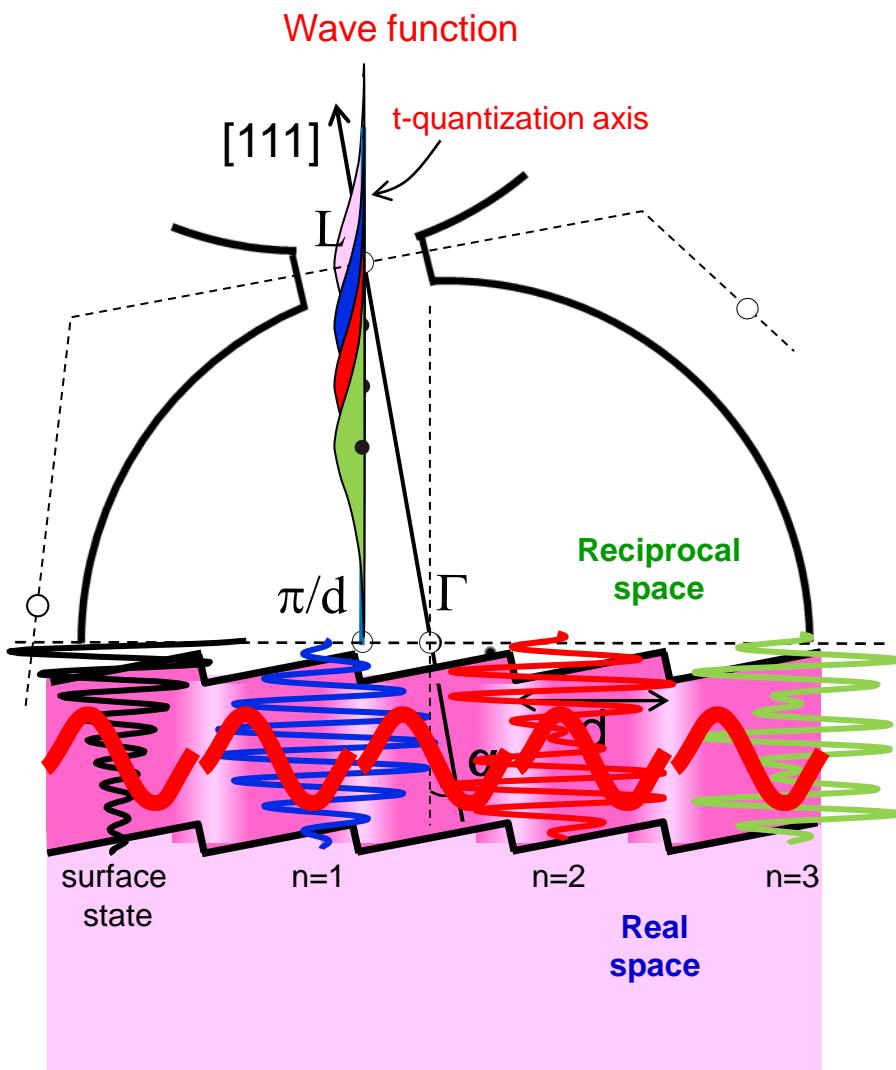


$t=13$ ML, $\alpha=3.7^\circ$ ($d = 36$ Å)

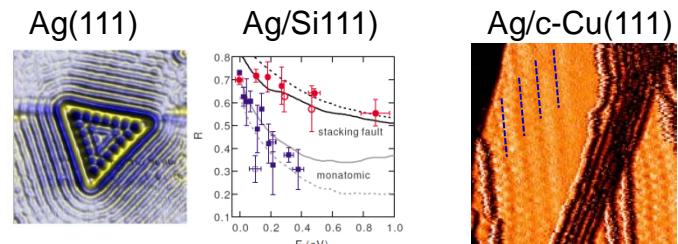


QW and surface states exhibit the same superlattice nature in thin stepped films

Stepped quantum wells



Lateral scattering at stacking faults?



Mugarza et al. (unpublished) and
 PRB 82, 113413 (2010)

Scattering in a Rashba surface alloy

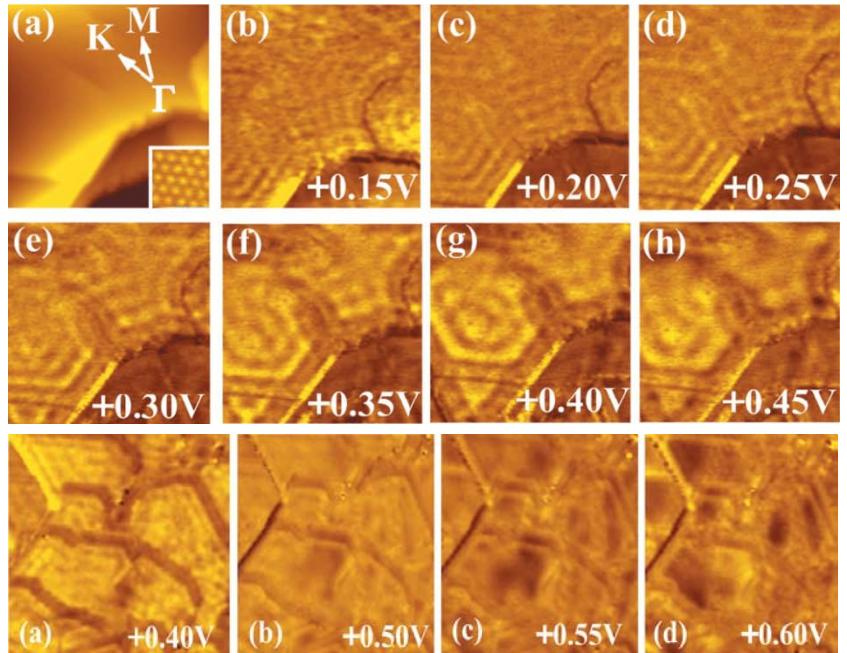
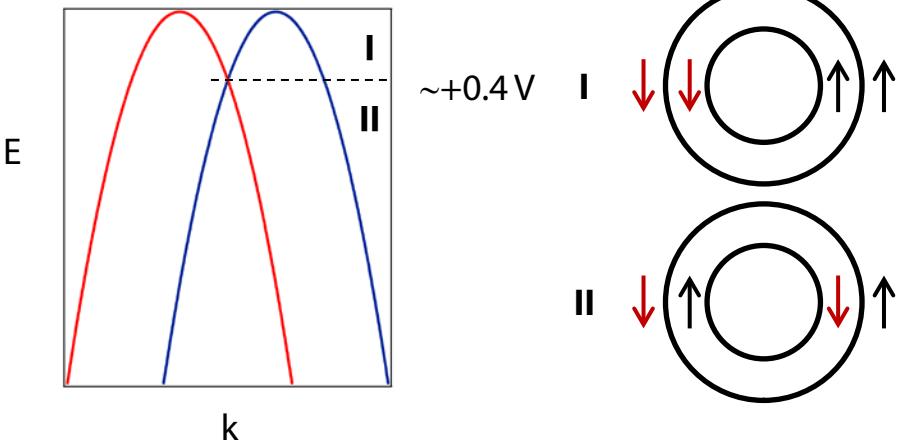
Quantum Interference of Rashba-Type Spin-Split Surface State Electrons

Hiroyuki Hirayama,* Yuki Aoki, and Chiaki Kato

Department of Materials Science and Engineering, Tokyo Institute of Technology,
 J1-3, 4259 Nagatsuda, Midori-ku, Yokohama 226-8502, Japan

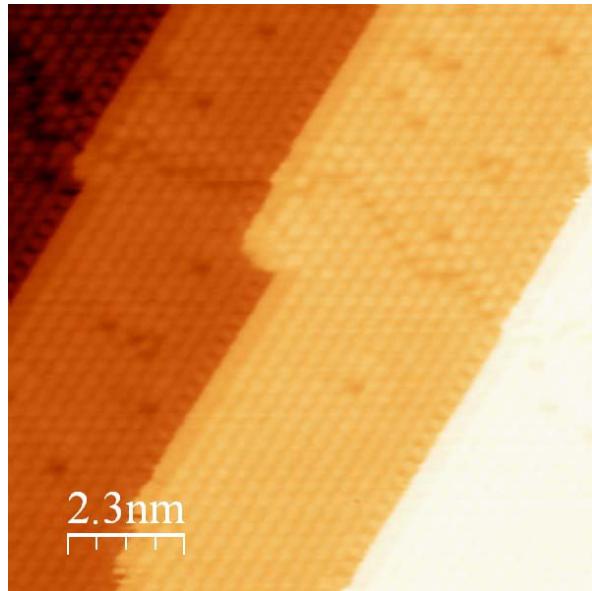
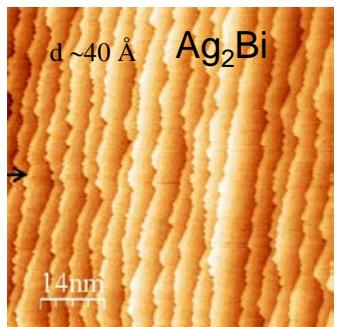
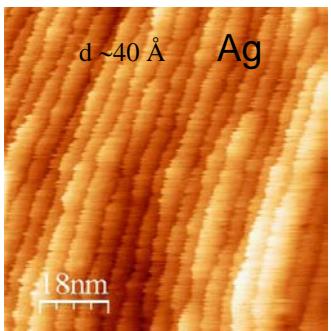
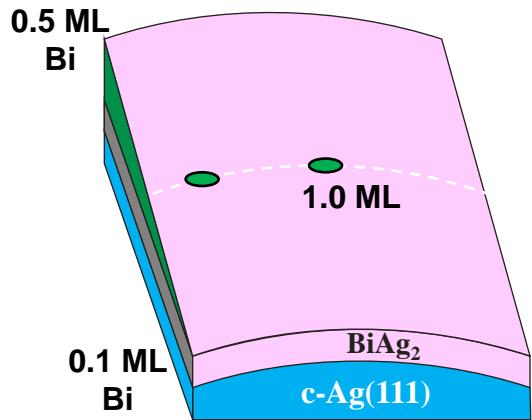
(Received 2 April 2011; published 7 July 2011)

We studied the quantum interference of electrons in the Bi (p_x, p_y) orbital-derived $j = 1/2$ spin-split surface states at Bi/Ag(111) $\sqrt{3} \times \sqrt{3}$ surfaces of 10 monolayer thick Ag(111) films on Si(111) substrates. Surface electron standing waves were observed clearly at the energy (E) below the intersection of the two spin-split downward dispersing parabola bands (E_x). The E dependence of the standing wave pattern reveals the dispersion as the average of the two spin-split surface bands due to the interference between $| (k + \Delta), \uparrow \rangle$ and $| - (k - \Delta), \uparrow \rangle$ [or $| (k - \Delta), \downarrow \rangle$ and $| - (k + \Delta), \downarrow \rangle$] states. In contrast, it was impossible to deduce the dispersion from the standing wave pattern at $E \geq E_x$ because the surface electron cannot find its backscattered state with the same spin polarization.

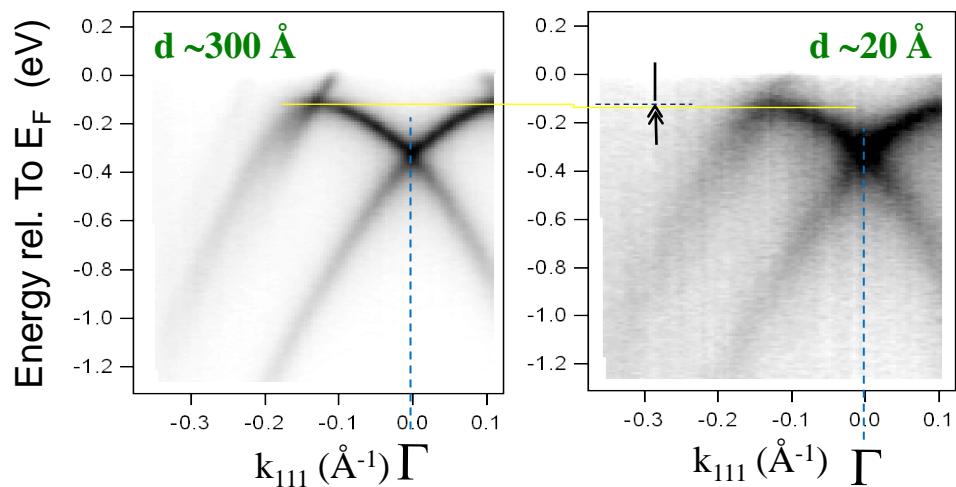


- No standing waves observed in region I ($E > +0.4$ V)
- Standing waves in region II ($E < +0.4$ V)

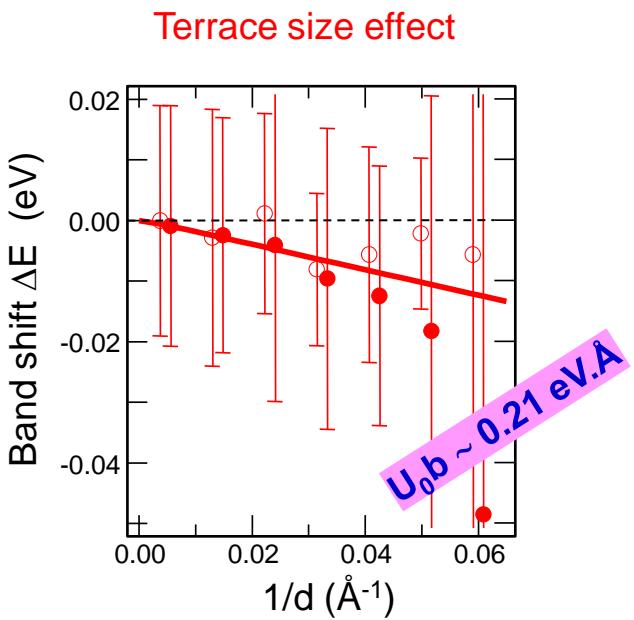
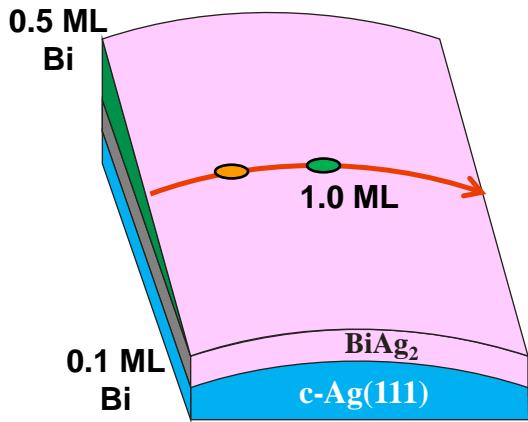
BiAg₂/c-Ag(111)



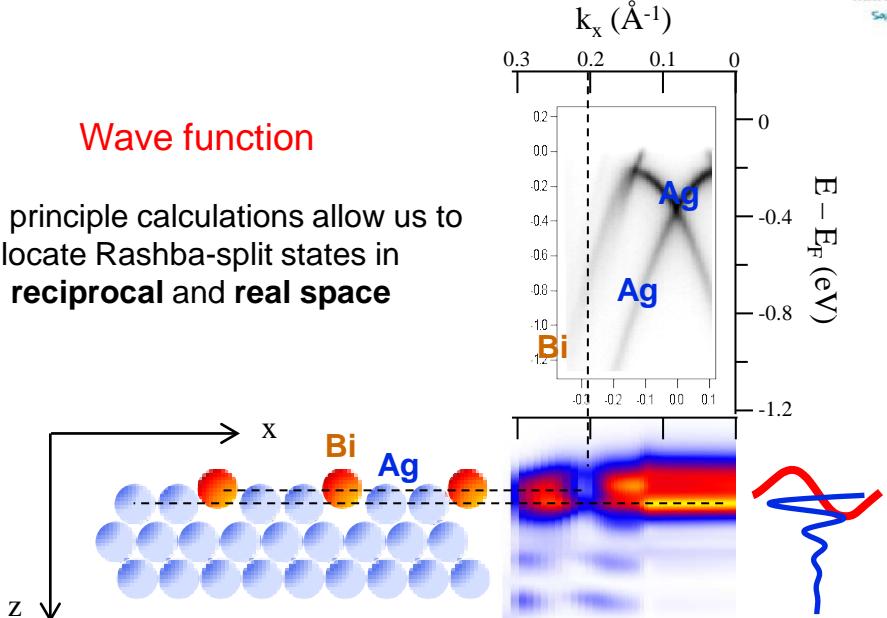
(111)-like band with a small downward shift for a high step density



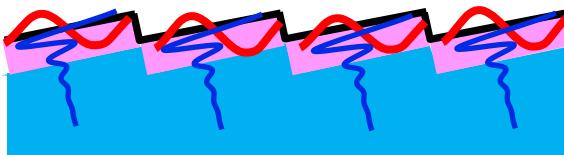
BiAg₂/Ag(111)



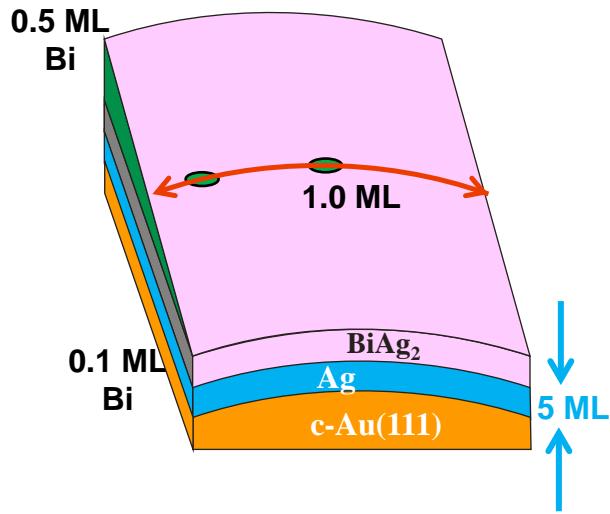
First principle calculations allow us to locate Rashba-split states in **reciprocal** and **real space**



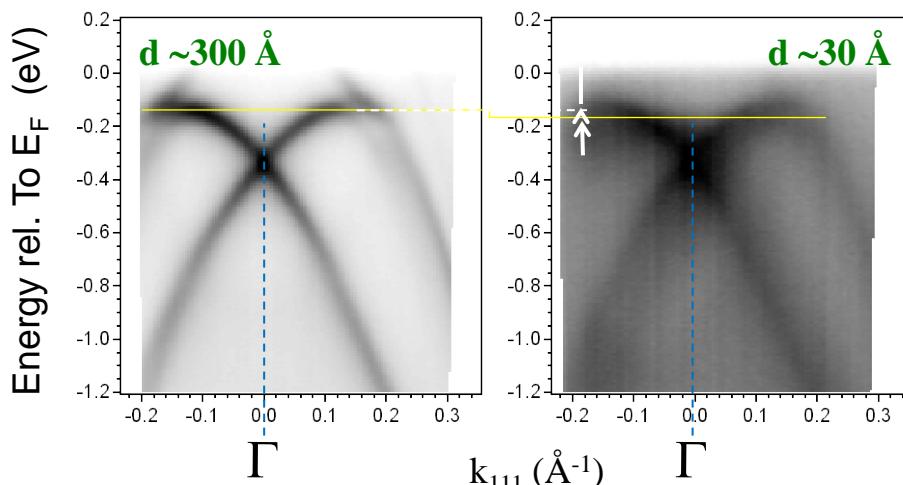
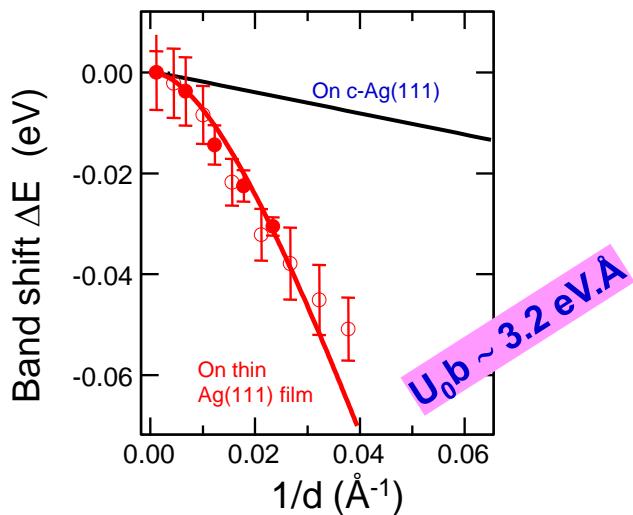
Spin split states around Γ and below E_F are mostly located at Ag atoms under the surface plane
 \Rightarrow weak sensitivity to step dipoles



BiAg₂/Ag/c-Au(111)

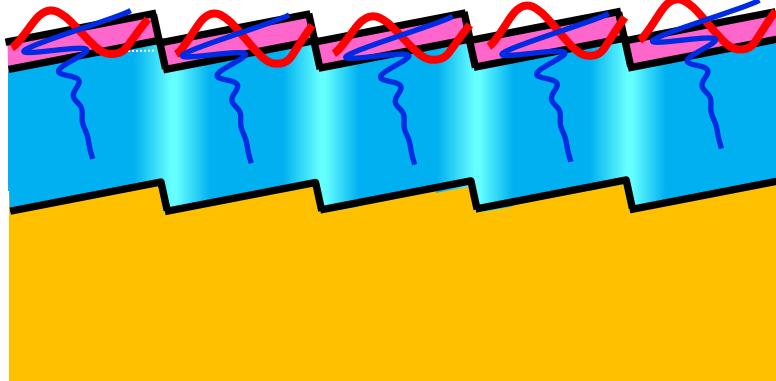


Terrace size effect



Wave function

Lateral scattering increases in BiAg₂ grown on thin films due to vertical stacking faults below surface steps

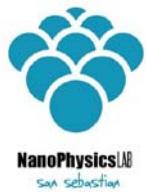


Summary

We probe lateral scattering of electrons, and wave function properties at stepped systems using ARPES

1. *In noble metal surfaces, we observe weak scattering at step lattices.*
2. *QW states of thin stepped Ag films exhibit scattering at stacking fault planes of the film*
3. *Rashba split states of the BiAg₂ monolayer are quasi transparent to monatomic steps, but undergo scattering at stacking faults.*

Acknowledgments



J. Lobo, F. Schiller (Nanophysics Laboratory, San Sebastian)
E. E. Krasovskii (Donostia International Physics Center, San Sebastian)
A. Mugarza, S. Schirone, P. Gambardella (ICN, Barcelona)

Thanks for your attention!!

